

# Einstein Metrics and Mostow Rigidity

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## Abstract

Using the new diffeomorphism invariants of Seiberg and Witten, a uniqueness theorem is proved for Einstein metrics on compact quotients of irreducible 4-dimensional symmetric spaces of non-compact type. The proof also yields a Riemannian version of the Miyaoka-Yau inequality.

A smooth Riemannian manifold  $(M, g)$  is said [1] to be *Einstein* if its Ricci curvature is a constant multiple of  $g$ . Any irreducible locally-symmetric space is Einstein, and, in light of Mostow rigidity [5], it is natural to ask whether, up to diffeomorphisms and rescalings, the standard metric is the only Einstein metric on any compact quotient of an irreducible symmetric space of non-compact type and dimension  $> 2$ . For example, any Einstein 3-manifold has constant curvature, so the answer is certainly affirmative in dimension 3. In dimension  $\geq 4$ , however, solutions to Einstein's equations can be quite non-trivial. Nonetheless, the following 4-dimensional result was recently proved by means of an entropy comparison theorem [2]:

**Theorem 1 (Besson-Courtois-Gallot)** *Let  $M^4$  be a smooth compact quotient of hyperbolic 4-space  $\mathcal{H}^4 = SO(4,1)/SO(4)$ , and let  $g_0$  be its standard metric of constant sectional curvature. Then every Einstein metric  $g$  on  $M$  is of the form  $g = \lambda\varphi^*g_0$ , where  $\varphi : M \rightarrow M$  is a diffeomorphism and  $\lambda > 0$  is a constant.*

In this note, we will prove the analogous result for the remaining 4-dimensional cases:

**Theorem 2** *Let  $M^4$  be a smooth compact quotient of complex-hyperbolic 2-space  $\mathcal{CH}_2 = SU(2,1)/U(2)$ , and let  $g_0$  be its standard complex-hyperbolic metric. Then every Einstein metric  $g$  on  $M$  is of the form  $g = \lambda\varphi^*g_0$ , where  $\varphi : M \rightarrow M$  is a diffeomorphism and  $\lambda > 0$  is a constant.*

In contrast to Theorem 1, the proof of this result is based on the new 4-manifold invariants [4] recently introduced by Witten [6].

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# 1 Seiberg-Witten Invariants

While the results in this section are largely due to Edward Witten [6], the crucial sharp form of the scalar-curvature inequality was pointed out to the author by Peter Kronheimer.

Let  $(M, g)$  be a smooth compact Riemannian manifold, and suppose that  $M$  admits an almost-complex structure. Then the given component of the almost-complex structures on  $M$  contains almost-complex structures  $J : TM \rightarrow TM$ ,  $J^2 = -1$  which are compatible with  $g$  in the sense that  $J^*g = g$ . Fixing such a  $J$ , the tangent bundle  $TM$  of  $M$  may be given the structure of a rank-2 complex vector bundle  $T^{1,0}$  by defining scalar multiplication by  $i$  to be  $J$ . Setting  $\wedge^{0,p} := \overline{\wedge^p T^{1,0}}^* \cong \wedge^p T^{1,0}$ , we may then then define rank-2 complex vector bundles  $V_{\pm}$  on  $M$  by

$$V_+ = \wedge^{0,0} \oplus \wedge^{0,2} \tag{1}$$

$$V_- = \wedge^{0,1}, \tag{2}$$

and  $g$  induces canonical Hermitian inner products on these bundles.

As described, these bundles depend on the choice of a particular almost-complex structure  $J$ , but they have a deeper meaning [3] that depends only on the homotopy class  $c$  of  $J$ ; namely, if we restrict to a contractible open set  $U \subset M$ , the bundles  $V_{\pm}$  may be canonically identified with  $\mathbf{S}_{\pm} \otimes L^{1/2}$ , where  $\mathbf{S}_{\pm}$  are the left- and right-handed spinor bundles of  $g$ , and  $L^{1/2}$  is a complex line bundle whose square is the ‘anti-canonical’ line-bundle  $L = \overline{\wedge^{0,2}} \cong (\wedge^{0,2})^*$ . For each connection  $A$  on  $L$  compatible with the  $g$ -induced inner product, we can thus define a corresponding Dirac operator

$$D_A : C^\infty(V_+) \rightarrow C^\infty(V_-).$$

If  $J$  is parallel with respect to  $g$ , so that  $(M, g, J)$  is a Kähler manifold, and if  $A$  is the Chern connection on the anti-canonical bundle  $L$ , then  $D_A = \sqrt{2}(\bar{\partial} \oplus \bar{\partial}^*)$ , where  $\bar{\partial} : C^\infty(\wedge^{0,0}) \rightarrow C^\infty(\wedge^{0,1})$  is the  $J$ -antilinear part of the exterior differential  $d$ , acting on complex-valued functions, and where  $\bar{\partial}^* : C^\infty(\wedge^{0,2}) \rightarrow C^\infty(\wedge^{0,1})$  is the formal adjoint of the map induced by the exterior differential  $d$  acting on 1-forms; more generally,  $D_A$  will differ from  $\sqrt{2}(\bar{\partial} \oplus \bar{\partial}^*)$  by only  $0^{th}$  order terms.

The Seiberg-Witten equations

$$D_A \Phi = 0 \tag{3}$$

$$F_A^+ = i\sigma(\Phi). \tag{4}$$

are equations for an unknown smooth connection  $A$  on  $L$  and an unknown smooth section  $\Phi$  of  $V_+$ . Here the purely imaginary 2-form  $F_A^+$  is the self-dual part of the curvature of  $A$ , and, in terms of (1), the real-quadratic map

$\sigma : V_+ \rightarrow \Lambda_+^2$  is given by

$$\sigma(f, \phi) = (|f|^2 - |\phi|^2) \frac{\omega}{4} + \Im m(\bar{f}\phi),$$

where  $\omega(\cdot, \cdot) = g(J\cdot, \cdot)$  is the ‘Kähler’ form. Notice that  $|F^+| = 2^{-3/2}|\Phi|^2$ .

For each solution  $(A, \Phi)$  of (3) and (4) one can generate a new solution  $(A + 2d \log f, f\Phi)$  for any  $f : M \rightarrow S^1 \subset \mathbf{C}$ ; two solutions which are related in this way are called *gauge equivalent*, and may be considered to be geometrically identical. A solution is called *reducible* if  $\Phi \equiv 0$ ; otherwise, it is called *irreducible*.

A useful generalization of the Seiberg-Witten equations is obtained by replacing (4) with the equation

$$iF^+ + \sigma(\Phi) = \varepsilon \tag{5}$$

for an arbitrary  $\varepsilon \in C^\infty(\Lambda^+)$ . We can then consider the map which sends solutions of (3) and (5) to the corresponding  $\varepsilon \in C^\infty(\Lambda^+)$ , and define a solution to be *transverse* if it is a regular point of this map — i.e. if the linearization  $C^\infty(V_+ \oplus \Lambda^1) \rightarrow C^\infty(\Lambda_+^2)$  of the left-hand-side of (5), constrained by the linearization of (3), is surjective.

**Example** Let  $(M, g, J)$  be a Kähler surface of constant scalar curvature  $s < 0$ . Let  $\Phi = (\sqrt{-s}, 0) \in \Lambda^{0,0} \oplus \Lambda^{0,2}$ , and let  $A$  be the Chern connection on the anti-canonical bundle. Since  $F_A^+ = is\omega/4$ ,  $(\Phi, A)$  is an irreducible solution of the Seiberg-Witten equations (3) and (4).

The linearization of (3) at this solution is just

$$(\bar{\partial} \oplus \bar{\partial}^*)(u + \psi) = -\frac{\sqrt{-s}}{2}\alpha, \tag{6}$$

where  $(u, \psi) \in C^\infty(V_+)$  is the linearization of  $\Phi = (f, \phi)$  and  $\alpha \in \Lambda^{0,1}$  is the  $(0, 1)$ -part of the purely imaginary 1-form which is the linearization of  $A$ . Linearizing (5) at our solution yields the operator

$$(u, \psi, \alpha) \mapsto id^+(\alpha - \bar{\alpha}) + \frac{\sqrt{-s}}{2}(\Re u)\omega + \sqrt{-s}\Im m\psi.$$

Since the right-hand-side is a real self-dual form, it is completely characterized by its component in the  $\omega$  direction and its  $(0, 2)$ -part. The  $\omega$ -component of this operator is just

$$(u, \psi, \alpha) \mapsto \Re \left[ -\bar{\partial}^* \alpha + \frac{\sqrt{-s}}{2}u \right],$$

while the  $(0, 2)$ -component is

$$(u, \psi, \alpha) \mapsto i\bar{\partial}\alpha - i\frac{\sqrt{-s}}{2}\psi.$$

Substituting (6) into these expressions, we obtain the operator

$$\begin{aligned} C^\infty(\mathbf{C} \oplus \wedge^{0,2}) &\longrightarrow C^\infty(\mathbf{R} \oplus \wedge^{0,2}) \\ (u, \psi) &\longmapsto \left( \frac{1}{\sqrt{-s}} \Re \left[ \Delta - \frac{s}{2} \right] u, -\frac{i}{\sqrt{-s}} \left[ \Delta - \frac{s}{2} \right] \psi \right), \end{aligned}$$

which is surjective because  $s < 0$  is not in the spectrum of the Laplacian. The constructed solution is therefore transverse.  $\square$

Relative to  $c = [J]$ , a metric  $g$  will be called *excellent* if it admits only irreducible transverse solutions of (3) and (4). Relative to any excellent metric, the set of solutions of (3) and (4), modulo gauge equivalence, is finite [4, 6]. Notice that a metric  $g$  is automatically excellent if the corresponding equations (3) and (4) admit no solutions at all.

**Definition 1** *Let  $(M, c)$  be a compact 4-manifold equipped with a homotopy class  $c = [J]$  of almost-complex structures. Assume either that  $b_+(M) > 1$ , or that  $b_+ = 1$  and that  $(2\chi + 3\tau)(M) > 0$ . If  $g$  is an excellent metric on  $M$ , define the (mod 2) Seiberg-Witten invariant  $n_c(M) \in \mathbf{Z}_2$  to be*

$$n_c(M) = \#\{\text{gauge classes of solutions of (3) and (4)}\} \bmod 2$$

*calculated with respect to  $g$ .*

It turns out [4] that  $n_c(M)$  is actually metric-independent; when  $b_+ = 1$ , this fact depends on the assumption that  $c_1(L)^2 = 2\chi + 3\tau > 0$ , which guarantees that (3) and (4) cannot admit reducible solutions for any metric.

**Theorem 3** *Let  $(M, J)$  be a compact complex surface, where the underlying oriented 4-manifold  $M$  is as in Definition 1. Suppose that  $(M, J)$  admits a Kähler metric  $g$  of constant scalar curvature  $s < 0$ , and let  $c = [J]$ . Then  $n_c(M) = 1 \in \mathbf{Z}_2$ .*

**Proof.** With respect to  $g$  we shall show that, up to gauge equivalence, there is exactly one solution of the Seiberg-Witten equations, namely the one described in the above example. Indeed, the Weitzenböck formula for the twisted Dirac operator and equation (4) tell us that

$$0 = D_A^* D_A \Phi = \nabla^* \nabla \Phi + \frac{s}{4} \Phi + \frac{1}{4} |\Phi|^2 \Phi,$$

which implies [4] the  $C^0$  estimate  $|\Phi|^2 \leq -s$ , with equality only at points where  $\nabla \Phi = 0$ . Since

$$|F_A^+|^2 = \frac{1}{8} |\Phi|^4 \leq \frac{s^2}{8},$$

it follows that

$$\int_M |F_A^+|^2 d\mu \leq \int_M \left(\frac{s}{4}|\omega|\right)^2 d\mu = \int_M |\rho^+|^2 d\mu$$

where the Ricci form  $\rho$  is in the same cohomology class as the closed form  $F_A$ , namely  $2\pi c_1(L) = 2\pi c_1(M, J)$ . But since  $s$  is constant,  $\rho$  is harmonic, and we must therefore have that

$$\int_M |\rho^+|^2 d\mu = 2\pi^2 c_1(L)^2 + \frac{1}{2} \int_M |\rho|^2 d\mu \leq 2\pi^2 c_1(L)^2 + \frac{1}{2} \int_M |F_A|^2 d\mu = \int_M |F_A^+|^2 d\mu$$

because a harmonic form minimizes the  $L^2$  norm among closed forms in its deRham class. Hence  $F_A = \rho$ , and  $A$  differs from the Chern connection on  $L$  by twisting with a flat connection. But also  $|\Phi|^2 \equiv -s$ , which forces  $\nabla\Phi \equiv 0$ . Since  $c_1(L) \neq 0$ , the induced connection on  $\wedge^{0,2} \subset V_+$  has non-trivial curvature, and  $\Phi$  must therefore be a section of  $\wedge^{0,0}$ . Since  $\Phi$  is parallel, the induced connection on  $\wedge^{0,0}$  must not only be flat, but also have trivial holonomy. Thus  $A$  must exactly be the Chern connection on  $L$ , and our solution coincides, up to gauge transformation, with that of the example. In particular, every solution with respect to  $g$  is irreducible and transverse, so  $g$  is excellent. But since there is only one gauge class of solutions with respect to  $g$ , we conclude that  $n_c(M) = 1 \pmod 2$ .  $\blacksquare$

The following refinement an observation of Witten [6, §3] is the real key to the proof of Theorem 2.

**Theorem 4** *Let  $M$  be a smooth compact oriented 4-manifold with  $2\chi(M) + 3\tau(M) > 0$ . Suppose that there is an orientation-compatible class  $c = [J]$  of almost-complex structures for which the Seiberg-Witten invariant  $n_c(M) \in \mathbf{Z}_2$  is non-zero. Let  $g$  be a metric of constant scalar curvature  $s$  and volume  $V$  on  $M$ . Then*

$$s\sqrt{V} \leq -2^{5/2}\pi\sqrt{2\chi + 3\tau},$$

*with equality iff  $g$  is Kähler-Einstein with respect to some integrable complex structure  $J$  in the homotopy class  $c$ .*

**Proof.** For any given metric  $g$  on  $M$ , there must exist a solution of (3) and (4), since otherwise we would have  $n_c(M) = 0$ . But since  $|F_A|^2 = |\Phi|^4/8 \leq -s/8$ , with equality iff  $\nabla\Phi = 0$ , it follows that

$$2\chi + 3\tau = c_1(L)^2 = \frac{1}{4\pi^2} \int_M (|F_A^+|^2 - |F_A^-|^2) d\mu \leq \frac{1}{32\pi^2} \int_M s^2 d\mu,$$

with equality only if  $\nabla F_A^+ \equiv 0$  and  $F_A^- = 0$ . If equality holds, the parallel self-dual form  $\sqrt{2}F_A/|F_A|$  corresponds via  $g$  to a parallel almost-complex structure

$J$ , and the manifold is thus Kähler, with Kähler class  $8\pi/s$  times  $c_1(M, J) = c_1(L)$ . But since  $s$  is constant, the Ricci form is harmonic, and the manifold is Kähler-Einstein.

On the other hand, any Kähler-Einstein metric will saturate the bound in question, since the first Chern class of a Kähler-Einstein surface is  $[s\omega/8\pi]$ , and the metric volume form is  $d\mu = \omega^2/2$ .  $\blacksquare$

## 2 The Miyaoka-Yau Inequality

For any compact oriented Riemannian 4-manifold  $(M, g)$ , the Euler characteristic and signature can be expressed as

$$\begin{aligned}\chi(M) &= \frac{1}{8\pi^2} \int_M \left( |W_+|^2 + |W_-|^2 + \frac{s^2}{24} - \frac{|\text{ric}_0|^2}{2} \right) d\mu \\ \tau(M) &= \frac{1}{12\pi^2} \int_M (|W_+|^2 - |W_-|^2) d\mu\end{aligned}$$

where  $s$ ,  $\text{ric}_0$ ,  $W_+$  and  $W_-$  are respectively the scalar, trace-free Ricci, self-dual Weyl, and anti-self-dual Weyl parts of the curvature tensor; pointwise norms are calculated with respect to  $g$ , and  $d\mu$  is the metric volume form. If  $g$  is Einstein,  $\text{ric}_0 = 0$ , and  $M$  therefore satisfies

$$(2\chi \pm 3\tau)(M) = \frac{1}{4\pi^2} \int_M \left( 2|W_{\pm}|^2 + \frac{s^2}{24} \right) d\mu,$$

so the *Hitchin-Thorpe inequality*  $2\chi \geq 3|\tau|$  holds, with equality iff  $g$  is flat.

Now assume that  $M$  admits a homotopy class of almost-complex structures for which the Seiberg-Witten invariant is non-zero. If  $g$  is an Einstein metric on  $M$ , Theorem 4 then tells us that

$$\begin{aligned}2\chi + 3\tau &\leq \frac{1}{32\pi^2} \int_M s^2 d\mu \\ &\leq 3 \left[ \frac{1}{4\pi^2} \int_M \left( |2W_-|^2 + \frac{s^2}{24} \right) d\mu \right] \\ &= 3(2\chi - 3\tau)\end{aligned}$$

with equality iff the metric is Kähler and  $W_- = 0$ . But curvature operator of any Kähler manifold is an element of  $\wedge^{1,1} \otimes \wedge^{1,1}$ , and in real dimension 4 one also has  $\wedge^{1,1} = \wedge^- \oplus \mathbf{C}\omega$ , where  $\omega$  is the Kähler form; when  $W_- : \wedge_- \rightarrow \wedge_-$  and  $\text{ric}_0 : \wedge_- \rightarrow \wedge_+$  both vanish, the curvature operator must therefore be of the form

$$\mathcal{R} = s\omega \otimes \omega + \frac{s}{12} \mathbf{1}_{\wedge_-}$$

and so satisfy

$$\nabla \mathcal{R} = 0,$$

which is to say that  $(M, g)$  must be locally symmetric. If  $s$  is negative, the point-wise form of the curvature tensor then implies that the exponential map induces an isometry between the universal cover of  $(M, g)$  and a complex-hyperbolic space which has been rescaled so as to have the same scalar curvature. This proves the following generalization of the Miyaoka-Yau inequality [7]:

**Theorem 5** *Let  $(M, g)$  be a compact Einstein 4-manifold, and suppose that  $M$  admits an almost-complex structure  $J$  for which the Seiberg-Witten invariant is non-zero. Also assume that  $M$  is not finitely covered by the 4-torus  $T^4$ . Then, with respect to the orientation of  $M$  determined by  $J$ , the Euler characteristic and signature of  $M$  satisfy*

$$\chi \geq 3\tau,$$

*with equality iff the universal cover of  $(M, g)$  is complex-hyperbolic 2-space  $\mathbf{CH}_2 := SU(2, 1)/U(2)$  with a constant multiple of its standard metric.*

On the other hand, Theorem 3 tells us the Seiberg-Witten invariant of any complex hyperbolic 4-manifold  $M = \mathbf{CH}_2/\Gamma$  is actually non-zero. Theorem 5 and Mostow rigidity thus imply Theorem 2.

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