

LOOPS IN THE HAMILTONIAN GROUP: A SURVEY

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ABSTRACT. This note describes some recent results about the homotopy properties of Hamiltonian loops. We describe conditions under which a circle action does not contract in the Hamiltonian group, and construct an example of a loop γ of diffeomorphisms of a symplectic manifold M with the property that none of the loops smoothly isotopic to γ preserve any symplectic form on M . We also discuss some new conditions under which the Hamiltonian group has infinite Hofer diameter.

1. INTRODUCTION

Let (M, ω) be a symplectic manifold that is closed, i.e. compact and without boundary. We denote its group of diffeomorphisms by $\text{Diff} := \text{Diff } M$ and by $\text{Symp} := \text{Symp}(M, \omega)$ the subgroup of symplectomorphisms, i.e. diffeomorphisms that preserve the symplectic form. Its identity component $\text{Symp}_0 := \text{Symp}_0(M, \omega)$ has an important normal subgroup $\text{Ham} := \text{Ham}(M, \omega)$ consisting of all symplectomorphisms $\phi \in \text{Symp}_0$ with zero flux, or equivalently, of all time 1 maps ϕ_1^H of Hamiltonian flows ϕ_t^H , where $H : M \times S^1 \rightarrow \mathbb{R}$ is a (smooth) time dependent function. Basic information about these groups may be found in [37, 38, 51] and the survey articles [32, 33]. Note that $\text{Ham} = \text{Symp}_0$ when $H^1(M; \mathbb{R}) = 0$.

This survey is mostly concerned with questions about based Hamiltonian loops, i.e. smooth paths $\{\phi_t\}_{0 \leq t \leq 1}$ in the Hamiltonian group Ham for which $\phi_1 = \phi_0 = \text{id}$. Each such loop is the flow of some time dependent Hamiltonian H_t , and may be reparametrized so that $H_0 = H_1$ and the induced map $H : M \times S^1 \rightarrow \mathbb{R}$ is smooth. If this generating Hamiltonian is time independent then its flow $\phi_t^H, t \in S^1$, is a subgroup of Ham isomorphic to S^1 . (The function $H : M \rightarrow \mathbb{R}$ is then called the *moment map*.) Although these loops are the easiest to understand, there are still many unsolved questions about them.

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Our first group of questions concerns the structure of $\pi_1(\text{Symp})$ and $\pi_1(\text{Ham})$. Note that when $\dim M = 2$, Moser's homotopy argument implies that the symplectomorphism group of M is homotopy equivalent to its group of orientation preserving diffeomorphisms. Thus the homotopy type of the groups Symp and Ham are known. Ham is contractible unless $M = S^2$, in which case both it and Symp_0 are homotopic to $\text{SO}(3)$; while in higher genus Symp_0 is contractible except in the case of the torus, when it is homotopic to the torus $\mathbb{R}^2/\mathbb{Z}^2$. The group $\pi_0(\text{Symp})$ is the well-known mapping class group. Thus the questions listed below are not interesting in this case.

The case when $\dim M = 4$ is also moderately well understood. In particular, every 4-manifold with a Hamiltonian S^1 action is the blow up of a rational or ruled 4-manifold. Moreover the homotopy type of Ham is understood when $M = \mathbb{C}P^2$ or $S^2 \times S^2$ or a one point blow up of such; see for example Gromov [13], Abreu–McDuff [2], Abreu–Granja–Kitchloo [1] and Lalonde–Pinsonnault [25]. For work on nonHamiltonian S^1 actions in dimensions 4 and above see Bouyakoub [4], Duistermaat–Pelayo [5] and Pelayo [47]. There is information on smooth circle actions in 4-dimensions, see Fintushel [9] and Baldrige [3]. But almost nothing is known about the diffeomorphism group of a manifold of dimension ≥ 4 ; for example it is not known whether $\text{Diff } \mathbb{C}P^2$ is homotopy equivalent to the projective unitary group $PU(3)$ as is $\text{Symp}(\mathbb{C}P^2)$. On the other hand it follows from [2] that $\text{Ham}(S^2 \times S^2, \omega)$ is not homotopy equivalent to $\text{Diff}(S^2 \times S^2)$ for any symplectic form $[\omega]$.

Question 1.1. *When does a circle subgroup γ of Symp represent a nonzero element in $\pi_1(\text{Symp})$, or even one of infinite order?*

(Entov–Polterovich [7] call circles of infinite order *incompressible*.)

The first problem here is to decide when a loop is Hamiltonian, i.e. is in the kernel of the Flux homomorphism. Recall from [37] that Flux is defined on the universal cover $\widetilde{\text{Symp}}_0$ of Symp_0 by

$$(1.1) \quad \text{Flux} : \widetilde{\text{Symp}}_0 \rightarrow H^1(M; \mathbb{R}), \quad \tilde{\phi} \mapsto \int_0^1 [\omega(\dot{\phi}_t, \cdot)] dt,$$

where $\tilde{\phi} = (\phi_1, \{\phi_t\}) \in \widetilde{\text{Symp}}_0$. The (symplectic) Flux group Γ_ω is defined to be the image of $\pi_1(\text{Symp})$ under Flux, so that there is an induced homomorphism

$$\text{Flux} : \pi_1(\text{Symp}) \rightarrow H^1(M; \mathbb{R})/\Gamma_\omega$$

with kernel Ham . Ono [45] recently proved that Γ_ω is discrete.

Unfortunately there seem to be no good techniques for understanding when $\text{Flux}(\gamma)$ is trivial. Since Hamiltonian S^1 actions always have fixed points at the critical points of the moment map $H : M \rightarrow \mathbb{R}$, a first guess might be that every symplectic action with fixed points is Hamiltonian. However McDuff [28] shows that this is not the case except in dimension 4. In fact, the following basic problem is still unsolved in dimensions > 4 .

Question 1.2. *Suppose that S^1 acts symplectically on the closed symplectic manifold (M, ω) with a finite but nonzero number of fixed points. Is the action Hamiltonian?*

If the action is semifree (i.e. the stabilizer of a point is either the identity or the whole group) Tolman–Weitsman [57] show that the answer to Question 1.2 is affirmative by computing various equivariant cohomology classes. Some other information on this question has been obtained by Feldman [8] and Pelayo–Tolman [48]. One might hope to extend the Tolman–Weitsman result to semifree actions with more general conditions on the fixed point components, for example that they are simply connected; note that these cannot be arbitrary because of the example in [28] of a semifree but nonHamiltonian action on a 6-manifold with fixed point sets that are 2-tori.

In the current discussion we will largely ignore this problem, for the most part considering only Hamiltonian loops and their images in $\pi_1(\text{Ham})$.

Question 1.3. *To what extent are $\pi_1(\text{Ham})$ and $\pi_1(\text{Symp})$ generated by symplectic S^1 actions?*

This question is a measure of our ignorance. S^1 actions do generate $\pi_1(\text{Symp})$ in very special cases such as $\mathbb{C}P^2$ or $(S^2 \times S^2, pr_1^*(\sigma) + pr_2^*(\sigma))$ (note that the factors have equal area). Indeed in these cases Symp itself is known to have the homotopy type of a compact Lie group (see [32]). However, as we see below, this does not hold in general.

Question 1.4. *What can one say about the relation between $\pi_1(\text{Ham})$, $\pi_1(\text{Symp})$ and $\pi_1(\text{Diff})$? For example, under what circumstances is the map $\pi_1(\text{Symp}) \rightarrow \pi_1(\text{Diff})$ injective or surjective?*

The symplectic Flux group Γ_ω is the quotient of $\pi_1(\text{Symp})$ by $\pi_1(\text{Ham})$ and hence precisely measures their difference. By Kędra–Kotschick–Morita [19], this group vanishes in many cases. Much of their paper in fact applies to the volume¹ flux group Γ_{vol} , which is in principle of a more topological nature than Γ_ω ; it would be interesting to find conditions for the vanishing of Γ_ω that involve symplectic geometry at a deeper level.

In this note, we begin by describing some classical methods for exploring the above questions, the first based on properties of the *action functional* \mathcal{A}_H and the second using volume. These methods give rather good information in the following cases:

- Question 1.1 for toric manifolds (see Corollary 2.4 below)
- Questions 1.3 and 1.4 for pointwise blow ups \widetilde{M} of arbitrary symplectic manifolds M (see Proposition 2.7 and its corollaries).

If M is noncompact and one considers the group $\text{Ham}^c := \text{Ham}^c M$ of compactly supported Hamiltonian symplectomorphisms of M , then there is another classical homomorphism called the *Calabi homomorphism*:

$$\text{Cal} : \pi_1(\text{Ham}^c) \rightarrow \mathbb{R}, \quad \gamma \mapsto \frac{1}{n!} \int_0^1 \left(\int_M H_t \omega^n \right) dt,$$

¹The volume flux is defined by equation (1.1), but ω should be understood as a volume form and the homomorphism takes values in $H^{m-1}(M)$, where $m := \dim M$. Accordingly, Γ_{vol} is the image of $\pi_1(\text{Diff}_{\text{vol}})$ under the flux.

where H_t is the generating Hamiltonian for γ , normalized to have compact support. We explain briefly in Lemma 3.9 why

- the Calabi homomorphism need not vanish.

As we shall see, this question, though classical in origin, is very closely related to questions about the Seidel representation in quantum homology.

One might wonder if it is possible to get better information about the above questions by using more modern (i.e. quantum) techniques. In fact, Question 1.1 first arose in McDuff–Slimowitz [39], a paper that uses Floer theoretic techniques to study paths in Ham that are geodesic with respect to the Hofer norm.² An unexpected consequence of the ideas developed there is that semifree Hamiltonian circle actions do not contract in Ham, though they might have finite order (for example, a rotation of S^2 .) The main tool that has proved useful in this context is an extension of the action homomorphism due to Seidel [55], that is called the *Seidel representation*; see §3.1. This homomorphism assigns to every $\gamma \in \pi_1(\text{Ham})$ a unit (i.e. invertible element) $\mathcal{S}(\gamma)$ in the (small) quantum homology $QH_*(M)$. Corollary 3.2 gives some more results on the above questions obtained using \mathcal{S} . Because \mathcal{S} is usually very hard to calculate, the classical methods often work better in specific examples. Nevertheless, \mathcal{S} is a key tool in other contexts.

One very interesting question is the following. Note that in two dimensions, the only symplectic manifold with a Hamiltonian S^1 action is S^2 , while T^2 has nonHamiltonian actions and higher genus surfaces Σ_g have none.

Question 1.5. *Is there a meaningful extension of the classification of Riemann surfaces into spheres, tori and higher genus to higher dimensional symplectic manifolds? If so, is any aspect of it reflected in the properties of $\pi_1(\text{Ham})$?*

This would be an analog of *minimal model* theory in algebraic geometry. A first step, accomplished by Ruan and his coworkers Hu and T.-J. Li [53, 16], is to understand what it means for two symplectic manifolds to be birationally equivalent. Their results imply that a reasonable class of manifolds to take as the analog of spheres are the *symplectically uniruled* manifolds. These are the manifolds for which there is a nonzero genus zero Gromov–Witten invariant $\langle a_1, a_2, \dots, a_k \rangle_\beta^M$ (for some $k \geq 1$, $a_i \in H_*(M)$ and $\beta \in H_2(M)$) with one of the constraints a_i equal to a point. This class includes all projective manifolds that are uniruled in the sense of algebraic geometry. In this case $\omega(\beta) \neq 0$ and $c_1(\beta) \neq 0$. At the other extreme are the *symplectically aspherical* manifolds for which the restriction $\omega|_{\pi_2(M)}$ of $[\omega]$ to $\pi_2(M)$ vanishes, and possibly also (depending on the author) the restriction $c_1|_{\pi_2(M)}$ of the first Chern class c_1 of (M, ω) . These manifolds have no J -holomorphic curves at all, and hence all nontrivial (i.e. $\beta \neq 0$) Gromov–Witten invariants vanish.

To a first approximation one can characterize symplectically uniruled manifolds in terms of the quantum homology $QH_*(M)$ in the following terms.

²This is defined in §3.2 below.

If (M, ω) is not uniruled then all invertible elements in $QH_{2n}(M)$ have the form $\mathbb{1} \otimes \lambda + x$ where λ is invertible in the coefficient field Λ and $x \in H_{<2n}M \otimes \Lambda$.

(This is nearly an iff statement and can be improved to such: cf. the appendix to [35]. Here $\mathbb{1}$ denotes the fundamental class $[M] \in QH_*(M)$; it is the identity element in $QH_*(M)$. The notation is explained in more detail below.) Therefore another version of the previous question is:

Question 1.6. *To what extent is the structure of the quantum homology ring $QH_*(M)$ reflected in the algebraic/topological/geometric structure of Ham ?*

For example, according to Polterovich [52],

if $\omega|_{\pi_2(M)} = 0$, Ham contains no elements of finite order except for id .

Does this property continue to hold if the condition on ω is weakened to the vanishing of all Gromov–Witten invariants? Since S^1 -manifolds certainly support nontrivial symplectomorphisms of finite order (such as a half turn), it is natural to ask whether these manifolds have nontrivial Gromov–Witten invariants. Since the Seidel element $\mathcal{S}(\gamma)$ is a unit in quantum homology one might also expect to see traces of the uniruled/non-uniruled dichotomy in its properties. This was the guiding idea in my recent proof [35] that

every closed symplectic manifold that supports a Hamiltonian S^1 action is uniruled.

The argument in [35] applies more generally to manifolds with Hamiltonian loops that are nondegenerate Hofer geodesics. This opens up many interesting questions of a more dynamical flavor.

Here we shall discuss the following basic (but unrelated) problem, which is still open in many cases.

Question 1.7. *Does Ham have infinite diameter with respect to the Hofer metric?*

One expects the answer to be positive always. However the proofs for spheres (due to Polterovich [50]) and other Riemann surfaces (due to Lalonde–McDuff [22]) are very different. In fact, as noted by Ostrover [46] one can use the spectral invariants of Schwarz [54] and Oh [43, 44] (see also Usher [58]) to show that the universal cover $\widetilde{\text{Ham}}$ of Ham always has infinite diameter with respect to the induced (pseudo)metric. Therefore the question becomes: when does this result transfer down to Ham ? Schwarz shows in [54] that this happens when ω and c_1 both vanish on $\pi_2(M)$. I recently extended his work in [36], showing that the asymptotic spectral invariants descend to Ham if, for example, all nontrivial genus zero Gromov–Witten invariants vanish and $\text{rank } H_2(M; \mathbb{R}) > 1$. As we explain in §3.2 below, his argument hinges on the properties of the Seidel elements $\mathcal{S}(\gamma)$ of $\gamma \in \pi_1(\text{Ham})$. Here we shall sketch a different extension of his result. In particular we show:

Ham has infinite diameter when M is a “small” one point blow up of $\mathbb{C}P^2$.

This manifold M is of course uniruled (and the spectral invariants do not descend). To my knowledge, it is not yet known whether Ham has infinite diameter for all one point blow ups of $\mathbb{C}P^2$, though it does for $\mathbb{C}P^2$ itself (and indeed for any $\mathbb{C}P^n$) by the results of Entov–Polterovich [6]. For further results on this problem see McDuff [36].

2. CLASSICAL METHODS

We consider two classical methods to detect elements in $\pi_1(\text{Ham})$, the first based on the action homomorphism and the second on considerations of volume. One can also use homological methods as in Kędra–McDuff [20] but these give more information on the higher homotopy groups $\pi_k(\text{Ham})$, $k > 1$.

2.1. The action homomorphism. First of all, we consider Weinstein’s *action homomorphism*

$$\mathcal{A}_\omega : \pi_1(\text{Ham}) \rightarrow \mathbb{R}/\mathcal{P}_\omega,$$

where $\mathcal{P}_\omega := \{\int_c \omega | c \in \pi_2(M)\} \subset \mathbb{R}$ is the group of spherical periods of $[\omega]$. In defining \mathcal{A}_ω , we shall use the following sign conventions. The flow $\{\phi_t^H\}_{t \in [0,1]}$ of a Hamiltonian $H : M \times S^1 \rightarrow \mathbb{R}$ satisfies the equation

$$\omega(\dot{\phi}_t^H(p), \cdot) = -dH_t(\phi_t^H(p)), \quad p \in M,$$

and the corresponding action functional on the space $\mathcal{L}_0 M$ of contractible loops in M is

$$\mathcal{A}_H : \mathcal{L}_0 M \rightarrow \mathbb{R}/\mathcal{P}_\omega, \quad \mathcal{A}_H(x) := - \int_{D^2} u^* \omega + \int_{S^1} H_t(x(t)) dt,$$

where $u : D^2 \rightarrow M$ is any extension of the (contractible) loop $x : S^1 \rightarrow M$, and H_t is assumed to be *mean normalized*, i.e. $\int_M H_t \omega^n = 0$ for all t . Note that any Hamiltonian H_t can be normalized (without changing its flow) by subtracting a suitable *normalization constant* $c_t := \int_M H_t \omega^n / \int_M \omega^n$.

One very important property of \mathcal{A}_H is that its critical points are the closed orbits $x(t) := \phi_t^H(p)$ of the corresponding flow. This is classical: the orbits of a Hamiltonian flow minimize (or, more correctly, are critical points of) the action. If $\{\phi_t^H\}_{t \in S^1}$ is a loop, then all its orbits are closed. Moreover they are contractible.³ Hence these orbits all have the same action. It is not hard to see that this value depends only on the homotopy class $\gamma := [\{\phi_t^H\}]$ of the loop. Call it $\mathcal{A}_\omega(\gamma)$. We therefore get a map $\mathcal{A}_\omega : \pi_1(\text{Ham}) \rightarrow \mathbb{R}/\mathcal{P}_\omega$, which is easily seen to be a homomorphism.

If the loop $\{\phi_t\}_{t \in S^1}$ is the circle subgroup γ_K of $\text{Ham}(M, \omega)$ generated by the mean normalized Hamiltonian $K : M \rightarrow \mathbb{R}$ then we may take p to be a critical point of K and u to be the constant disc. Hence $\mathcal{A}_\omega(\gamma_K)$ is the image in $\mathbb{R}/\mathcal{P}_\omega$ of any critical value of K . This is well defined because the difference $K(p) - K(q)$ between two critical values

³This folk theorem is proved, for example, in [38, Ch 9.1] in the case when (M, ω) is semipositive. The general case follows from the proof of the Arnold conjecture: see [10, 26]. A simpler proof is sketched in [33]. Since all known proofs use quantum methods the existence of the homomorphism \mathcal{A}_ω is not entirely “classical” in the case when $H^1(M) \neq \{0\}$. However, toric manifolds are simply connected, and so our discussion of that case is “classical”.

is the integral of ω over the 2-sphere formed by rotating an arc from p to q by the S^1 -action.

Polterovich noticed that when (M, ω) is spherically monotone (i.e. there is a constant κ such that $\kappa[\omega]$ and $c_1(M)$ induce the same homomorphism $\pi_2(M) \rightarrow \mathbb{R}$) then one can use the Maslov index of a fixed point to get rid of the indeterminacy of \mathcal{A}_ω , hence lifting it to the *combined Action–Maslov homomorphism* $\pi_1(\text{Ham}) \rightarrow \mathbb{R}$; cf. [49, 7].

In general, the action homomorphism is hard to calculate because there are few good ways of understanding the normalization constant for an arbitrary Hamiltonian. However, as pointed out in McDuff–Tolman [41], this calculation is possible in the toric case and, as we now explain, one gets quite useful information from it.

The toric case.

A $2n$ -dimensional symplectic manifold is said to be *toric* if it supports an effective Hamiltonian action of the standard torus T^n . The image of the corresponding moment map $\Phi : M \rightarrow \mathbb{R}^n$ is a convex polytope Δ (called the *moment polytope*) that, as shown by Delzant, completely determines the initial symplectic manifold which we therefore denote $(M_\Delta, \omega_\Delta)$. These manifolds have a compatible complex structure J_Δ , also determined by Δ , and hence are Kähler; cf. Guillemin [14]. Denote by $\text{Isom}_0 M_\Delta$ the identity component of the corresponding Lie group of Kähler isometries. Then $T^n \subseteq \text{Isom}_0 M_\Delta$ and it is natural to try to understand the resulting homomorphisms

$$(2.1) \quad \pi_1(T^n) \rightarrow \pi_1(\text{Isom}_0 M_\Delta) \rightarrow \pi_1(\text{Ham } M_\Delta).$$

For example, if $M = \mathbb{C}P^n$ (complex projective space), Δ is an n -simplex and $\text{Isom}_0 M_\Delta$ is the projective unitary group $= PU(n+1)$. Hence the first map has image $\mathbb{Z}/(n+1)\mathbb{Z}$. Seidel pointed out in his thesis [55] that this finite subgroup $\mathbb{Z}/(n+1)\mathbb{Z}$ injects into $\pi_1(\text{Ham } \mathbb{C}P^n)$. We will prove this here by calculating the action homomorphism.⁴

Denote by \mathfrak{t} the Lie algebra of T^n , and by ℓ its integral lattice. Each vector $H \in \ell$ exponentiates to a circle subgroup $\gamma_H \subset T^n$. Strictly speaking, the moment map should be considered to take values in the dual \mathfrak{t}^* of the Lie algebra \mathfrak{t} , and its definition implies that the corresponding S^1 action on M is generated by the function

$$p \mapsto \langle H, \Phi(p) \rangle, \quad p \in M,$$

where $\langle \cdot, \cdot \rangle$ denotes the natural pairing between \mathfrak{t} and its dual $\mathfrak{t}^* \supset \text{Im } \Phi$. The corresponding mean normalized Hamiltonian is

$$K := \langle H, \Phi \rangle - \langle H, c_\Delta \rangle,$$

where c_Δ denotes the center of mass of Δ . Since vertices of Δ correspond to fixed points in M_Δ , this implies that

$$\mathcal{A}_\omega(\gamma_H) = \langle H, v \rangle - \langle H, c_\Delta \rangle \in \mathbb{R}/\mathcal{P}_\omega,$$

where v is any vertex of Δ .

To go further, recall that the moment polytope Δ is *rational*, that is, the outward conormals to the facets $F_i, 1 \leq i \leq N$, of Δ are rational and so have unique primitive

⁴Although our proof looks different from Seidel's, it is essentially the same; cf. the description (given below) of $\mathcal{A}_\omega(\gamma)$ in terms of the fibration $P_\gamma \rightarrow S^2$.

representatives $\eta_i \in \ell \subset \mathfrak{t}$, $1 \leq i \leq N$. (This implies that slopes of the edges of Δ are rational, but its vertices need not be.) Thus Δ may be described as the solution set of a system of linear inequalities:

$$\Delta = \bigcap_{i=1}^N \{x \in \mathfrak{t}^* : \langle \eta_i, x \rangle \leq \kappa_i\}.$$

The constants $(\kappa_1, \dots, \kappa_N)$ are called the *support numbers* of Δ and determine the symplectic form ω_Δ . They may be slightly varied without changing the diffeomorphism type of M_Δ . We shall write $\kappa := (\kappa_1, \dots, \kappa_N)$. Thus we may think of the moment polytope $\Delta := \Delta(\kappa)$ and its center of gravity $c_\Delta(\kappa)$ as functions of κ . An element $H \in \ell$ is said to be a *mass linear function* on Δ if the quantity $\langle H, c_\Delta(\kappa) \rangle$ depends linearly on κ .

Here is a foundational result from [41]. The proof that the α_i are integers uses the fact that M_Δ is smooth, i.e. for each vertex of Δ the conormals of the facets meeting at that point form a basis for the integer lattice ℓ .

Proposition 2.1. *Let $H \in \ell$. If γ_H contracts in $\text{Ham}(M_\Delta, \omega_{\Delta(\kappa_0)})$ then H is mass linear. More precisely, there are integers α_i such that for all κ sufficiently near κ_0*

$$\langle H, c_\Delta(\kappa) \rangle = \sum \alpha_i \kappa_i.$$

Sketch of proof. Clearly if γ_H vanishes in $\pi_1(\text{Ham } M_\Delta)$, then $\mathcal{A}_\omega(\gamma_H) = 0$. It is immediate from Moser's homotopy argument that if γ_H contracts in $\text{Ham}(M_\Delta, \omega_\Delta)$, then it also contracts for all sufficiently small perturbations of ω_Δ . Since varying κ corresponds to varying the symplectic form on M_Δ , for any vertex v of Δ the image of γ_H under the action homomorphism

$$\mathcal{A}_{\omega(\kappa)}(\gamma_H) = \langle H, v \rangle - \langle H, c_\Delta(\kappa) \rangle$$

must lie in $\mathcal{P}_{\omega(\kappa)}$ for all κ sufficiently close to $\kappa_0 \in \mathbb{R}^N$. But $\mathcal{P}_{\omega(\kappa)}$ is generated by the (affine) lengths⁵ of the edges of the polytope Δ , and so it is a finitely generated subgroup of \mathbb{R} whose generators are linear functions of the κ_i with integer coefficients. Similarly, for each vertex $v = v(\kappa)$ of $\Delta(\kappa)$ the function $\langle H, v \rangle$ is linear with respect to the κ_i with integer coefficients. Since the function $\kappa \mapsto \mathcal{A}_{\omega(\kappa)}(\gamma_H)$ is continuous, it follows that the function $\kappa \mapsto \langle H, c_\Delta(\kappa) \rangle$ is also a linear function of the κ_i with integer coefficients as κ varies in some open set. \square

We now can sketch the proof of the result mentioned just after equation (2.1) that $\pi_1(PU(n+1))$ injects into $\pi_1(\text{Ham } \mathbb{C}P^n)$.

Corollary 2.2. *If $M_\Delta = \mathbb{C}P^n$, the map $\pi_1(\text{Isom}_0 M_\Delta) \rightarrow \pi_1(\text{Ham } M_\Delta)$ is injective.*

⁵Since Δ is rational its edges e have rational slopes and so for each e there is an integral affine transformation of $(\mathfrak{t}^*, \ell^*) \cong (\mathbb{R}^n, \mathbb{Z}^n)$ taking it to the x_1 axis. The affine length of e is defined to be the usual length of its image on the x_1 -axis.

Proof. Since Δ is a simplex, we can take it to be $\Delta := \{x \in \mathbb{R}^n | x_i \geq 0, \sum x_i = 1\}$ with edges of affine length 1. Hence $\mathcal{P}_\omega = \mathbb{Z}$ and the center of gravity is $(\frac{1}{n+1}, \dots, \frac{1}{n+1})$. Next, note that $\pi_1(\text{Isom}_0 M_\Delta)$ is generated by the circle action

$$\gamma : [z_0 : z_1 : \dots : z_n] \mapsto [e^{-2\pi i \theta} z_0 : z_1 : \dots : z_n].$$

The corresponding vector $H \in \mathfrak{t}$ is one of the outward conormals of the moment polytope Δ . One can easily check that $\langle H, c_\Delta \rangle = -\frac{1}{n+1} \in \mathbb{R}/\mathcal{P}_\omega = \mathbb{R}/\mathbb{Z}$, which has order $n+1$. Hence the order of γ in $\pi_1(\text{Ham})$ is divisible by $n+1$, and so the map must be injective. \square

It turns out that most moment polytopes do not admit mass linear H . More precisely, we call a facet *pervasive* if it meets all the other facets and *flat* if the conormals of all the facets meeting it lie in a hyperplane in \mathfrak{t} . Note that one can destroy pervasive and flat facets by suitable blow ups.

Proposition 2.3 ([41]). *Suppose that Δ is a moment polytope with no pervasive or flat facets. Then every mass linear function H on Δ vanishes.*

Corollary 2.4. *If Δ has no pervasive or flat facets, the map $\pi_1(T^n) \rightarrow \pi_1(\text{Ham } M_\Delta)$ is injective.*

On the other hand we also show:

Proposition 2.5 ([41]). *The only toric manifolds for which the image of $\pi_1(T^n)$ in $\pi_1(\text{Ham } M_\Delta)$ is finite are products of projective spaces.*

It is also easy to characterize mass linear H corresponding to circles γ_H that contract in $\text{Isom}_0 M_\Delta$. (Such H are called *inessential*. The other mass linear functions are called *essential*.) The papers [41, 42] classify all pairs (Δ, H) consisting of a moment polytope Δ of dimension less than or equal to four together with an essential mass linear H on it. It turns out that there are many interesting families of examples when $n = 4$. However when $n \leq 3$ there is only one.

Proposition 2.6. *Let $(M_\Delta, \omega_\Delta)$ be a toric manifold of dimension $2n \leq 6$ such that $\pi_1(\text{Isom}_0 M_\Delta) \rightarrow \pi_1(\text{Ham } M_\Delta)$ is not injective. Then M is a $\mathbb{C}P^2$ -bundle over $\mathbb{C}P^1$.*

2.2. Applications using Volume. This section is based on the paper [34], which uses various methods, both classical and quantum, to explore the homotopy groups of $\text{Ham } \widetilde{M}$, where \widetilde{M} is the one point blow up of M . We shall concentrate here on questions concerning $\pi_1(\text{Ham } \widetilde{M})$. Our arguments do not use the geometry provided by the form ω , but just the fact that its cohomology class $a := [\omega]$ has $a^n \neq 0$. Thus in this section we will suppose that (M, a) is a *c-symplectic manifold* (short for cohomologically symplectic), which simply means that $a \in H^2(M)$ has the property that $a^n \neq 0$, where, as usual, $2n := \dim M$. Note that such a manifold is oriented, and so for each point $p \in M$ has a well defined blow up \widetilde{M} at p obtained by choosing a complex structure near p that is compatible with the orientation and then performing the usual complex blow up at p . If ω is a closed form on M that is symplectic in the

neighborhood U_p of p , then one can obtain a family of closed forms $\tilde{\omega}_\varepsilon$ on \tilde{M} for small $\varepsilon > 0$, by thinking of the blow up as the manifold obtained from (M, ω) by cutting out a symplectic ball in U_p of radius $\sqrt{\varepsilon/\pi}$ and identifying the boundary of this ball with the exceptional divisor via the Hopf map: see [37].

Here is a simplified version of one of the main results in [34].

Proposition 2.7. *Let (M, a) be a c -symplectic manifold. Then there is a homomorphism $\tilde{f}_* : \pi_2(M) \oplus \mathbb{Z} \rightarrow \pi_1(\text{Diff } \tilde{M})$ whose kernel is contained in the kernel of the rational Hurewicz homomorphism $\pi_2(M) \rightarrow H_2(M; \mathbb{Q})$.*

This result has several consequences that throw light on Questions 1.3 and 1.4. The first was observed by Keřdra [18] (and was one of the motivations for [34]).

Corollary 2.8 ([18]). *There are symplectic 4-manifolds (M, ω) such that $\pi_1(\text{Ham } M)$ is nonzero, but that do not support any S^1 action.*

Proof. There are many blow up manifolds that do not admit circle actions. For example, in dimension 4 Baldrige [3] has shown that if X has $b^+ > 1$ and admits a circle action with a fixed point then its Seiberg–Witten invariants must vanish. Since manifolds that admit fixed point free circle actions must have zero Euler characteristic, this implies that no simply connected Kähler surface with $b^+ > 1$ admits a circle action. Thus the blow up of a K3 surface has no circle actions, but does have nontrivial $\pi_1(\text{Ham})$ and $\pi_1(\text{Diff})$. \square

When M is symplectic rather than c -symplectic, the map \tilde{f}_* of Proposition 2.1 takes the factor $\pi_2(M)$ into $\pi_1(\text{Ham})$. However, the elements in the image of the \mathbb{Z} factor are constructed using a twisted blow up construction and so most probably do not lie in $\pi_1(\text{Ham})$ for any ω . When $\dim M = 4$ there are some cases when one can actually prove this.

Corollary 2.9. *Let X be the blow up of $\mathbb{C}P^2$ at one point or the blow up of T^4 at k points for some $k \geq 1$. Then $\pi_1(\text{Diff } X)$ is not generated by the images of $\pi_1(\text{Symp}(X, \omega))$, as ω varies over the space of all symplectic forms on X .*

This result should be compared with Gromov’s observation in [13] that the map $\pi_1(\text{Ham}(M, \omega)) \rightarrow \pi_1(\text{Diff})$ is not surjective when $M = S^2 \times S^2$ and each sphere factor has the same ω -area. Seidel [56] extended this to $\mathbb{C}P^m \times \mathbb{C}P^n$. However, the elements in $\pi_1(\text{Diff})$ that they consider are in the image of $\pi_1(\text{Ham}(M, \omega'))$ for some other ω' , and so they did not establish the stronger statement given above.

We now sketch the proof of Proposition 2.7 in the case when $H^1(M; \mathbb{R}) = 0$. The first task is to define the map $\tilde{f}_* : \pi_2(M) \oplus \mathbb{Z} \rightarrow \pi_1(\text{Diff } \tilde{M})$. We shall identify $\pi_1(\text{Diff } \tilde{M})$ with $\pi_2(B\text{Diff } \tilde{M})$, i.e. for each pair $(\alpha, \ell) \in \pi_2(M) \oplus \mathbb{Z}$ we shall construct a smooth bundle $P_{(\alpha, \ell)} \rightarrow S^2$ with fiber \tilde{M} , defining $\tilde{f}_*(\alpha, \ell)$ to be the homotopy class of the corresponding clutching function $S^1 \rightarrow \text{Diff } \tilde{M}$.

First suppose that (M, ω) is symplectic. Then the diagonal is a symplectic submanifold of $(M \times M, \omega \oplus \omega)$ and so we can construct a symplectic bundle $(\tilde{Q}, \tilde{\Omega}_\varepsilon) \rightarrow M$

with fiber $(\widetilde{M}, \widetilde{\omega}_\varepsilon)$ by blowing up normal to the diagonal with weight ε . The bundle $\widetilde{P}_\alpha \rightarrow S^2$ is then defined to be the pullback of $\widetilde{Q} \rightarrow M$ by $\alpha : S^2 \rightarrow M$:

$$\begin{array}{ccc} \widetilde{P}_\alpha & \longrightarrow & \widetilde{Q} \\ \downarrow & & \downarrow \\ S^2 & \xrightarrow{\alpha} & M. \end{array}$$

Thus we get a homomorphism

$$\widetilde{f}_* : \pi_2(M) \rightarrow \pi_1(\text{Symp}(\widetilde{M}, \widetilde{\omega}_\varepsilon)).$$

(This construction is basically due to Kędra [17]. The image of the map actually lies in $\pi_1(\text{Ham})$ since $\widetilde{\omega}_\varepsilon$ has the closed extension $\widetilde{\Omega}_\varepsilon$; cf. Lalonde–McDuff [23].)

Alternatively, instead of constructing the universal model, we can start with the product

$$(S^2 \times M, \Omega), \quad \text{where } \Omega = pr_1^*(\beta) + pr_2^*\omega,$$

for some area form β on S^2 and then blow this up along the graph $gr_\alpha := \{(z, \alpha(z)), z \in S^2\}$ of α . To make this work we just need to choose β so that Ω is symplectic near gr_α and restricts on it to a nondegenerate form. In particular, we do *not* need ω to be symplectic, and so can take it to be any representative of the class $a \in H^2(M)$. But then there is another degree of freedom. If Ω is not symplectic everywhere, there is no longer a canonical choice of complex structure on the normal bundle ν_α to gr_α . In fact, there is \mathbb{Z} s worth of choices of complex structure J on ν_α .

We define $\widetilde{f}_*(\alpha, \ell)$ to be the element of $\pi_1(\text{Diff } \widetilde{M})$ obtained by blowing up along gr_α with complex structure chosen so that $c_1(\nu_\alpha, J) = \ell$. (This only depends on ℓ not on the particular choice of J .) For example, we get a “nonsymplectic” element in $\pi_1(\widetilde{M})$ by blowing up $S^2 \times M$ along the trivial section $S^2 \times \{pt\}$ using a nontrivial complex structure. (When $M = \mathbb{C}P^2$ we prove Corollary 2.9 by showing that the corresponding loop is not homotopic to any loop of symplectomorphisms.)

Proof of Proposition 2.7. If α is not in the kernel of the Hurewicz map, there is some class in $H^2(M)$ that does not vanish on it. Hence, by slightly perturbing a if necessary we may suppose that $\lambda := a(\alpha) \neq 0$. The bundle $\widetilde{P}_{\alpha, \ell} \rightarrow S^2$ constructed above is equipped with a closed form $\widetilde{\omega}_\varepsilon$ that is symplectic near the exceptional divisor. We show below that

$$(2.2) \quad \text{vol}(\widetilde{P}_{\alpha, \ell}, \widetilde{\Omega}_\varepsilon) = \mu_0 V - v_\varepsilon \left(\mu_0 + \lambda - \frac{\ell}{n+1} \varepsilon \right),$$

where λ is as above, $\mu_0 = \int_{S^2} \beta$ is the area of the base with respect to the chosen area form β , $V = \frac{1}{n!} \int_M a^n$ is the volume of (M, a) , and $v_\varepsilon := \frac{\varepsilon^n}{n!}$ is the volume of a ball of capacity ε . Thus $\mu_0 V$ is the volume of $(S^2 \times M, \Omega)$.

Now observe that the underlying smooth bundle $\widetilde{P} \rightarrow S^2$ does not depend on the choice of ε . Therefore, if $\widetilde{P}_{\alpha, \ell} \rightarrow S^2$ were a smoothly trivial fibration, then, for all ε , the volume of $\widetilde{P}_{\alpha, \ell}$ would be the product of $V - v_\varepsilon := \text{vol}(\widetilde{M}, \widetilde{\omega}_\varepsilon)$ with the “size” μ of the base. Since μ could be measured by integrating $\widetilde{\Omega}_\varepsilon$ over a section of \widetilde{P} which is

the same for all ε , $\mu = \mu_1 + k\varepsilon$ would be a linear function⁶ of ε . Therefore, the two polynomial functions $\text{vol}(\tilde{P}_{\alpha,\ell}, \tilde{\Omega}_\varepsilon)$ and $(V - v_\varepsilon)(\mu_1 + k\varepsilon)$ would have to be equal. This is possible only if $\lambda = 0$ and also $k = \ell = 0$. The result follows.

It remains to derive the formula for $\text{vol}(\tilde{P}_{\alpha,\ell}, \tilde{\Omega}_\varepsilon)$. First assume that $\ell = 0$. Then the section gr_α has trivial (complex) normal bundle in $S^2 \times M$. Hence it has a neighborhood $U_\varepsilon \subset U$ that is symplectomorphic to a product $gr_\alpha \times B^{2n}(\varepsilon)$. Thus, the volume of U_ε with respect to the form Ω is $\text{vol}(U_\varepsilon) = v_\varepsilon(\mu_0 + \lambda)$ where $\mu_0 := \int_{S^2} \beta$. (Recall that β is an arbitrary area form on the base S^2 .) Since we construct the blow up by cutting out U_ε from (P, Ω) and identifying the boundary via the Hopf map, we have

$$\text{vol}(\tilde{P}_{\alpha,0}, \tilde{\Omega}_\varepsilon) = \mu_0 V - v_\varepsilon(\mu_0 + \lambda),$$

as claimed.

Now consider the case when $\ell \neq 0$. Then the normal bundle to gr_α in $M \times S^2$ is isomorphic to the product $\mathbb{C}^{n-1} \oplus L_\ell$, where $L_\ell \rightarrow S^2$ is the holomorphic line bundle with $c_1 = \ell$. Therefore, we can choose Ω so that it restricts in some neighborhood of gr_α to the product of a ball in \mathbb{C}^{n-1} with a δ -neighborhood $\mathcal{N}_\delta(L_\ell)$ of the zero section of L_ℓ . Identifying $\mathcal{N}_\delta(L_\ell)$ with part of the 4-dimensional symplectic toric manifold $\mathbb{P}(L_\ell \oplus \mathbb{C})$, we can see that its volume is $h\delta - \ell\delta^2/2$, where $h = \text{area of zero section}$ and $\delta = \pi r^2$ is the capacity (or area) of the disc normal to gr_α . (Recall that this volume is just the area of a small neighborhood of the appropriate edge of the moment polygon.) Therefore, since $h = \mu_0 + \lambda$ here,

$$\begin{aligned} \text{vol}(U_\varepsilon) &= \int_0^{\sqrt{\frac{\varepsilon}{\pi}}} \text{vol}(S^{2n-3}(r)) \cdot \text{vol}(\mathcal{N}_{\varepsilon-\pi r^2}(L_\ell)) \, dr \\ &= \frac{(\mu_0 + \lambda)\varepsilon^n}{n!} - \frac{\ell\varepsilon^{n+1}}{(n+1)!} \\ &= v_\varepsilon\left(\mu_0 + \lambda - \frac{\ell}{n+1}\varepsilon\right). \end{aligned}$$

Everything in the previous calculation remains valid except that we have to add $\frac{\ell}{n+1}\varepsilon v_\varepsilon$ to the volume of $\tilde{P}_{\alpha,\ell}$. This completes the proof. \square

3. QUANTUM METHODS

By quantum methods we really mean the use of J -holomorphic curves and quantum homology. The main tool from this theory that has been used to understand $\pi_1(\text{Ham})$ is the Seidel representation. Readers may consult the survey articles [32, §2.4] or [29] for a brief introduction to J -holomorphic curves, or the book [37] for a more detailed presentation.

⁶We cannot assume that it is independent of ε since the diffeomorphism may not converge to a product as $\varepsilon \rightarrow 0$. For example, consider the (trivial) bundle obtained by blowing up $T^2 \times T^2$ along the diagonal.

3.1. The Seidel representation. Consider the small quantum homology $QH_*(M) := H_*(M) \otimes \Lambda$ of M . Here $\Lambda := \Lambda^{univ}[q, q^{-1}]$ where q is a polynomial variable of degree 2 and Λ^{univ} denotes the generalized Laurent series ring with elements $\sum_{i \geq 1} r_i t^{\kappa_i}$, where $r_i \in \mathbb{Q}$ and $\kappa_i \in \mathbb{R}$ is a strictly decreasing sequence that tends to $-\infty$. We write the elements of $QH_*(M)$ as infinite sums $\sum_{i \geq 1} a_i \otimes q^{d_i} t^{\kappa_i}$, where $a_i \in H_*(M; \mathbb{Q}) =: H_*(M)$, $|d_i|$ is bounded and κ_i is as before. The term $a \otimes q^{d} t^{\kappa}$ has degree $2d + \deg a$.

The quantum product $a * b$ of the elements $a, b \in H_*(M) \subset QH_*(M)$ is defined as follows. Let $\xi_i, i \in I$, be a basis for $H_*(M)$ and write $\xi_i^*, i \in I$, for the basis of $H_*(M)$ that is dual with respect to the intersection pairing, that is $\xi_j^* \cdot \xi_i = \delta_{ij}$. Then

$$(3.1) \quad a * b := \sum_{i, \beta \in H_2(M; \mathbb{Z})} \langle a, b, \xi_i \rangle_{\beta}^M \xi_i^* \otimes q^{-c_1(\beta)} t^{-\omega(\beta)},$$

where $\langle a, b, \xi_i \rangle_{\beta}^M$ denotes the Gromov–Witten invariant in M that counts curves in class β through the homological constraints a, b, ξ_i . Note that if $(a * b)_{\beta} := \sum_i \langle a, b, \xi_i \rangle_{\beta}^M \xi_i^*$, then $(a * b)_{\beta} \cdot c = \langle a, b, c \rangle_{\beta}^M$. Further, $\deg(a * b) = \deg a + \deg b - 2n$, and the identity element is $\mathbb{1} := [M]$. The product is extended to $H_*(M) \otimes \Lambda$ by linearity over Λ .

The *Seidel representation* is a homomorphism \mathcal{S} from $\pi_1(\text{Ham}(M, \omega))$ to the degree $2n$ multiplicative units $QH_{2n}(M)^{\times}$ of the small quantum homology ring first considered by Seidel in [55]. One way of thinking of it (which corresponds to the formula given below) is to say that it "counts" all the sections of the bundle $P_{\gamma} \rightarrow S^2$ associated to γ . However, it can also be considered as the Floer continuation map around the loop γ in the Hamiltonian group, which is Seidel's original point of view. This second point of view makes it more clear why it is a homomorphism; the connection between them is discussed in Lalonde–McDuff–Polterovich [24].

To define \mathcal{S} , observe that each loop $\gamma = \{\phi_t\}$ in Ham gives rise to an M -bundle $P_{\gamma} \rightarrow S^2$ defined⁷ by the clutching function γ :

$$P_{\gamma} := (D_+ \times M) \cup (D_- \times M) / \sim, \quad (e^{2\pi i t}, \phi_t(x))_+ \sim (e^{2\pi i t}, x)_-$$

where D_{\pm} are copies of the unit disc in \mathbb{C} . Because the loop γ is Hamiltonian, the fiberwise symplectic form ω extends to a closed form Ω on P_{γ} , that we can arrange to be symplectic by adding to it the pullback of a suitable form on the base S^2 .

In the case of a circle action with normalized moment map $K : M \rightarrow \mathbb{R}$ we may simply take (P_{γ}, Ω) to be the quotient $(S^3 \times_{S^1} M, \Omega_c)$, where S^1 acts diagonally on S^3 and on its product with M and Ω_c pulls back to $\omega + d((c - K)\alpha)$. Here α is the standard contact form on S^3 normalized so that it descends to an area form on S^2 with total area 1, and c is any constant larger than the maximum K_{\max} of K . (This last condition implies that Ω_c is nondegenerate.) Points p_{\max}, p_{\min} in the fixed point sets

⁷Different papers have different sign conventions. Here we use those of [40, §2.1]. Thus we define the Hamiltonian vector field X_H of H_t by the equation $\omega(X_t, \cdot) = -dH_t$. If $K : S^2 \rightarrow \mathbb{R}$ is the height function and we take the form $dx_3 \wedge d\theta$ on S^2 then the induced S^1 action is generated by $\partial/\partial\theta$. Moreover the signs have been chosen so that K has positive weights at its minimum and negative weights at its maximum.

F_{\max} and F_{\min} give rise to sections $s_{\max} := S^2 \times \{p_{\max}\}$ and $s_{\min} := S^2 \times \{p_{\min}\}$. Note that our orientation conventions are chosen so that the integral of Ω over the section s_{\min} is *larger* than that over s_{\max} . For example, if $M = S^2$ and γ is a full rotation, P_γ can be identified with the one point blow up of $\mathbb{C}P^2$, and s_{\max} is the exceptional divisor. In the following we denote particular sections as s_{\max} or s_{\min} , while writing $\sigma_{\max}, \sigma_{\min}$ for the homology classes they represent.

The bundle $P_\gamma \rightarrow S^2$ carries two canonical cohomology classes, the first Chern class c_1^{Vert} of the vertical tangent bundle and the coupling class u_γ , the unique class that extends the fiberwise symplectic class $[\omega]$ and is such that $u_\gamma^{n+1} = 0$. Then we define

$$(3.2) \quad \mathcal{S}(\gamma) := \sum_{\sigma, i} \langle \xi_i \rangle_\sigma^P \xi_i^* \otimes q^{-c_1^{\text{Vert}}(\sigma)} t^{-u_\gamma(\sigma)} \in H_*(M) \otimes \Lambda,$$

Cf. [38, Def. 11.4.1] and [30]. Here $\langle \xi_i \rangle_\sigma^P$ denotes the Gromov–Witten invariant that counts curves in class σ though the single constraint ξ_i that we imagine represented in a fiber of P ; similarly, below $\langle b, \xi_i \rangle_\sigma^P$ counts curves through two fiberwise constraints. The sum is taken over all section classes in P and over the basis ξ_i for $H_*(M)$.

Further for all $b, c \in H_*(M)$

$$(3.3) \quad \mathcal{S}(\gamma) * b = \sum_{\sigma, i} \langle b, \xi_i \rangle_\sigma^P \xi_i^* \otimes q^{-c_1^{\text{Vert}}(\sigma)} t^{-u_\gamma(\sigma)}.$$

In general it is very hard to calculate $\mathcal{S}(\gamma)$. The following result is essentially due to Seidel [55]; see also McDuff–Tolman [40]. We denote by F_{\max} (resp. F_{\min}) the fixed point component on which the normalized Hamiltonian $K : M \rightarrow \mathbb{R}$ that generates the action takes its maximum (resp. minimum), and by K_{\max} (resp. K_{\min}) the maximum (resp. minimum) value of K .

Proposition 3.1. *Suppose that the Hamiltonian circle action γ is generated by the normalized Hamiltonian K . Then*

$$\mathcal{S}(\gamma) := a_{\max} \otimes q^{m_{\max}} t^{K_{\max}} + \sum_{\beta \in H_2(M; \mathbb{Z}), \omega(\beta) > 0} a_\beta \otimes q^{m_{\max} - c_1(\beta)} t^{K_{\max} - \omega(\beta)},$$

where $m_{\max} := -c_1^{\text{Vert}}(\sigma_{\max})$ and $a_{\max} \in H_*(F_{\max}) \subset H_*(M)$. Moreover, if the action is locally semifree near F_{\max} then $a_{\max} = [F_{\max}]$.

The first statement below is an immediate consequence of the above result, but the second requires some work.

Corollary 3.2 ([40]). *Let γ be a Hamiltonian circle action on M .*

(i) *If the action is semifree near either F_{\max} or F_{\min} then γ represents a nonzero element in $\pi_1(\text{Ham})$.*

(ii) *Suppose that γ is contractible in Ham , and choose k so that for all $p \in M$ the order of the stabilizer subgroup at p is at most k . Then, after replacing γ by γ^{-1} if necessary,*

$$K_{\max} \leq |K_{\min}| \leq (k-1)K_{\max}.$$

In particular, if no stabilizer subgroup has order > 2 then $K_{\max} = -K_{\min}$.

Before leaving this topic, we remark that the Seidel representation can sometimes be used to give information about $QH_*(M)$ itself. Tolman–Weitsman [57] show that if M supports a semifree S^1 action with isolated fixed points then there is a natural isomorphism between the equivariant cohomology of M and that of a product of 2-spheres. In [11, 12] Gonzalez used the Seidel representation to improve this, showing that:

Under these hypotheses M has the same quantum homology as a product of 2-spheres. Moreover, if $\dim M \leq 6$ then M is equivariantly symplectomorphic to a product of 2-spheres.

3.2. Applications to Hofer geometry. We give a very brief introduction to Hofer geometry. For more details see [15, 51, 38, 31]. The *Hofer length* $\mathcal{L}(\{\phi_t^H\})$ of a Hamiltonian path ϕ_t^H from $\phi_0^H = id$ to $\phi_1^H := \phi$ is defined to be

$$\mathcal{L}(\{\phi_t\}) = \int_0^1 \left(\max_{x \in M} H_t(x) - \min_{x \in M} H_t(x) \right) dt.$$

The Hofer (pseudo)norm of an element $\tilde{\phi} = (\phi, \{\phi_t\})$ in the universal cover $\widetilde{\text{Ham}}$ of Ham is then defined to be the infimum of the lengths of the paths to ϕ that are homotopic to $\{\phi_t\}$.

Question 3.3. *Is the Hofer pseudonorm a norm?*

It is known that this pseudonorm descends to a norm $\|\cdot\|$ on Ham , i.e. $\|\phi\| = 0 \Leftrightarrow \phi = id$. (It is obvious that $\|id\| = 0$ but the other implication is quite hard, see [15, 21].) Therefore this question is equivalent to the following:

Question 3.4. *Are there any nontrivial elements $\gamma \in \pi_1(\text{Ham})$ whose Hofer length $\|\gamma\|$ is zero?*

Since the Hofer norm is conjugation invariant, the corresponding metric given by

$$\rho(\phi, \psi) := \|\phi\psi^{-1}\|$$

is biinvariant. There are many interesting open questions in Hofer geometry. Here we focus on the question of whether the diameter of Ham can ever be finite.

One way to estimate the length of a Hamiltonian path is to use the Schwarz–Oh *spectral invariants*; see [54, 43, 58] and [38, Ch 12.4]. For each element $\tilde{\phi} \in \widetilde{\text{Ham}}$ and each element $a \in QH_*(M)$ one gets a number $c(a, \tilde{\phi}) \in \mathbb{R}$ with the following properties:

$$(3.4) \quad -\|\tilde{\phi}\| \leq c(a, \tilde{\phi}) \leq \|\tilde{\phi}\|;$$

$$(3.5) \quad c(\tilde{\phi}, \lambda a) = c(a, \tilde{\phi}) + \nu(\lambda) \quad \text{for all } \lambda \in \Lambda,$$

$$(3.6) \quad c(a, \tilde{\phi} \circ \gamma) = c(\mathcal{S}(\gamma) * a, \tilde{\phi}) \quad \text{for all } \gamma \in \pi_1(\text{Ham}),$$

$$(3.7) \quad c(a * b, \tilde{\phi}\tilde{\psi}) \leq c(a, \tilde{\phi}) + c(b, \tilde{\psi}) \quad \text{for all } \tilde{\phi}, \tilde{\psi} \in \widetilde{\text{Ham}}, a, b \in QH_*(M).$$

In (3.5), $\nu : QH_*(M) \rightarrow \mathbb{R}$ is the valuation given by

$$\nu(b) := \max\{-\kappa_i | b_i \neq 0\}, \quad \text{where } b = \sum b_i q^{d_i} t^{-\kappa_i}, b_i \in H_*(M).$$

These invariants are defined by looking at the filtered Floer complex⁸ $CF_*(H, J)$ of the generating (normalized) Hamiltonian H , and by [44, 58] turn out to be particular critical values of the corresponding action functional \mathcal{A}_H .⁹ Thus each invariant $c(a, \tilde{\phi})$ corresponds to a particular fixed point of the endpoint $\phi_1 \in \text{Ham}$ of $\tilde{\phi}$. Property (3.6) above explains how they depend on the path $\tilde{\phi}$.

It is usually very hard to calculate these numbers. However, if $\tilde{\phi}^H$ is generated by a C^2 -small mean normalized Morse function $H : M \rightarrow \mathbb{R}$, then the invariants $c(a, \tilde{\phi}^H)$ for $a \in H_*(M)$ are the same as the corresponding invariants $c_M(a, \tilde{\phi}^H)$ obtained from the Morse complex $CM_*(H)$ of H . These are defined as follows. For each $\kappa \in \mathbb{R}$, denote by $CM_*^\kappa(M, H)$ the subcomplex of the Morse complex generated by the critical points p with critical values $H(p) \leq \kappa$. Denote by ι_κ the inclusion of the homology H_*^κ of this subcomplex into $H_*^\infty \cong H_*(M)$. Then for each $a \in H_*(M)$

$$(3.8) \quad c_M(a, \tilde{\phi}^H) := \inf\{\kappa : a \in \text{Im } \iota_\kappa\}.$$

Moreover, in this case the filtered Floer complex $CF_*(H, J)$ is simply the tensor product $CM_*(H) \otimes \Lambda$ with the obvious product filtration. Hence

$$(3.9) \quad c\left(\sum a_i \otimes q^{d_i} t^{\kappa_i}, \tilde{\phi}^H\right) = \sup_i (c_M(a_i, \tilde{\phi}^H) + \kappa_i)$$

As pointed out by Ostrover [46], one can use the continuity properties of the $c(a, \tilde{\phi})$ to prove that $\widetilde{\text{Ham}}$ has infinite diameter. Indeed, let H be a small Morse function as above, choose an open set U that is displaced by ϕ_1^H (i.e. $\phi_1^H(U) \cap U = \emptyset$) and let $F : M \rightarrow U$ be a function with support in U and with nonzero integral $I := \int_M F \omega^n$. Denote the flow of F by f_t and consider the path $\tilde{\phi}_s := \{f_{ts} \phi_t^H\}_{t \in [0,1]}$ for $s \rightarrow \infty$. This is generated by the Hamiltonian

$$F_s \# H := F_s + H \circ f_{st}.$$

The corresponding normalized Hamiltonian is

$$K_s := F_s + H \circ f_{st} - sI.$$

By construction, $f_s \phi_1^H$ has the same fixed points as ϕ_1^H , namely the critical points of H . Hence the continuity and spectrality properties of $c(a, \tilde{\phi}_s)$ imply that for each $a \in QH_*(M)$ the fixed point p_a whose critical value is $c(a, \tilde{\phi}_s)$ remains unchanged as s increases. But the spectral value does change. In fact, if $a \in H_*(M)$, then when $s = 0$ there is a critical point p_a of H such that $c(a, \tilde{\phi}_0) = c_M(a, \phi_1^H) = H(p_a)$. Hence

$$(3.10) \quad c(a, \tilde{\phi}_s) = K_s(p_a) = H(p_a) - sI, \quad \text{for all } s \in \mathbb{R}, a \in H_*(M).$$

By (3.4) it follows that $\widetilde{\text{Ham}}$ has infinite Hofer diameter.

⁸The differential in this complex depends on the choice of a suitable family of almost complex structures J_t , but the spectral values are independent both of this choice and of the choice of H .

⁹This is essentially the same as the action functional used in §2.1. However, because we want it to take values in \mathbb{R} we must now define it on the universal cover $\widetilde{\mathcal{L}_0 M}$ of the loop space $\mathcal{L}_0 M$.

Schwarz proved the following result in [54, §4.3] under the assumption that both ω and c_1 vanish on $\pi_2(M)$. However, his argument (which is based on an idea due to Seidel) works equally well if one just assumes that ω vanishes.

Proposition 3.5. *If ω vanishes on $\pi_2(M)$ then every Seidel element has the form*

$$\mathcal{S}(\gamma) = (\mathbb{1} + x) \otimes \lambda, \quad \text{where } \nu(\lambda) = 0, x \in H_{<2n}(M) \otimes \Lambda.$$

Moreover the spectral invariants descend to Ham. In particular, Ham has infinite diameter.

Sketch of proof. If $\omega = 0$ on $\pi_2(M)$, quantum multiplication is the same as the intersection product. Hence all units in $QH_*(M)$ have the form $\mathbb{1} \otimes \lambda + x$ where $\lambda \in \Lambda$ is nonzero and $x \in H_{<2n}(M) \otimes \Lambda$. In particular there must be a section class σ_0 of P_γ with $c_1^{\text{vert}}(\sigma_0) = 0$ that contributes to the coefficient of $\mathbb{1}$. Moreover, because all sections of P_γ have the same energy, the Seidel element has the form $\mathcal{S}(\gamma) = (\mathbb{1} + x) \otimes t^\kappa$ where $x \in H_{<2n}(M)[q, q^{-1}]$. Therefore it suffices to show that $\kappa := -\nu(\sigma_0) = 0$.

Now choose a generic almost complex structure J on P that is compatible with the fibration $P \rightarrow S^2$ and consider the moduli space \mathcal{M} of J holomorphic sections in class σ_0 , parametrized as sections. This space is always compact because there can be no fiberwise bubbles, and is a manifold for generic J . Hence there is a commutative diagram

$$\begin{array}{ccc} S^2 \times \mathcal{M} & \xrightarrow{ev} & P \\ \downarrow & & \downarrow \\ S^2 & \xrightarrow{=} & S^2, \end{array}$$

where the top horizontal map is the evaluation map. Since ev maps each fiber $\{z\} \times \mathcal{M}$ to the corresponding fiber $\pi^{-1}(z)$ with positive degree, the identity $(u_\gamma)^{n+1} = 0$ implies that the class $ev^*(u_\gamma) \in H^2(S^2 \times \mathcal{M})$ is pulled back from \mathcal{M} . Hence

$$\kappa := - \int_{\sigma_0} u_\gamma = \int_{S^2 \times \{pt\}} ev^*(u_\gamma) = 0.$$

This proves the first statement. It follows that $\nu(\mathcal{S}(\gamma)) \leq 0$ for all γ . But it is impossible to have $\nu(\mathcal{S}(\gamma)) < 0$ for some γ because $\mathcal{S}(\gamma) * \mathcal{S}(\gamma^{-1}) = \mathbb{1}$ and $\nu(a * b) \leq \max(\nu(a), \nu(b))$. Hence $\nu(\mathcal{S}(\gamma)) = 0$ for all γ . But then the multiplicative relation (3.7) implies that

$$c(a, \tilde{\phi} \circ \gamma) \leq c(a, \tilde{\phi}) + c(\mathbb{1}, \gamma) = c(a, \tilde{\phi}) + \nu(\mathcal{S}(\gamma)) = c(a, \tilde{\phi}),$$

for all $\gamma \in \pi_1(\text{Ham})$. But then $c(a, \tilde{\phi}) = c(a, \tilde{\phi} \circ \gamma \circ \gamma^{-1}) \leq c(a, \tilde{\phi} \circ \gamma)$. Hence $c(a, \tilde{\phi}) = c(a, \tilde{\phi} \circ \gamma)$ for all γ , and so the spectral invariants descend. \square

These arguments are significantly extended in [36]. For example, if $\dim H_2(M) > 1$, the condition $[\omega] = 0$ on $\pi_2(M)$ in the above proposition can be replaced by the much weaker condition that all the nontrivial genus zero 3-point Gromov–Witten invariants of M vanish.

Here is a somewhat different approach to this question that gives information about the Hofer diameter even when the spectral invariants do not descend. We shall say that $\pi_1(\text{Ham})$ is *spectrally asymmetric* if

$$\liminf_{\|\gamma\| \rightarrow \infty} \frac{|\nu(\mathcal{S}(\gamma)) + \nu(\mathcal{S}(\gamma^{-1}))|}{\|\gamma\|} \geq \varepsilon > 0.$$

Roughly speaking, this means that for loops γ that are long with respect to the Hofer norm $\|\gamma\|$, their spectral “norms” $\nu(\mathcal{S}(\gamma))$ and $\nu(\mathcal{S}(\gamma^{-1}))$ are not comparable. For example, if $\pi_1(\text{Ham}) \cong \mathbb{Z}$ with generator α , this holds if

$$\lim_{k \rightarrow \infty} \frac{\nu(\mathcal{S}(\alpha^k))}{k} \neq - \lim_{k \rightarrow \infty} \frac{\nu(\mathcal{S}(\alpha^{-k}))}{k}$$

because $\|\alpha^k\| \leq k\|\alpha\|$.

Lemma 3.6. *Ham has infinite diameter whenever $\pi_1(\text{Ham})$ is spectrally asymmetric.*

Sketch of proof. Consider the path $\tilde{\phi}_s$ of equation (3.10). Then by equation (3.9) there is for any $\gamma \in \pi_1(\text{Ham})$ a critical point p_γ of H such that

$$\begin{aligned} c(\mathbb{1}, \gamma \tilde{\phi}_s) &= c(\mathcal{S}(\gamma), \tilde{\phi}_s) \\ &= c(\mathcal{S}(\gamma), f_s \phi_1^H) \\ &= c(\mathcal{S}(\gamma), \phi_1^H) - sI \\ &= H(p_\gamma) + \nu(\mathcal{S}(\gamma)) - sI. \end{aligned}$$

Now suppose that Ham has diameter D . Because the Hofer norm is symmetric, i.e. $\|\phi\| = \|\phi^{-1}\|$, for each s there must be a loop γ_s such that

$$|c(\mathbb{1}, \gamma_s \tilde{\phi}_s)| \leq D, \quad |c(\mathbb{1}, (\gamma_s \tilde{\phi}_s)^{-1})| \leq D.$$

Therefore,

$$|c(\mathcal{S}(\gamma_s), \phi_1^H) - sI| \leq D, \quad |c(\mathcal{S}(\gamma_s^{-1}), \phi_1^{-H}) + sI| \leq D.$$

But by (3.4) $|\nu(\mathcal{S}(\gamma_s))| = |c(\mathbb{1}, \gamma_s)| \leq \|\gamma_s\|$. Hence $\|\gamma_s\| \rightarrow \infty$ as $s \rightarrow \infty$, and so

$$(3.11) \quad \lim_{s \rightarrow \infty} \frac{|c(\mathcal{S}(\gamma_s), \phi_1^H) + c(\mathcal{S}(\gamma_s^{-1}), \phi_1^{-H})|}{\|\gamma_s\|} = 0.$$

But we saw above that there are critical points p, q of H such that

$$\begin{aligned} c(\mathcal{S}(\gamma_s), \phi_1^H) &= H(p) + \nu(\mathcal{S}(\gamma_s)), \\ c(\mathcal{S}(\gamma_s^{-1}), \phi_1^{-H}) &= -H(q) + \nu(\mathcal{S}(\gamma_s^{-1})). \end{aligned}$$

Hence the LHS of (3.11) has the same limit as

$$\frac{|\nu(\mathcal{S}(\gamma_s)) + \nu(\mathcal{S}(\gamma_s^{-1}))|}{\|\gamma_s\|},$$

and this is $\neq 0$ by assumption. This contradiction shows that Ham cannot have finite diameter. \square

Now consider the one point blow up M of $\mathbb{C}P^2$ with symplectic form ω_a that takes the value πa on the exceptional divisor and π on the line. (M, ω_a) is a “small” blow up if $3a^2 < 1$, i.e. the exceptional divisor is smaller than it is in the monotone case.

Corollary 3.7. $\text{Ham}(M, \omega_a)$ has infinite diameter when $3a^2 < 1$.

Proof. By [2] $\pi_1(\text{Ham}(M, \omega_a)) = \mathbb{Z}$. By [31, Prop. 5.3], it has a generator α with $\mathcal{S}(\alpha) = Q \otimes t^\mu$ where $\mu = \pi \frac{(1-a^2)^2}{12(1+a^2)}$. Moreover when $3a^2 < 1$ there are positive constants c_1, c_2 such that

$$\nu(Q^{-k}) \geq c_1 k, \quad \nu(Q^k) \geq c_2 k, \quad k \geq 1.$$

Hence, because $\|\alpha^k\| \leq k\|\alpha\|$,

$$\liminf_{k \rightarrow \infty} \frac{|\nu(\mathcal{S}(\alpha^k)) + \nu(\mathcal{S}(\alpha^{-k}))|}{\|\alpha^k\|} \geq \liminf_{k \rightarrow \infty} \frac{|(c_1 - \mu)k + (c_2 + \mu)k|}{k\|\alpha\|} > 0.$$

Thus $\text{Ham}(M, \omega_a)$ is spectrally asymmetric. \square

It is shown in [31] that $|\nu(Q^k)|$ is bounded when $3a^2 \geq 1$. Therefore in this case our methods give no information on the Hofer diameter.

Remark 3.8. In the above Lemma 3.6 we exploited the symmetric definition of the Hofer norm. As pointed out in [31] there are variants of this norm with more asymmetric definitions; for example one can measure the size of $\phi \in \text{Ham}$ by separately minimizing the positive and negative parts of the Hofer lengths over all paths in Ham from id to ϕ . This seminorm is not known to be nondegenerate for all M , though it is for the one point blow up of $\mathbb{C}P^2$. Moreover the argument in Lemma 3.6 would not apply to it.

3.3. The Calabi homomorphism. Finally, we show how to calculate the Calabi homomorphism $\pi_1(\text{Ham}^c) \rightarrow \mathbb{R}$ for certain loops on noncompact manifolds M , and also explain its relation to the Seidel element.

Suppose that the loop $\gamma \in \pi_1(\text{Ham}^c M)$ is supported in an open subset $U \subset M$ such that (U, ω) can be symplectically embedded into some closed manifold (X, ω_X) . Consider the corresponding bundle $X \rightarrow P_{\gamma_X} \rightarrow S^2$ and the corresponding Seidel element $\mathcal{S}(\gamma_X) \in QH_*(X)$. Then P_{γ_X} has a family of flat sections $S^2 \times \{p\}, p \in X \setminus U$. Suppose now that these sections contribute nontrivially to the coefficient of $\mathbb{1} \otimes t^\kappa$ in $\mathcal{S}(\gamma_X)$.

Lemma 3.9. $\text{Cal}(\gamma) = -\kappa \text{Vol}(X, \omega_X)$.

Sketch of Proof. Let $\gamma_X = \{\phi_t\}$ and suppose that $H_t : X \rightarrow \mathbb{R}$ is the corresponding mean normalized Hamiltonian. Then each H_t equals some constant c_t outside U so

that

$$\begin{aligned}
\text{Cal}(\gamma) &= \frac{1}{n!} \int_0^1 \left(\int_U (H_t - c_t) \omega^n \right) dt \\
&= \frac{1}{n!} \int_0^1 \left(\int_X (H_t - c_t) \omega_X^n \right) dt \\
&= -\text{Vol}(X, \omega_X) \int_0^1 c_t dt.
\end{aligned}$$

On the other hand, an easy calculation shows that the coupling class u_{γ_X} is represented by a form that we can take to be $pr^*(\omega)$ on $D_- \times X$ and $pr^*(\omega) - d(\beta(r)H_t dt)$ on $D_+ \times X$ for some cut off function β . Then u_{γ_X} restricts to zero on $D_- \times \{p\}$ and so by the definition of $\mathcal{S}(\gamma)$:

$$\begin{aligned}
-\kappa &:= \int_{S^2 \times \{p\}} u_\gamma = - \int_{D_+ \times \{p\}} d(\beta(r)H_t) dt \\
&= - \int_{D_+} \beta'(r) dr H_t(x_0) dt \\
&= - \int_{S^1} H_t(x_0) dt = - \int_0^1 c_t dt.
\end{aligned}$$

The result follows. \square

Remark 3.10. (i) In constructing explicit examples where $\text{Cal}(\gamma)$ does not vanish, one can proceed the other way around. Suppose given a loop γ_X for which the corresponding bundle $P_{\gamma_X} \rightarrow S^2$ has a section s_X with trivial normal bundle and with $u_{\gamma_X}(s_X) = -\kappa \neq 0$. (One way to find such a section is to show that the term $\mathbb{1} \otimes t^k$ appears with nontrivial coefficient in $\mathcal{S}(\gamma_X)$.) Then we may trivialize the bundle $P_{\gamma_X} \rightarrow S^2$ near s_X and take γ to be the loop in $\text{Ham}^c(X \setminus \{pt\})$ corresponding to the fibration $(P_{\gamma_X} \setminus s_X) \rightarrow S^2$. If $\mathcal{S}(\gamma) = \mathbb{1} \otimes \lambda + x$ where $\lambda = \sum r_i t^{\varepsilon_i}$ has infinitely many nonzero terms then one would get infinitely many different loops in $\text{Ham}^c(X \setminus \{pt\})$. However, there would be algebraic relations between them stemming from the long exact homotopy sequence of the fibration

$$\text{Ham}^{id}(X, p) \longrightarrow \text{Ham } X \xrightarrow{ev} \text{Fr}(X),$$

where $\text{Fr}(X)$ is the symplectic frame bundle of X and $\text{Ham}^{id}(X, p)$ denotes the subgroup of $\text{Ham } X$ consisting of elements ϕ that fix the point p and have $d\phi_p = id$. (A standard Moser type argument shows that the inclusion $\text{Ham}^c(X \setminus \{p\}) \rightarrow \text{Ham}^{id}(X, p)$ is a homotopy equivalence.)

(ii) The manifolds $M = X \setminus \{pt\}$ obtained in (i) have finite volume and are concave at infinity. However, if one started with a bundle $\widetilde{M} \rightarrow \widetilde{P} \rightarrow S^2$ formed by blowing up the product bundle along the section gr_α then this bundle is trivial over $\Phi^{-1}(S^2 \times (M \setminus V))$, where $\Phi: \widetilde{P} \rightarrow P = S^2 \times M$ is the blowdown map and V is a neighborhood of $\text{im } \alpha$ in M . Hence we could get examples in $\pi_1(\text{Ham}^c W)$ where W is any symplectic extension

of (V, ω) . Note also that one can of course calculate $\text{Cal}(\gamma)$ for these blow up bundles without any reference to the Seidel element. All one needs is the formula for the coupling class given in the proof of Lemma 3.6.

Remark 3.11 (Questions about non-compact manifolds). The above discussion of the diameter of Ham assumes that (M, ω) is closed. If M is not closed then one can ask a similar question about Ham^c . One really should divide into two cases here, depending on whether M has finite or infinite volume.¹⁰

If $\text{Vol } M$ is finite, then the universal cover $\widetilde{\text{Ham}}^c$ of Ham^c has infinite Hofer diameter. To see this, observe first that the Calabi homomorphism defined earlier on $\pi_1(\text{Ham}^c)$ extends to $\widetilde{\text{Ham}}^c$ by:

$$\text{Cal} : \widetilde{\text{Ham}}^c \rightarrow \mathbb{R}, \quad \tilde{\phi} \mapsto \int_0^1 \int_M H_t \omega^n dt,$$

where H_t is any compactly supported functions that generates $\tilde{\phi}$. Next note the obvious fact that $\text{Cal}(\tilde{\phi}) \leq \text{Vol}(M) \|\tilde{\phi}\|$.

However I know of no methods to help one decide whether the *kernel* of Cal has infinite diameter: Ostrover's argument on spectral drift does not work here because all Hamiltonians are normalized to vanish at infinity. Since this question is unsolved there is not much point in wondering about the corresponding subgroup of Ham^c . (Banyaga showed that this is a perfect group; hence in many ways this kernel plays the same role for noncompact M as does Ham in the closed case.)

Nevertheless there are interesting open questions about the homomorphism $\text{Cal} : \pi_1(\text{Ham}^c) \rightarrow \mathbb{R}$ itself. It is known to vanish when ω is exact; cf. [37, Ch 10]. On the other hand we saw above that it need not vanish in general. One obvious question is whether it always has discrete image.¹¹ By Lemma 3.9 this is closely related to the question of the possible values of $\nu(\mathcal{S}(\gamma))$ for $\gamma \in \pi_1(\text{Ham } X)$ where X is a closed manifold containing M .¹² (But note that the function $\gamma \rightarrow \nu(\mathcal{S}(\gamma))$ need not be a homomorphism.) Of course if Cal does not vanish on $\pi_1(\text{Ham}^c)$ then one can also ask if Ham^c itself has infinite diameter.

¹⁰Observe that it is known by a rather intricate geometric argument that the Hofer seminorm is a norm on Ham^c for *every* open manifold M , whether nice at infinity or not; see [21].

¹¹Here one should no doubt restrict to M of finite type, eg those that are the interior of a compact manifold \overline{M} with boundary. But even this topological restriction on M may not be enough. Because symplectic forms exist on open manifolds provided the most mild topological conditions are satisfied, one may need to ask that ω extend to \overline{M} in order to have a hope of discreteness here. For example, a similar question concerning the flux homomorphism was explored in [27], and examples of open manifolds M were constructed for which $\text{Flux}(\pi_1(\text{Symp}^c M))$ is not discrete. However no control at infinity was imposed. Presumably Ono's argument in [45] would extend if the control at infinity were strong enough.

¹²One could try to define some kind of analog to the Seidel representation, though one would have to specify exactly which quantum homology group is the target of this homomorphism. However, it is not at all clear how such a representation could help understand the questions at hand. Lemma 3.9 suggests to me that in the noncompact case the role of the map $\gamma \mapsto \nu(\mathcal{S}(\gamma))$ is played by the Calabi homomorphism.

Finally we remark that in the case when M has infinite volume, there is no direct connection between the Hofer norm and the Calabi homomorphism. Therefore, one should tackle the question of the Hofer diameter by other means, for example by seeing if some version of the energy–capacity inequality applies.

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