

Chinese dragons and mating trees

Jason Miller and Scott Sheffield

Massachusetts Institute of Technology

October, 2014

Overview

Part I: Cast of Characters

1. **Fractals from complex dynamics:** background, motivation, Julia sets, matings
2. **Canonical random trees:** Brownian motion, continuum random tree
3. **Canonical random surfaces:** quantum gravity, planar maps, string theory
4. **Canonical random paths:** walks, interfaces, Schramm-Loewner evolution
5. **Canonical random growth:** Eden model, DLA, DBM

Part II: Drama

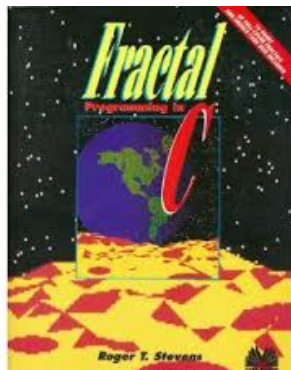
1. **Welding random surfaces:** a calculus of random surfaces and SLE seams
2. **Mating random trees:** tree plus tree (conformally mated) equals surface plus path
3. **Random growth on random surfaces:** dendrites, dragons, surprising tractability

References:

1. *Conformal weldings of random surfaces: SLE and the quantum gravity zipper* (2010)
2. *Imaginary Geometry I-IV* (Miller, S., 2012-2013)
3. *Quantum Loewner Evolution* (Miller, S. 2013)
4. *Liouville quantum gravity as a mating of trees* (Duplantier, Miller, S. 2014)

FRACTALS FROM COMPLEX DYNAMICS

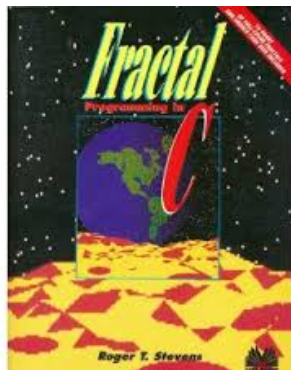
- ▶ Julia sets (Julia, 1918), popularized in 1980's



Published 1989, by Roger T. Stevens

FRACTALS FROM COMPLEX DYNAMICS

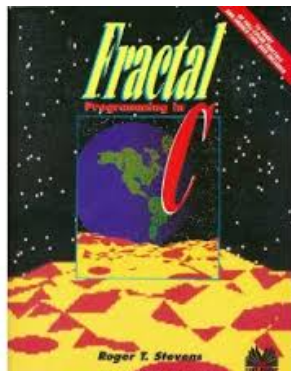
- ▶ Julia sets (Julia, 1918), popularized in 1980's
- ▶ Consider map $\phi(z) = z^2$.



Published 1989, by Roger T. Stevens

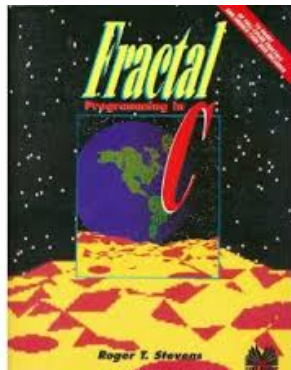
FRACTALS FROM COMPLEX DYNAMICS

- ▶ Julia sets (Julia, 1918), popularized in 1980's
- ▶ Consider map $\phi(z) = z^2$.
- ▶ Maps $\mathbf{C} \setminus \overline{D}$ conformally to self (2 to 1) where D is unit disc. Repeated iteration takes points in $\mathbf{C} \setminus \overline{D}$ to ∞ , leaves others bounded.



Published 1989, by Roger T. Stevens

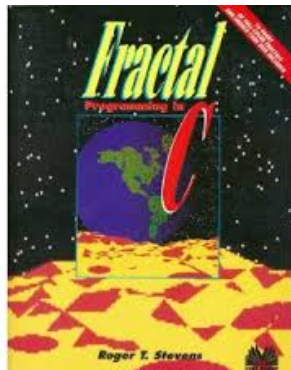
FRACTALS FROM COMPLEX DYNAMICS



Published 1989, by Roger T. Stevens

- ▶ Julia sets (Julia, 1918), popularized in 1980's
- ▶ Consider map $\phi(z) = z^2$.
- ▶ Maps $\mathbf{C} \setminus \overline{D}$ conformally to self (2 to 1) where D is unit disc. Repeated iteration takes points in $\mathbf{C} \setminus \overline{D}$ to ∞ , leaves others bounded.
- ▶ If K is another compact set with connected hull, can construct a similar (2 to 1) conformal map ϕ_K from $\mathbf{C} \setminus \overline{K}$ to itself.

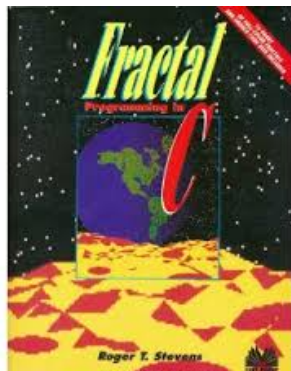
FRACTALS FROM COMPLEX DYNAMICS



Published 1989, by Roger T. Stevens

- ▶ Julia sets (Julia, 1918), popularized in 1980's
- ▶ Consider map $\phi(z) = z^2$.
- ▶ Maps $\mathbf{C} \setminus \overline{D}$ conformally to self (2 to 1) where D is unit disc. Repeated iteration takes points in $\mathbf{C} \setminus \overline{D}$ to ∞ , leaves others bounded.
- ▶ If K is another compact set with connected hull, can construct a similar (2 to 1) conformal map ϕ_K from $\mathbf{C} \setminus \overline{K}$ to itself.
- ▶ Might expect more intricate sets K to yield more intricate maps. But suppose we take $\phi_K(z) = z^2 + c$ and let K be set of points remaining bounded under repeated iteration.

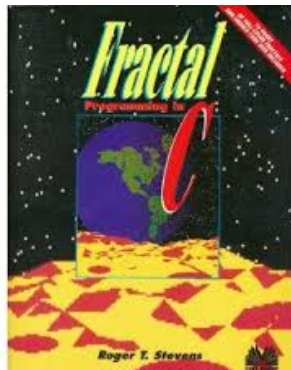
FRACTALS FROM COMPLEX DYNAMICS



Published 1989, by Roger T. Stevens

- ▶ Julia sets (Julia, 1918), popularized in 1980's
- ▶ Consider map $\phi(z) = z^2$.
- ▶ Maps $\mathbf{C} \setminus \overline{D}$ conformally to self (2 to 1) where D is unit disc. Repeated iteration takes points in $\mathbf{C} \setminus \overline{D}$ to ∞ , leaves others bounded.
- ▶ If K is another compact set with connected hull, can construct a similar (2 to 1) conformal map ϕ_K from $\mathbf{C} \setminus \overline{K}$ to itself.
- ▶ Might expect more intricate sets K to yield more intricate maps. But suppose we take $\phi_K(z) = z^2 + c$ and let K be set of points remaining bounded under repeated iteration.
- ▶ Look at Julia set images, Arnaud Chéritat's "mating" animations.

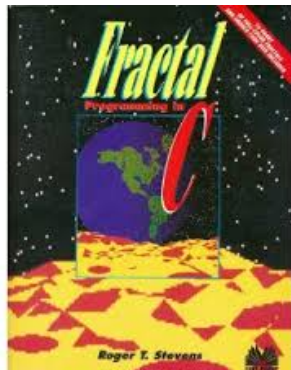
FRACTALS FROM COMPLEX DYNAMICS



Published 1989, by Roger T. Stevens

- ▶ Julia sets (Julia, 1918), popularized in 1980's
- ▶ Consider map $\phi(z) = z^2$.
- ▶ Maps $\mathbf{C} \setminus \overline{D}$ conformally to self (2 to 1) where D is unit disc. Repeated iteration takes points in $\mathbf{C} \setminus \overline{D}$ to ∞ , leaves others bounded.
- ▶ If K is another compact set with connected hull, can construct a similar (2 to 1) conformal map ϕ_K from $\mathbf{C} \setminus \overline{K}$ to itself.
- ▶ Might expect more intricate sets K to yield more intricate maps. But suppose we take $\phi_K(z) = z^2 + c$ and let K be set of points remaining bounded under repeated iteration.
- ▶ Look at Julia set images, Arnaud Chéritat's "mating" animations.
- ▶ Popular lexicon: chaos, butterfly effect, fractal, self-similar.

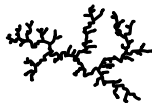
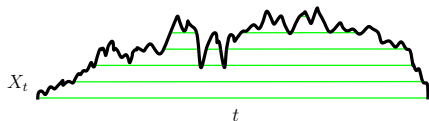
FRACTALS FROM COMPLEX DYNAMICS



Published 1989, by Roger T. Stevens

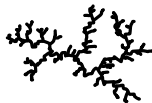
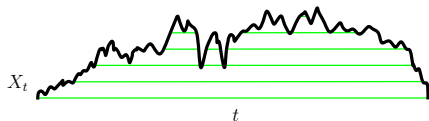
- ▶ Julia sets (Julia, 1918), popularized in 1980's
- ▶ Consider map $\phi(z) = z^2$.
- ▶ Maps $\mathbf{C} \setminus \overline{D}$ conformally to self (2 to 1) where D is unit disc. Repeated iteration takes points in $\mathbf{C} \setminus \overline{D}$ to ∞ , leaves others bounded.
- ▶ If K is another compact set with connected hull, can construct a similar (2 to 1) conformal map ϕ_K from $\mathbf{C} \setminus \overline{K}$ to itself.
- ▶ Might expect more intricate sets K to yield more intricate maps. But suppose we take $\phi_K(z) = z^2 + c$ and let K be set of points remaining bounded under repeated iteration.
- ▶ Look at Julia set images, Arnaud Chéritat's "mating" animations.
- ▶ Popular lexicon: chaos, butterfly effect, fractal, self-similar.
- ▶ What about *random* fractals, only self similar in law?

RANDOM TREES



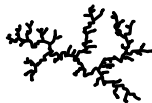
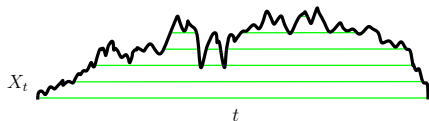
- ▶ This is the easiest random fractal to explain.

RANDOM TREES



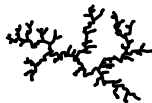
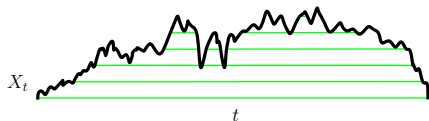
- ▶ This is the easiest random fractal to explain.
- ▶ Aldous (1993) constructs **continuum random tree** (CRT) from a Brownian excursion. To produce tree, start with graph of Brownian excursion and then identify points connected by horizontal line segment that lies below graph except at endpoints. Result is a random metric space.

RANDOM TREES



- ▶ This is the easiest random fractal to explain.
- ▶ Aldous (1993) constructs **continuum random tree** (CRT) from a Brownian excursion. To produce tree, start with graph of Brownian excursion and then identify points connected by horizontal line segment that lies below graph except at endpoints. Result is a random metric space.
- ▶ Discrete analog: Consider a tree embedded in the plane with n edges and a distinguished root. As one traces the outer boundary of the tree clockwise, distance from root performs a simple walk on \mathbf{Z}_+ with $2n$ steps, starting and ending at 0.

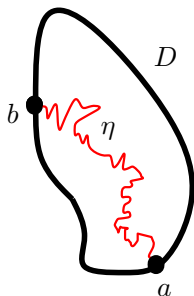
RANDOM TREES



- ▶ This is the easiest random fractal to explain.
- ▶ Aldous (1993) constructs **continuum random tree** (CRT) from a Brownian excursion. To produce tree, start with graph of Brownian excursion and then identify points connected by horizontal line segment that lies below graph except at endpoints. Result is a random metric space.
- ▶ Discrete analog: Consider a tree embedded in the plane with n edges and a distinguished root. As one traces the outer boundary of the tree clockwise, distance from root performs a simple walk on \mathbf{Z}_+ with $2n$ steps, starting and ending at 0.
- ▶ Simple bijection rooted planar trees and walks of this type.

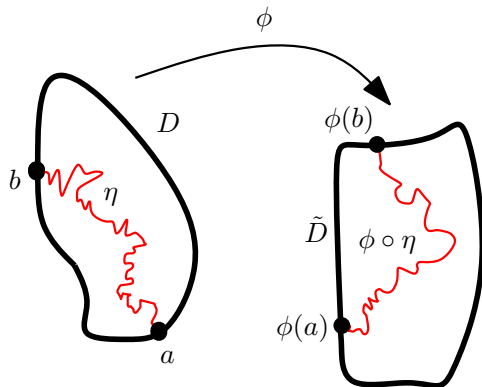
RANDOM PATHS

Given a simply connected planar domain D with boundary points a and b and a parameter $\kappa \in [0, \infty)$, the **Schramm-Loewner evolution** SLE_κ is a random non-self-crossing path in \overline{D} from a to b .



The parameter κ roughly indicates how “windy” the path is. Would like to argue that SLE is in some sense the “canonical” random non-self-crossing path. What symmetries characterize SLE?

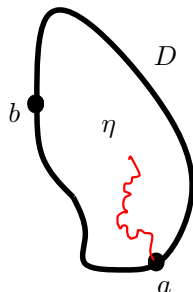
Conformal Markov property of SLE



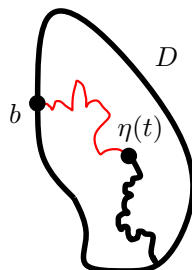
If ϕ conformally maps D to \tilde{D} and η is an SLE_κ from a to b in D , then $\phi \circ \eta$ is an SLE_κ from $\phi(a)$ to $\phi(b)$ in \tilde{D} .

Markov Property

Given η up to a stopping time t ...



law of remainder is SLE in $D \setminus \eta[0, t]$ from $\eta(t)$ to b .



Chordal Schramm-Loewner evolution (SLE)

- ▶ **THEOREM [Oded Schramm]:** Conformal invariance and the Markov property completely determine the law of SLE, up to a single parameter which we denote by $\kappa \geq 0$.

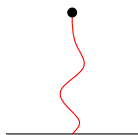
Chordal Schramm-Loewner evolution (SLE)

- ▶ **THEOREM [Oded Schramm]:** Conformal invariance and the Markov property completely determine the law of SLE, up to a single parameter which we denote by $\kappa \geq 0$.
- ▶ **Explicit construction:** An SLE path γ from 0 to ∞ in the complex upper half plane \mathbf{H} can be defined in an interesting way: given path γ one can construct conformal maps $g_t : \mathbf{H} \setminus \gamma([0, t]) \rightarrow \mathbf{H}$ (normalized to look like identity near infinity, i.e., $\lim_{z \rightarrow \infty} g_t(z) - z = 0$). In SLE_κ , one defines g_t via an ODE (which makes sense for each fixed z):

$$\partial_t g_t(z) = \frac{2}{g_t(z) - W_t}, \quad g_0(z) = z,$$

where $W_t = \sqrt{\kappa} B_t \stackrel{\text{LAW}}{=} B_{\kappa t}$ and B_t is ordinary Brownian motion.

SLE phases [Rohde, Schramm]



$$\kappa \leq 4$$



$$\kappa \in (4, 8)$$



$$\kappa \geq 8$$

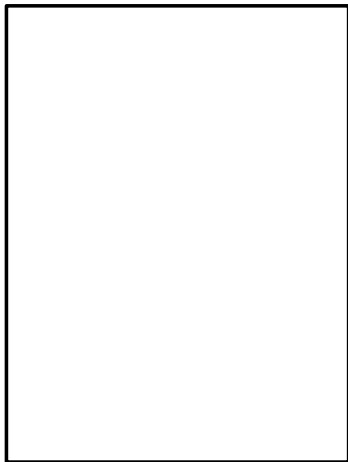
Radial Schramm-Loewner evolution (SLE)

- ▶ **Radial SLE:** $\partial_t g_t(z) = g_t(z) \frac{\xi_t + g_t(z)}{\xi_t - g_t(z)}$ where $\xi_t = e^{i\sqrt{\kappa}B_t}$.

Radial Schramm-Loewner evolution (SLE)

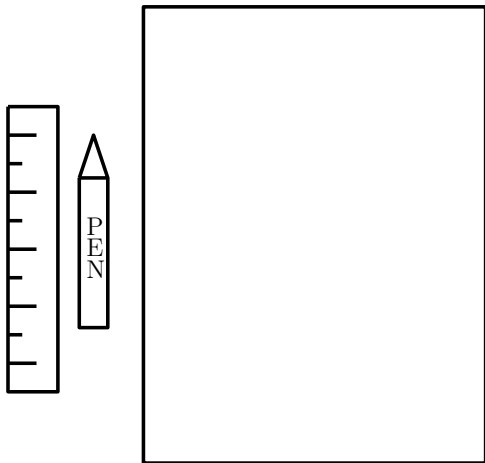
- ▶ **Radial SLE:** $\partial_t g_t(z) = g_t(z) \frac{\xi_t + g_t(z)}{\xi_t - g_t(z)}$ where $\xi_t = e^{i\sqrt{\kappa}B_t}$.
- ▶ **Radial measure-driven Loewner evolution:** $\partial_t g_t(z) = \int g_t(z) \frac{x + g_t(z)}{x - g_t(z)} dm_t(x)$ where, for each g , m_t is a measure on the complex unit circle.

RANDOM SURFACES



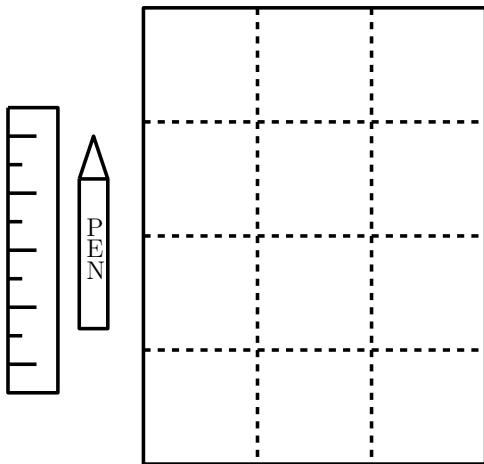
Start out with a sheet of paper

RANDOM SURFACES



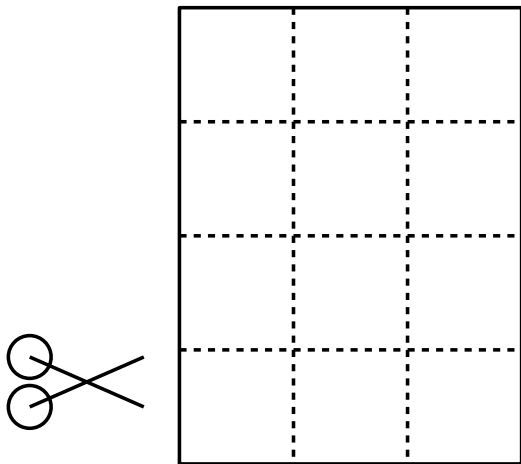
Get out pen and ruler

RANDOM SURFACES



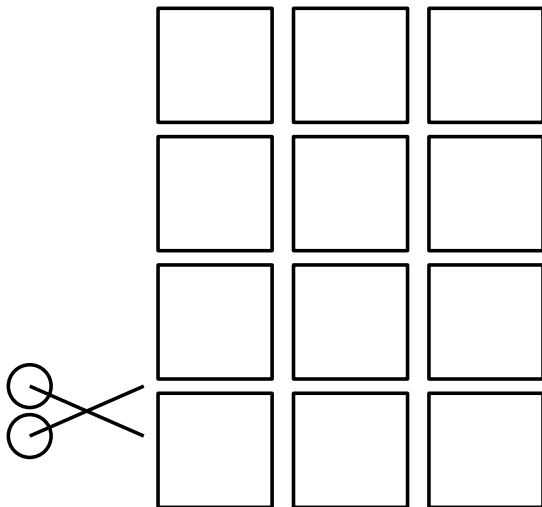
Measure and mark squares squares of equal size

RANDOM SURFACES



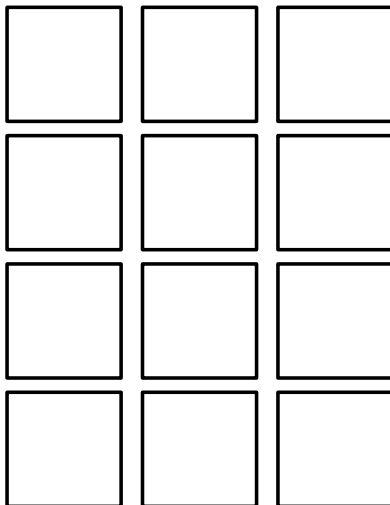
Get out scissors

RANDOM SURFACES



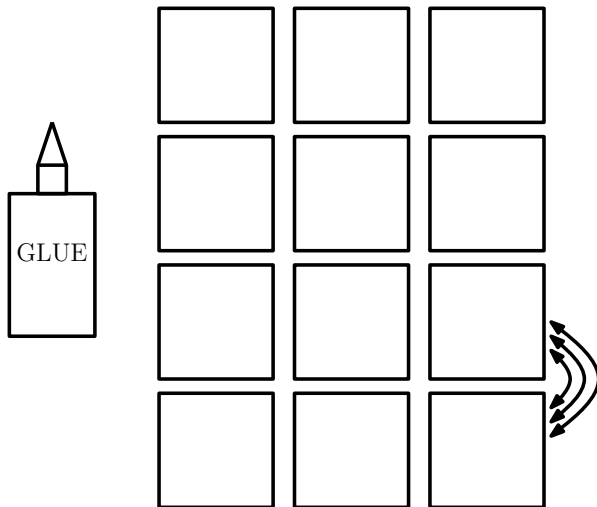
Cut into squares

RANDOM SURFACES

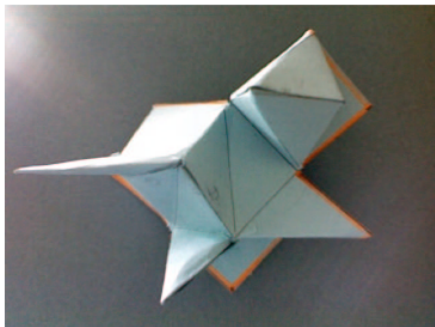
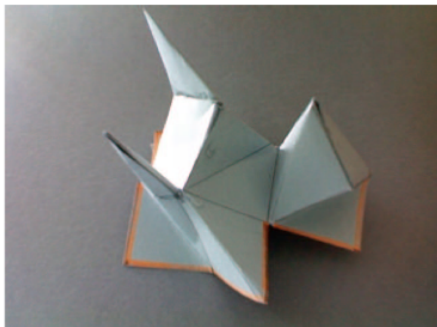


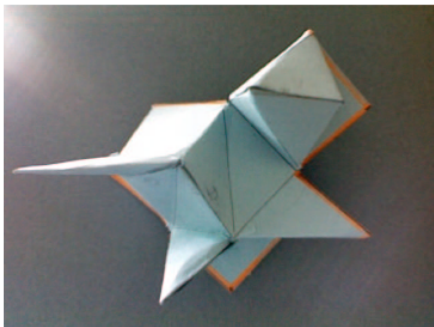
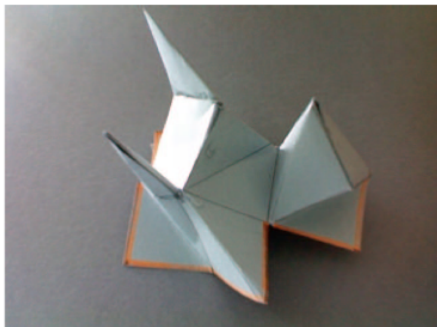
Get out bottle of glue

RANDOM SURFACES



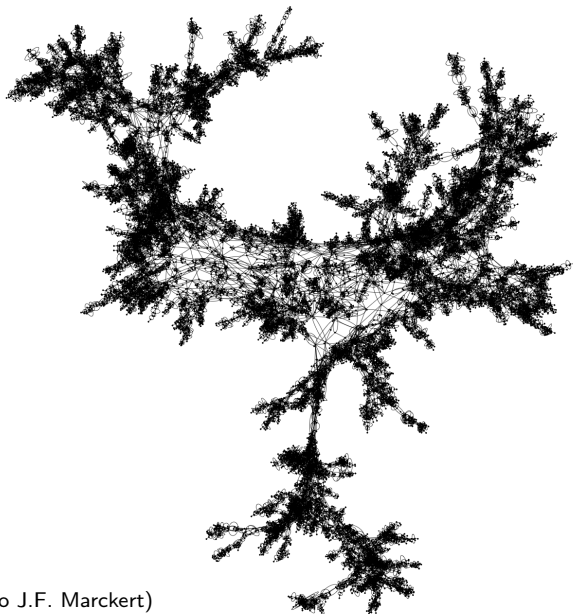
Attach squares along boundaries with glue to form a surface “without holes.”





What is the structure of a typical quadrangulation when the number of faces is large?

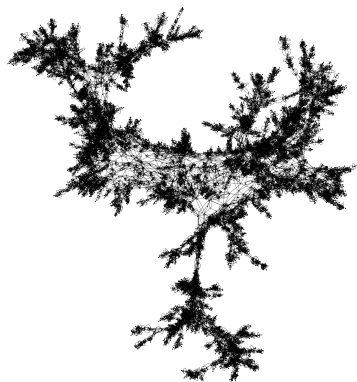
Random quadrangulation with 25,000 faces



(Simulation due to J.F. Marckert)

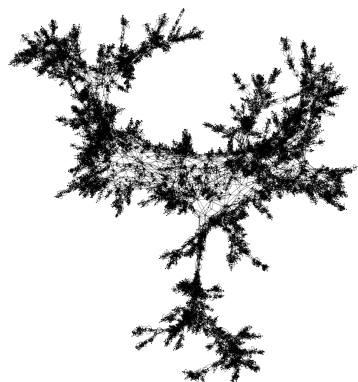
Background

1. First studied by Tutte in 1960s while working on the four color theorem.



(Simulation due to J.F. Marckert)

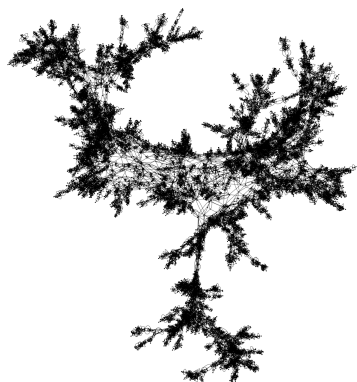
Background



(Simulation due to J.F. Marckert)

1. First studied by Tutte in 1960s while working on the four color theorem.
2. Many variants (triangulations, quadrangulations, etc.) Some come equipped with extra statistical physics structure (a distinguished spanning tree, a general distinguished edge subset, a “spin” function on vertices, etc.)

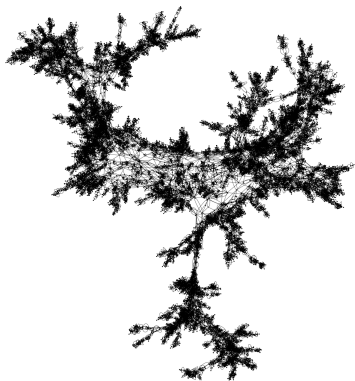
Background



(Simulation due to J.F. Marckert)

1. First studied by Tutte in 1960s while working on the four color theorem.
2. Many variants (triangulations, quadrangulations, etc.) Some come equipped with extra statistical physics structure (a distinguished spanning tree, a general distinguished edge subset, a “spin” function on vertices, etc.)
3. Can be interpreted as Riemannian manifolds with conical singularities.

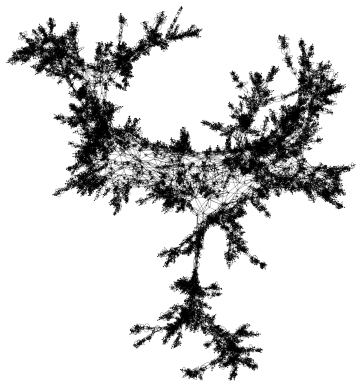
Background



(Simulation due to J.F. Marckert)

1. First studied by Tutte in 1960s while working on the four color theorem.
2. Many variants (triangulations, quadrangulations, etc.) Some come equipped with extra statistical physics structure (a distinguished spanning tree, a general distinguished edge subset, a “spin” function on vertices, etc.)
3. Can be interpreted as Riemannian manifolds with conical singularities.
4. Converges in law in Gromov-Hausdorff sense to random metric space called Brownian map, homeomorphic to the 2-sphere, Hausdorff dimension 4 (established in several works by subsets of Chaissang, Schaefer, Le Gall, Paulin, Miermont)

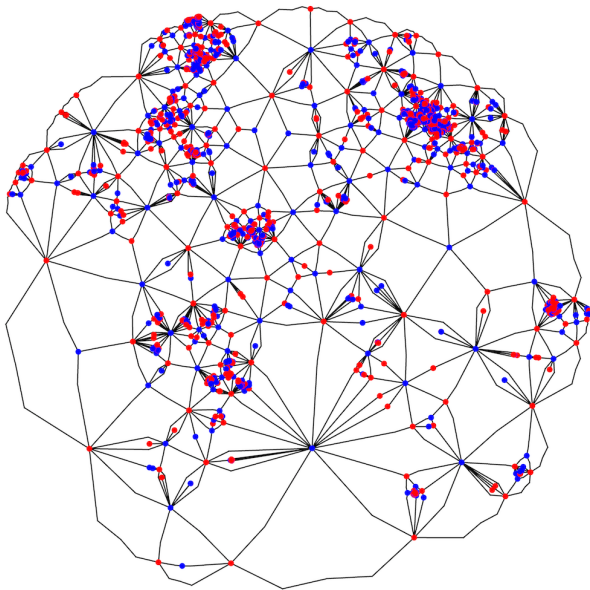
Background



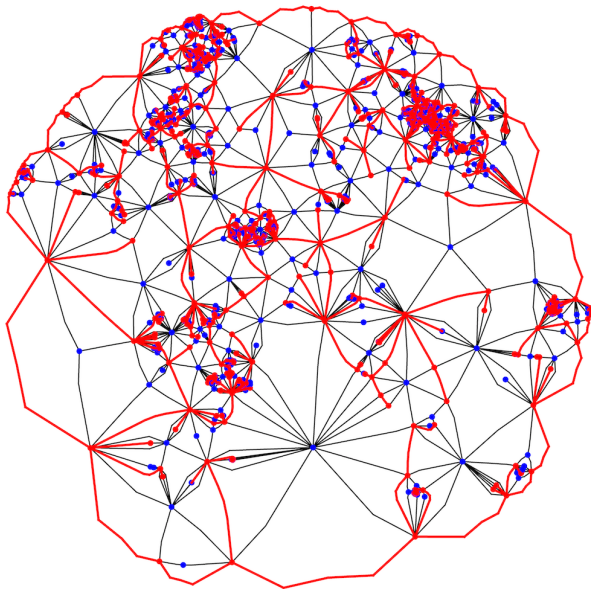
(Simulation due to J.F. Marckert)

1. First studied by Tutte in 1960s while working on the four color theorem.
2. Many variants (triangulations, quadrangulations, etc.) Some come equipped with extra statistical physics structure (a distinguished spanning tree, a general distinguished edge subset, a “spin” function on vertices, etc.)
3. Can be interpreted as Riemannian manifolds with conical singularities.
4. Converges in law in Gromov-Hausdorff sense to random metric space called Brownian map, homeomorphic to the 2-sphere, Hausdorff dimension 4 (established in several works by subsets of Chaissang, Schaefer, Le Gall, Paulin, Miermont)
5. Important tool: Bijections encoding surface via pair of trees.

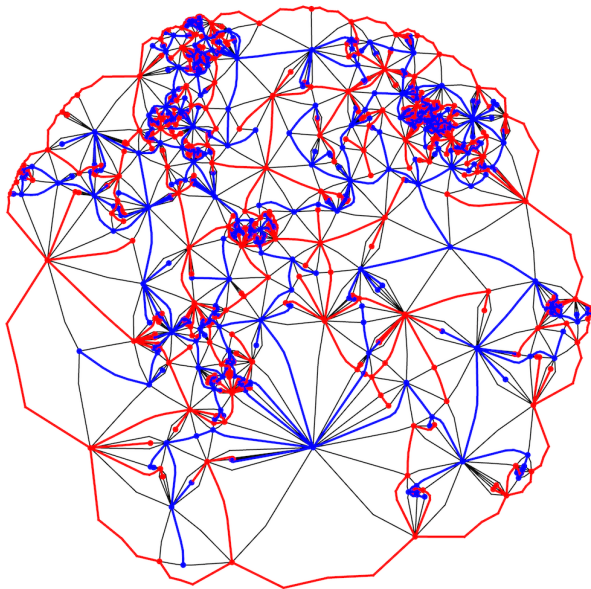
Random quadrangulation



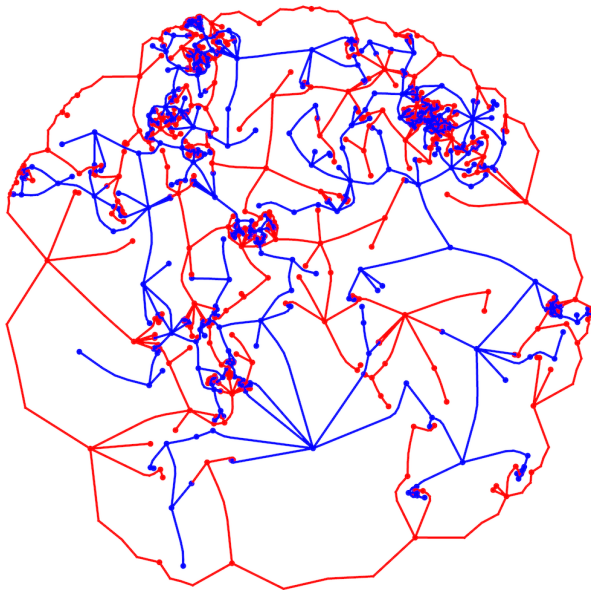
Red tree



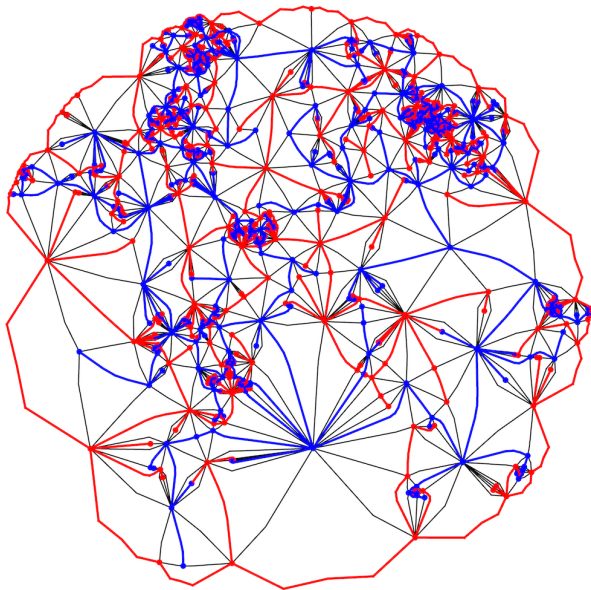
Red and blue trees



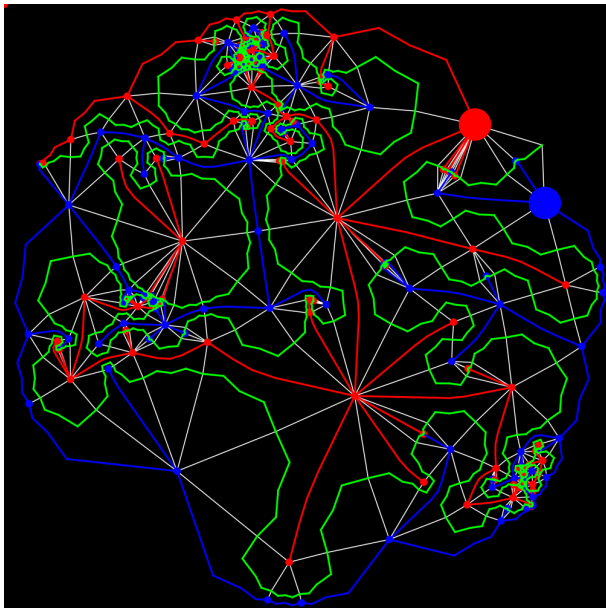
Red and blue trees alone do not determine the map structure



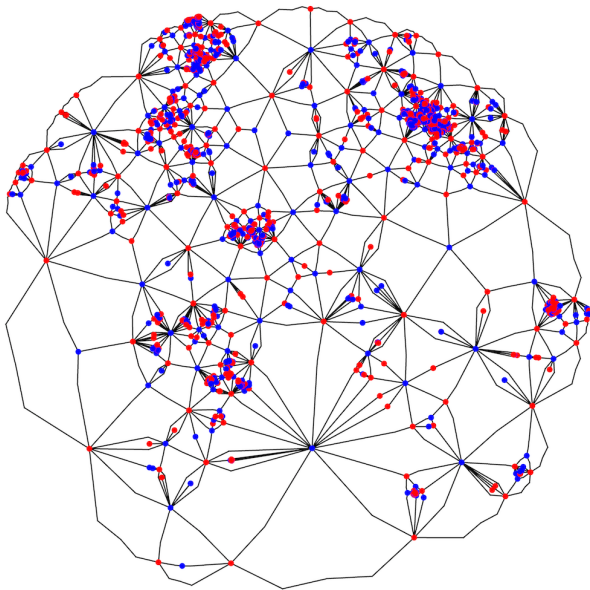
Random quadrangulation with red and blue trees



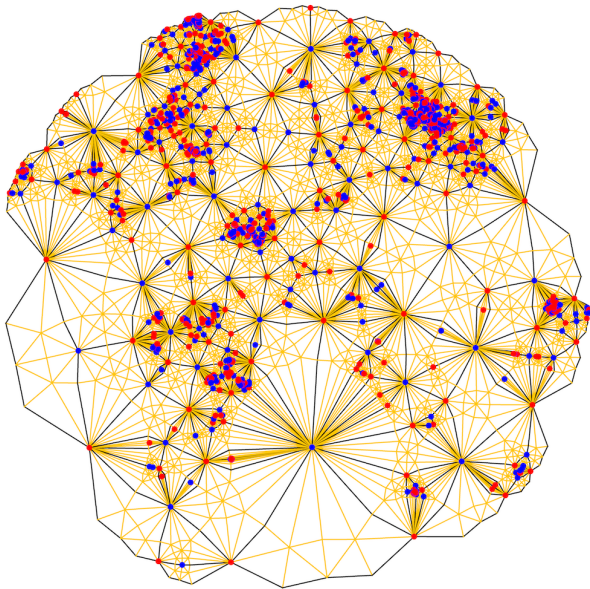
Path snaking between the trees. Encodes the trees and how they are glued together.



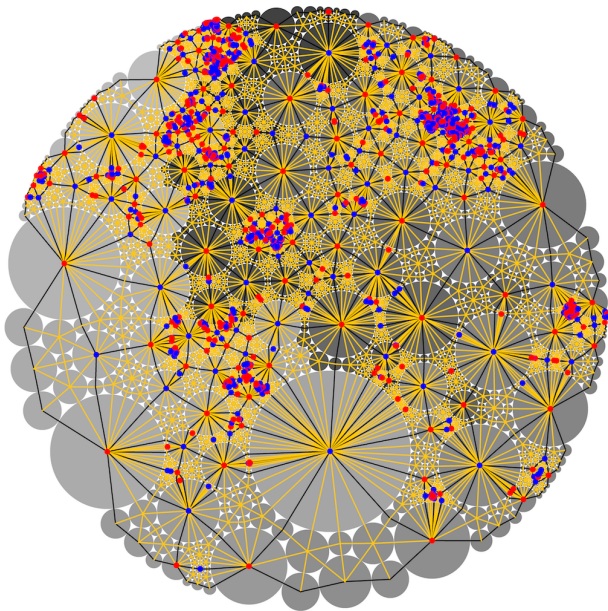
How was the graph embedded into \mathbf{R}^2 ?



Can subdivide each quadrilateral to obtain a triangulation without multiple edges.

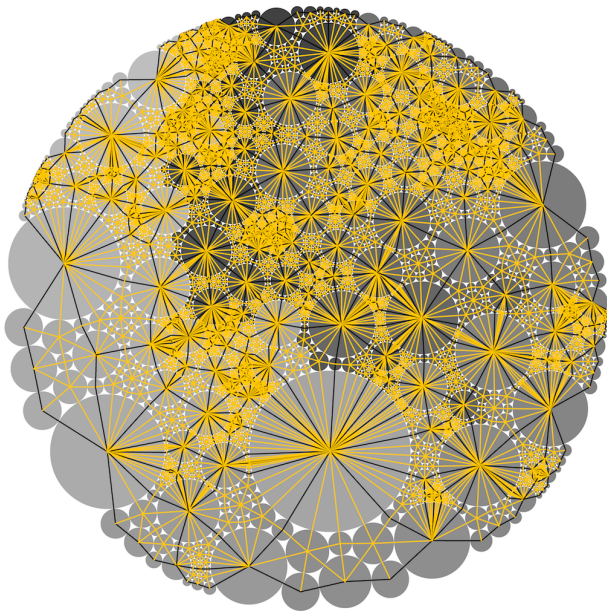


Circle pack the resulting triangulation.



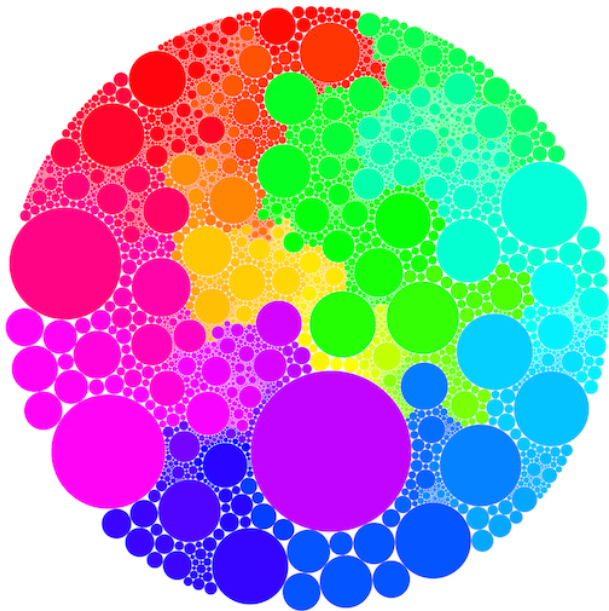
Packed with Stephenson's CirclePack.

Circle pack the resulting triangulation.



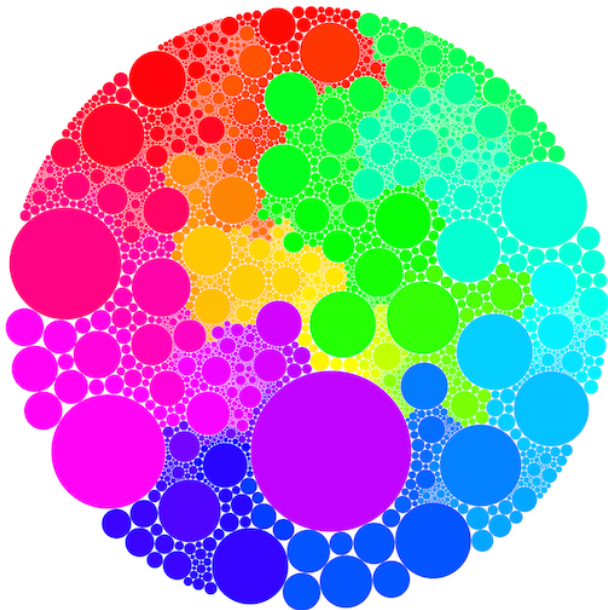
Packed with Stephenson's CirclePack.

Circle pack the resulting triangulation.



Packed with Stephenson's CirclePack.

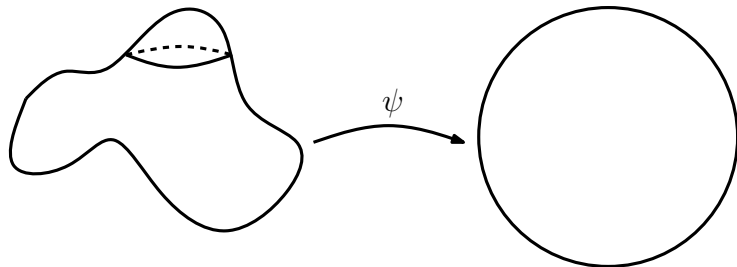
What is the “limit” of this embedding? Circle packings are related to conformal maps.



Packed with Stephenson's CirclePack.

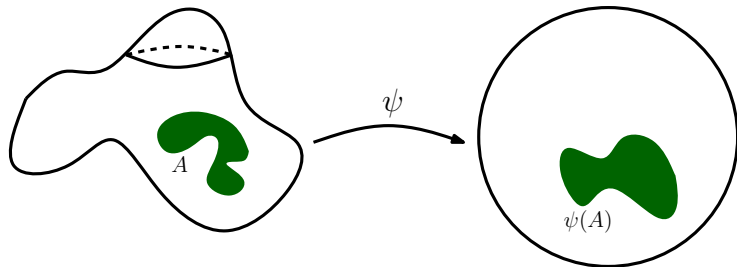
Picking a surface at random in the continuum

Uniformization theorem: every simply connected Riemannian surface can be conformally mapped to either the unit disk, the plane, or the sphere \mathbf{S}^2 in \mathbf{R}^3



Picking a surface at random in the continuum

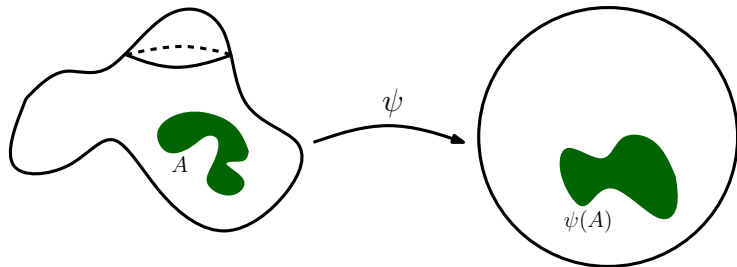
Uniformization theorem: every simply connected Riemannian surface can be conformally mapped to either the unit disk, the plane, or the sphere \mathbf{S}^2 in \mathbf{R}^3



Isothermal coordinates: Metric for the surface takes the form $e^{\rho(z)} dz$ for some smooth function ρ where dz is the Euclidean metric. (See David Gu's gallery.)

Picking a surface at random in the continuum

Uniformization theorem: every simply connected Riemannian surface can be conformally mapped to either the unit disk, the plane, or the sphere \mathbf{S}^2 in \mathbf{R}^3



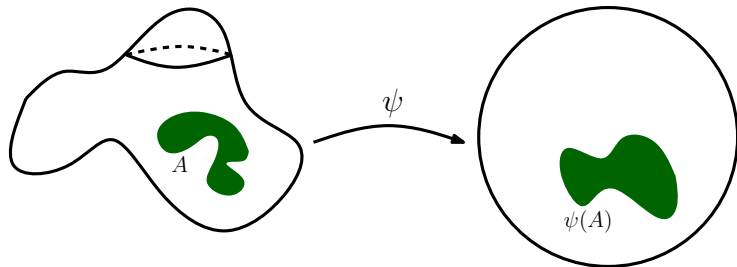
Isothermal coordinates: Metric for the surface takes the form $e^{\rho(z)} dz$ for some smooth function ρ where dz is the Euclidean metric. (See David Gu's gallery.)

⇒ Can parameterize the space of surfaces with smooth functions.

- ▶ If $\rho = 0$, get the same surface
- ▶ If $\Delta\rho = 0$, i.e. if ρ is harmonic, the surface described is flat

Picking a surface at random in the continuum

Uniformization theorem: every simply connected Riemannian surface can be conformally mapped to either the unit disk, the plane, or the sphere \mathbf{S}^2 in \mathbf{R}^3



Isothermal coordinates: Metric for the surface takes the form $e^{\rho(z)} dz$ for some smooth function ρ where dz is the Euclidean metric. (See David Gu's gallery.)

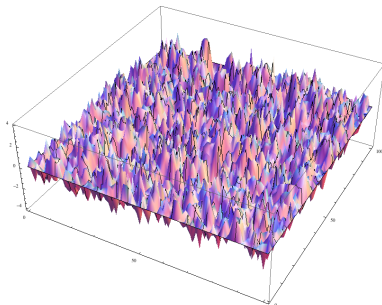
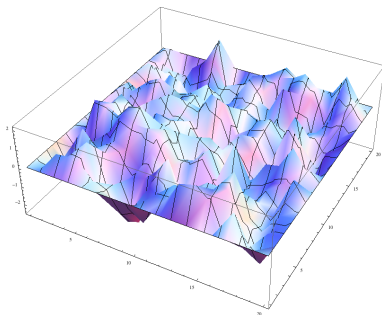
⇒ Can parameterize the space of surfaces with smooth functions.

- ▶ If $\rho = 0$, get the same surface
- ▶ If $\Delta\rho = 0$, i.e. if ρ is harmonic, the surface described is flat

Question: Which measure on ρ ? If we want our surface to be a perturbation of a flat metric, natural to choose ρ as the canonical perturbation of a harmonic function.

The Gaussian free field

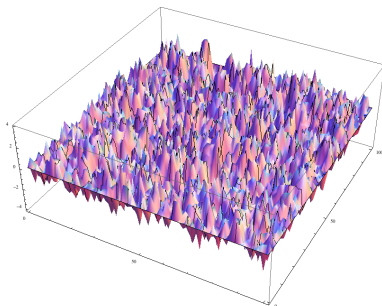
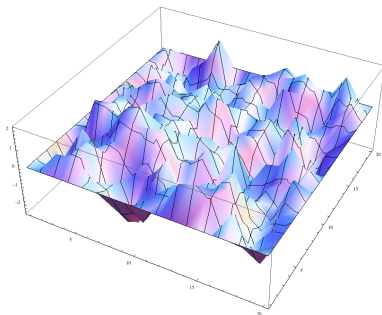
- ▶ The **discrete Gaussian free field** (DGFF) is a **Gaussian random surface** model.



The Gaussian free field

- ▶ The **discrete Gaussian free field** (DGFF) is a **Gaussian random surface** model.
- ▶ Measure on functions $h: D \rightarrow \mathbf{R}$ for $D \subseteq \mathbf{Z}^2$ and $h|_{\partial D} = \psi$ with density respect to Lebesgue measure on $\mathbf{R}^{|\mathcal{D}|}$:

$$\frac{1}{\mathcal{Z}} \exp \left(-\frac{1}{2} \sum_{x \sim y} (h(x) - h(y))^2 \right)$$

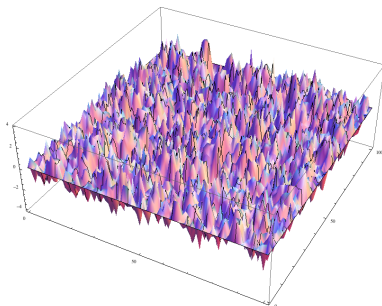
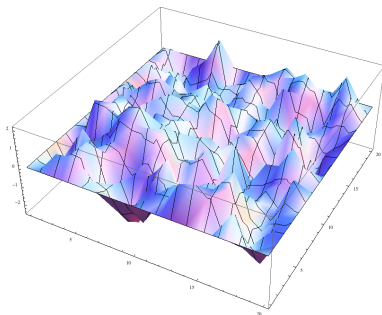


The Gaussian free field

- ▶ The **discrete Gaussian free field** (DGFF) is a **Gaussian random surface** model.
- ▶ Measure on functions $h: D \rightarrow \mathbf{R}$ for $D \subseteq \mathbf{Z}^2$ and $h|_{\partial D} = \psi$ with density respect to Lebesgue measure on $\mathbf{R}^{|\mathcal{D}|}$:

$$\frac{1}{\mathcal{Z}} \exp \left(-\frac{1}{2} \sum_{x \sim y} (h(x) - h(y))^2 \right)$$

- ▶ Natural perturbation of a harmonic function



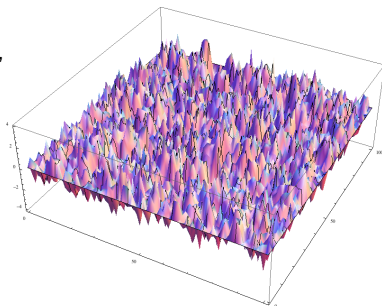
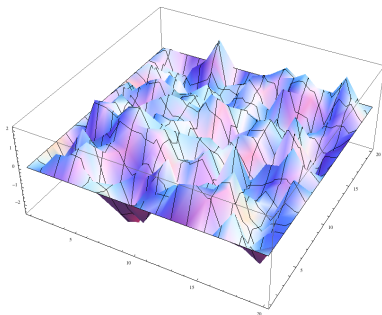
The Gaussian free field

- ▶ The **discrete Gaussian free field** (DGFF) is a **Gaussian random surface** model.
- ▶ Measure on functions $h: D \rightarrow \mathbf{R}$ for $D \subseteq \mathbf{Z}^2$ and $h|_{\partial D} = \psi$ with density respect to Lebesgue measure on $\mathbf{R}^{|D|}$:

$$\frac{1}{\mathcal{Z}} \exp \left(-\frac{1}{2} \sum_{x \sim y} (h(x) - h(y))^2 \right)$$

- ▶ Natural perturbation of a harmonic function
- ▶ Fine mesh limit: converges to the continuum GFF, i.e. the standard Gaussian wrt the **Dirichlet inner product**

$$(f, g)_{\nabla} = \frac{1}{2\pi} \int \nabla f(x) \cdot \nabla g(x) dx.$$



The Gaussian free field

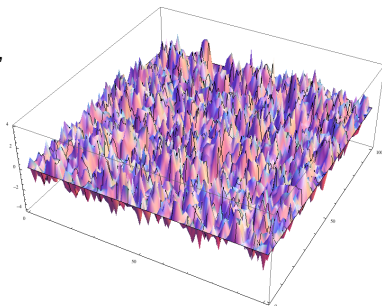
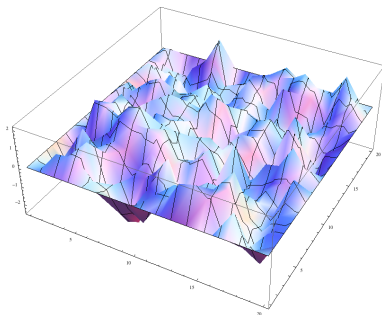
- ▶ The **discrete Gaussian free field** (DGFF) is a **Gaussian random surface** model.
- ▶ Measure on functions $h: D \rightarrow \mathbf{R}$ for $D \subseteq \mathbf{Z}^2$ and $h|_{\partial D} = \psi$ with density respect to Lebesgue measure on $\mathbf{R}^{|D|}$:

$$\frac{1}{\mathcal{Z}} \exp \left(-\frac{1}{2} \sum_{x \sim y} (h(x) - h(y))^2 \right)$$

- ▶ Natural perturbation of a harmonic function
- ▶ Fine mesh limit: converges to the continuum GFF, i.e. the standard Gaussian wrt the **Dirichlet inner product**

$$(f, g)_{\nabla} = \frac{1}{2\pi} \int \nabla f(x) \cdot \nabla g(x) dx.$$

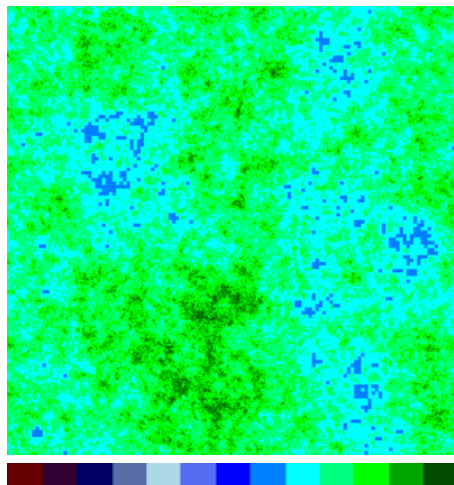
- ▶ Continuum GFF not a function — only a generalized function



Liouville quantum gravity

- ▶ Liouville quantum gravity: $e^{\gamma h(z)} dz$
where h is a GFF and $\gamma \in [0, 2)$

$$\gamma = 0.5$$

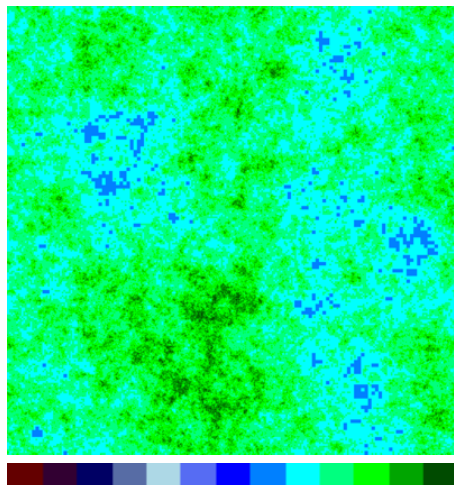


(Number of subdivisions)

Liouville quantum gravity

- ▶ Liouville quantum gravity: $e^{\gamma h(z)} dz$
where h is a GFF and $\gamma \in [0, 2)$
- ▶ Introduced by Polyakov in the 1980s

$$\gamma = 0.5$$

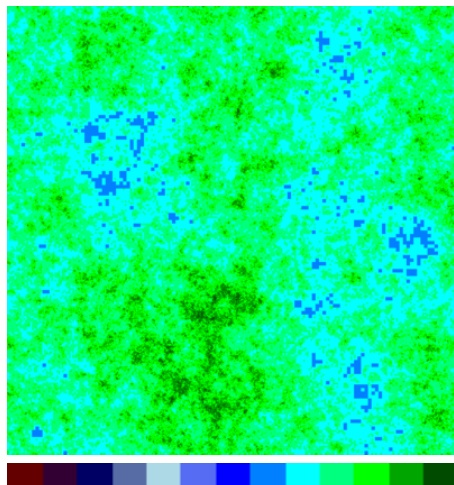


(Number of subdivisions)

Liouville quantum gravity

- ▶ Liouville quantum gravity: $e^{\gamma h(z)} dz$ where h is a GFF and $\gamma \in [0, 2)$
- ▶ Introduced by Polyakov in the 1980s
- ▶ Does not make literal sense since h takes values in the space of distributions

$$\gamma = 0.5$$

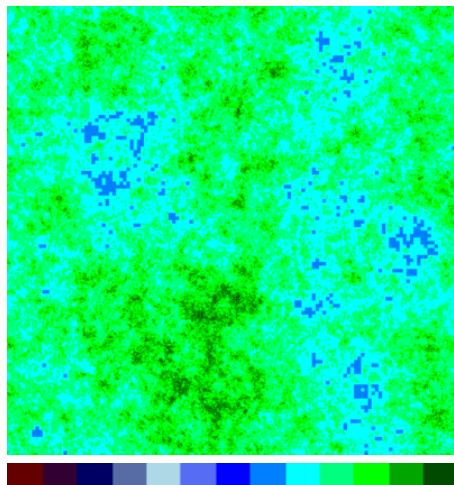


(Number of subdivisions)

Liouville quantum gravity

- ▶ Liouville quantum gravity: $e^{\gamma h(z)} dz$ where h is a GFF and $\gamma \in [0, 2)$
- ▶ Introduced by Polyakov in the 1980s
- ▶ Does not make literal sense since h takes values in the space of distributions
- ▶ Can be made sense of as a random area measure using a regularization procedure
 - ▶ Can compute areas of regions and lengths of curves

$$\gamma = 0.5$$

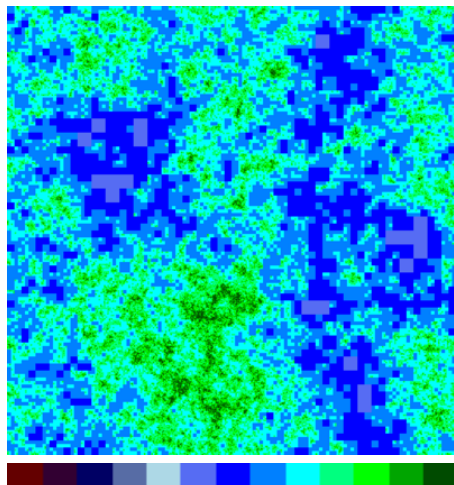


(Number of subdivisions)

Liouville quantum gravity

- ▶ Liouville quantum gravity: $e^{\gamma h(z)} dz$ where h is a GFF and $\gamma \in [0, 2)$
- ▶ Introduced by Polyakov in the 1980s
- ▶ Does not make literal sense since h takes values in the space of distributions
- ▶ Can be made sense of as a random area measure using a regularization procedure
 - ▶ Can compute areas of regions and lengths of curves

$$\gamma = 1.0$$

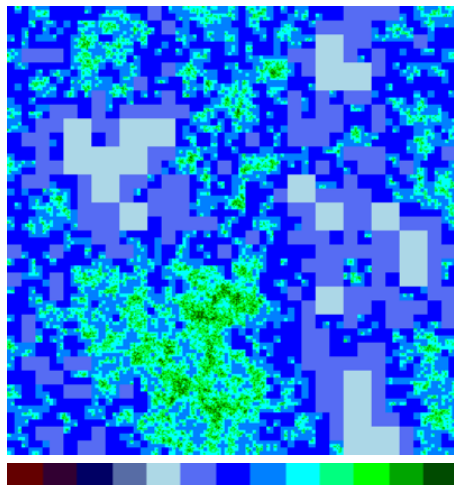


(Number of subdivisions)

Liouville quantum gravity

- ▶ Liouville quantum gravity: $e^{\gamma h(z)} dz$ where h is a GFF and $\gamma \in [0, 2)$
- ▶ Introduced by Polyakov in the 1980s
- ▶ Does not make literal sense since h takes values in the space of distributions
- ▶ Can be made sense of as a random area measure using a regularization procedure
 - ▶ Can compute areas of regions and lengths of curves

$$\gamma = 1.5$$

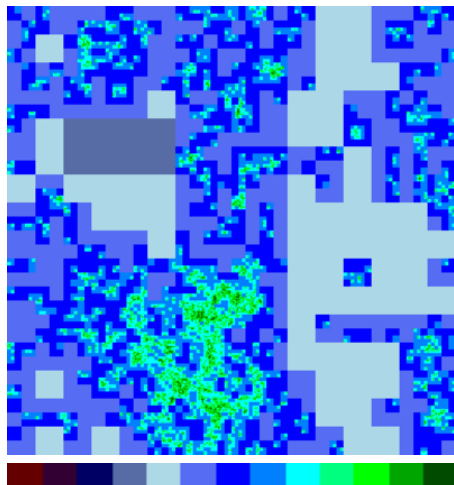


(Number of subdivisions)

Liouville quantum gravity

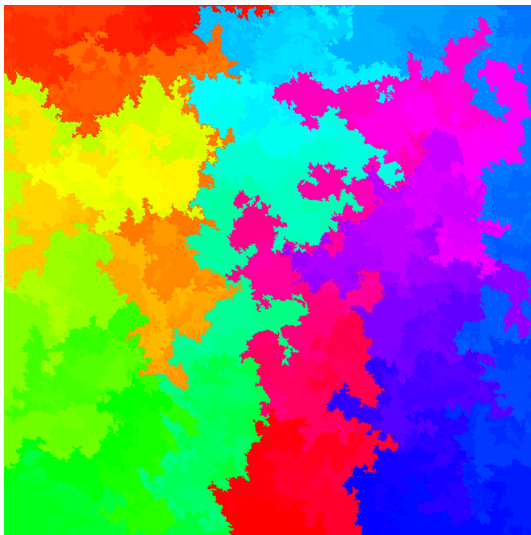
- ▶ Liouville quantum gravity: $e^{\gamma h(z)} dz$ where h is a GFF and $\gamma \in [0, 2)$
- ▶ Introduced by Polyakov in the 1980s
- ▶ Does not make literal sense since h takes values in the space of distributions
- ▶ Can be made sense of as a random area measure using a regularization procedure
 - ▶ Can compute areas of regions and lengths of curves

$$\gamma = 2.0$$



(Number of subdivisions)

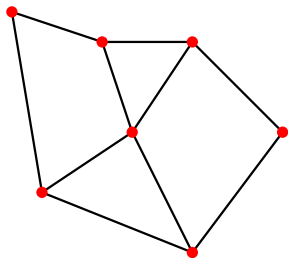
Continuum space-filling path



Space-filling SLE_6 on a LQG surface. Random path which encodes the limit of a RPM.

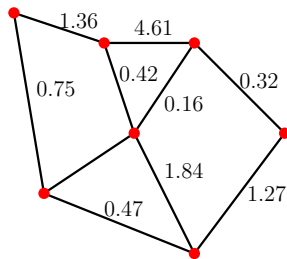
RANDOM GROWTH

- ▶ Associate with a graph (V, E) i.i.d. $\exp(1)$ edge weights



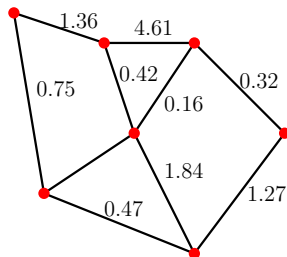
RANDOM GROWTH

- ▶ Associate with a graph (V, E) i.i.d. $\exp(1)$ edge weights



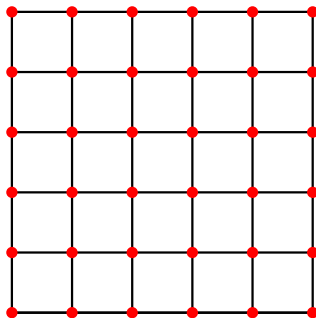
RANDOM GROWTH

- ▶ Associate with a graph (V, E) i.i.d. $\exp(1)$ edge weights
- ▶ Introduced by Eden (1961) and Hammersley and Welsh (1965)



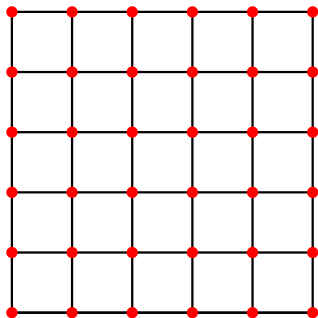
RANDOM GROWTH

- ▶ Associate with a graph (V, E) i.i.d. $\exp(1)$ edge weights
- ▶ Introduced by Eden (1961) and Hammersley and Welsh (1965)
- ▶ On \mathbf{Z}^2 ?



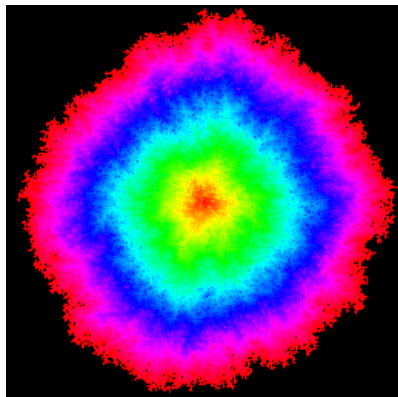
RANDOM GROWTH

- ▶ Associate with a graph (V, E) i.i.d. $\exp(1)$ edge weights
- ▶ Introduced by Eden (1961) and Hammersley and Welsh (1965)
- ▶ On \mathbf{Z}^2 ?
- ▶ **Question:** Large scale behavior of shape of ball wrt perturbed metric?



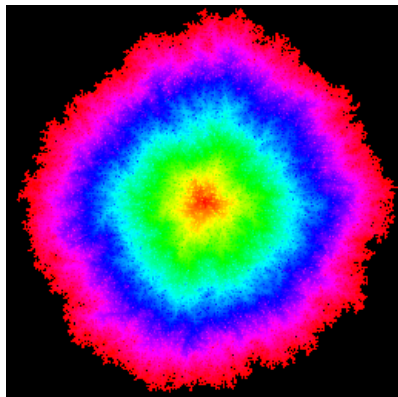
RANDOM GROWTH

- ▶ Associate with a graph (V, E) i.i.d. $\exp(1)$ edge weights
- ▶ Introduced by Eden (1961) and Hammersley and Welsh (1965)
- ▶ On \mathbf{Z}^2 ?
- ▶ **Question:** Large scale behavior of shape of ball wrt perturbed metric?



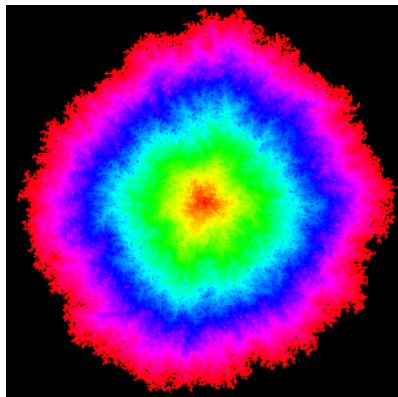
RANDOM GROWTH

- ▶ Associate with a graph (V, E) i.i.d. $\exp(1)$ edge weights
- ▶ Introduced by Eden (1961) and Hammersley and Welsh (1965)
- ▶ On \mathbf{Z}^2 ?
- ▶ **Question:** Large scale behavior of shape of ball wrt perturbed metric?
- ▶ Cox and Durrett (1981) showed that the macroscopic shape is convex



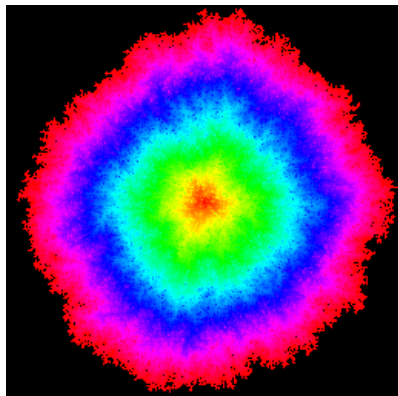
RANDOM GROWTH

- ▶ Associate with a graph (V, E) i.i.d. $\exp(1)$ edge weights
- ▶ Introduced by Eden (1961) and Hammersley and Welsh (1965)
- ▶ On \mathbf{Z}^2 ?
- ▶ **Question:** Large scale behavior of shape of ball wrt perturbed metric?
- ▶ Cox and Durrett (1981) showed that the macroscopic shape is convex
- ▶ Computer simulations show that it is not a Euclidean disk



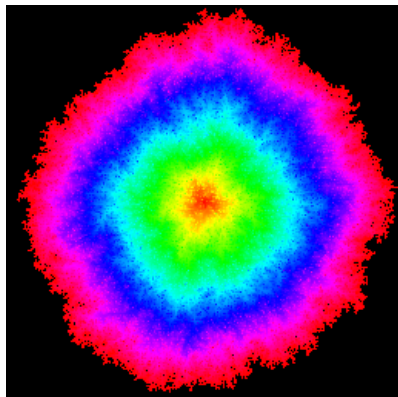
RANDOM GROWTH

- ▶ Associate with a graph (V, E) i.i.d. $\exp(1)$ edge weights
- ▶ Introduced by Eden (1961) and Hammersley and Welsh (1965)
- ▶ On \mathbf{Z}^2 ?
- ▶ **Question:** Large scale behavior of shape of ball wrt perturbed metric?
- ▶ Cox and Durrett (1981) showed that the macroscopic shape is convex
- ▶ Computer simulations show that it is not a Euclidean disk
- ▶ \mathbf{Z}^2 is not isotropic enough



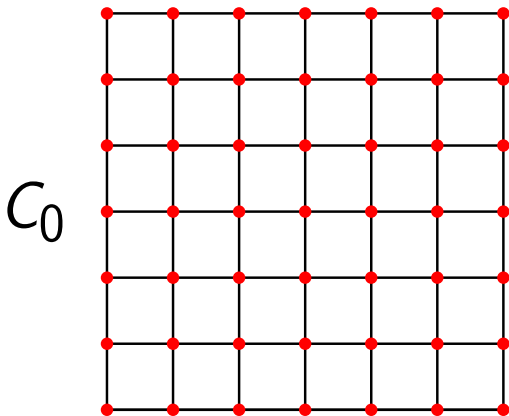
RANDOM GROWTH

- ▶ Associate with a graph (V, E) i.i.d. $\exp(1)$ edge weights
- ▶ Introduced by Eden (1961) and Hammersley and Welsh (1965)
- ▶ On \mathbf{Z}^2 ?
- ▶ **Question:** Large scale behavior of shape of ball wrt perturbed metric?
- ▶ Cox and Durrett (1981) showed that the macroscopic shape is convex
- ▶ Computer simulations show that it is not a Euclidean disk
- ▶ \mathbf{Z}^2 is not isotropic enough
- ▶ Vahidi-Asl and Weirmann (1990) showed that the rescaled ball converges to a disk if \mathbf{Z}^2 is replaced by the Voronoi tessellation associated with a Poisson process



Markovian formulation

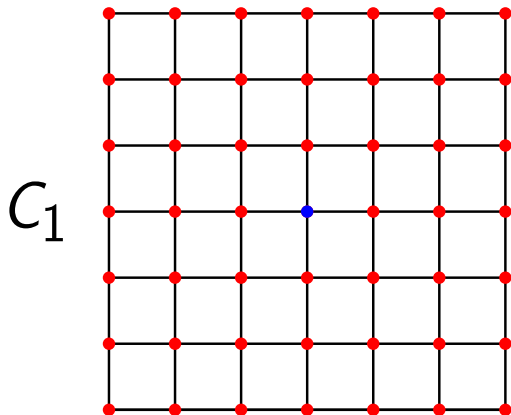
Eden exploration



Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).

Markovian formulation

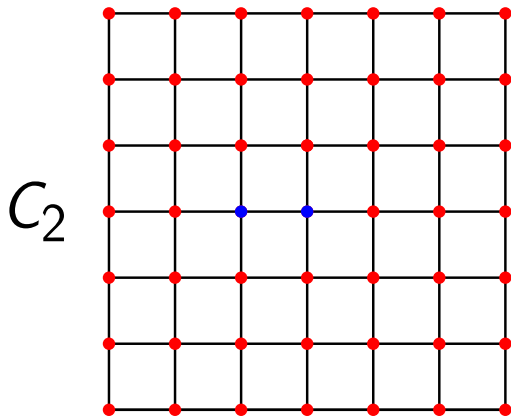
Eden exploration



Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).

Markovian formulation

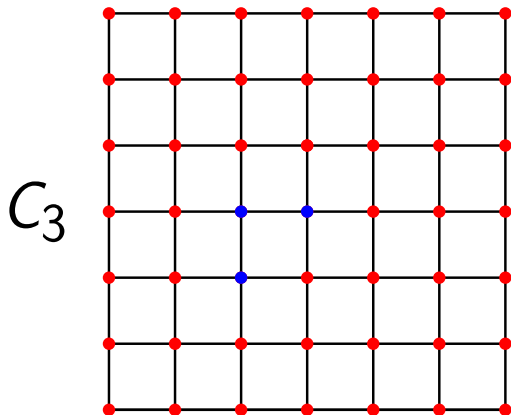
Eden exploration



Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).

Markovian formulation

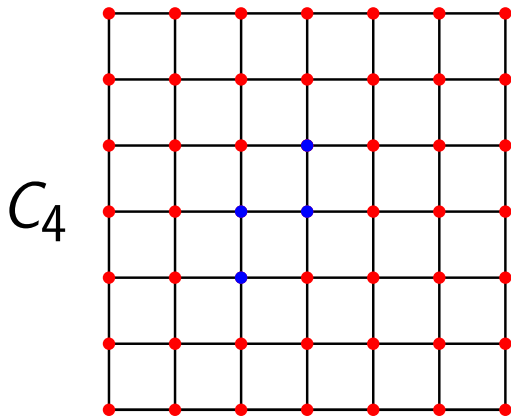
Eden exploration



Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).

Markovian formulation

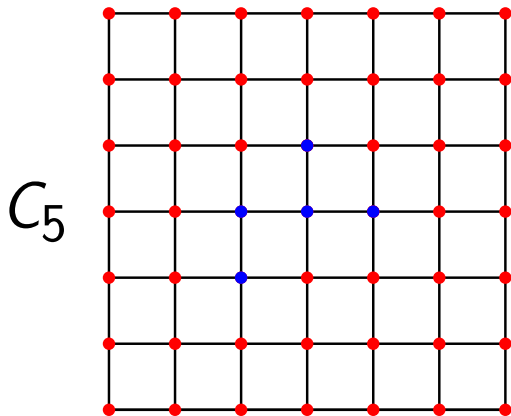
Eden exploration



Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).

Markovian formulation

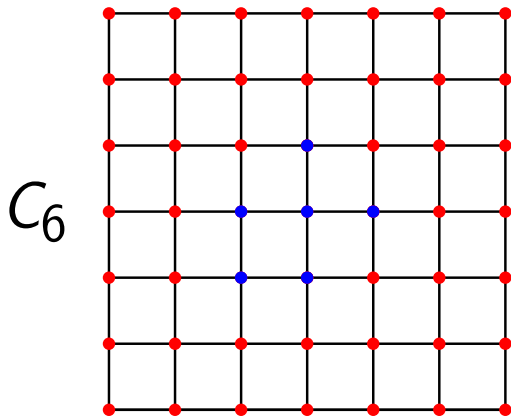
Eden exploration



Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).

Markovian formulation

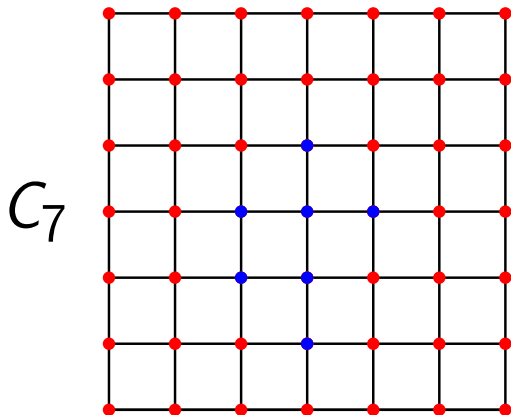
Eden exploration



Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).

Markovian formulation

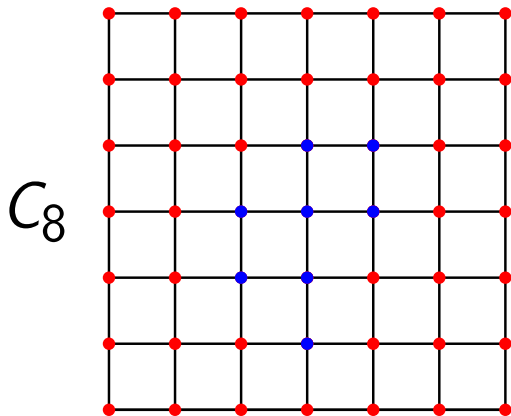
Eden exploration



Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).

Markovian formulation

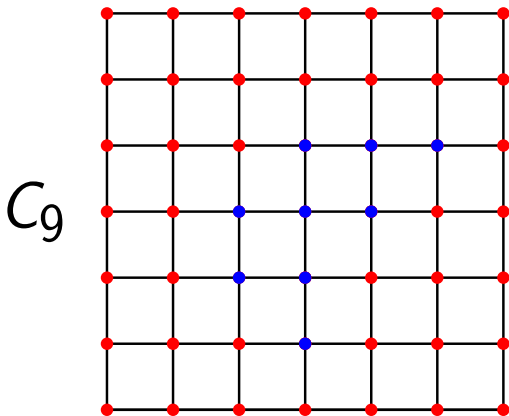
Eden exploration



Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).

Markovian formulation

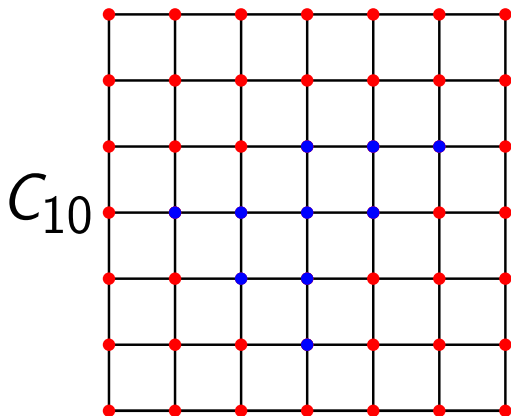
Eden exploration



Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).

Markovian formulation

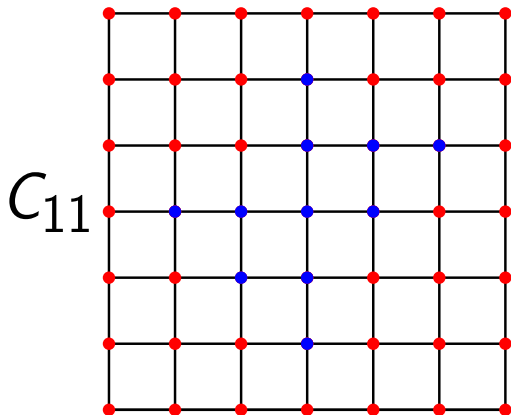
Eden exploration



Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).

Markovian formulation

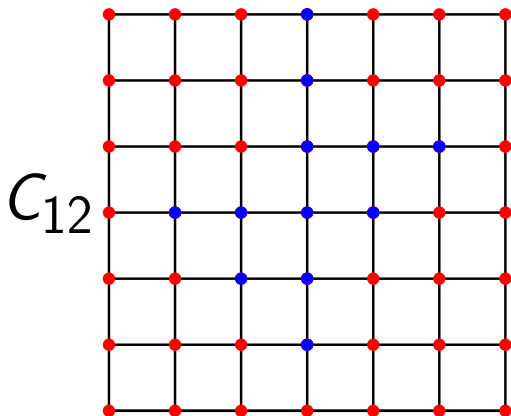
Eden exploration



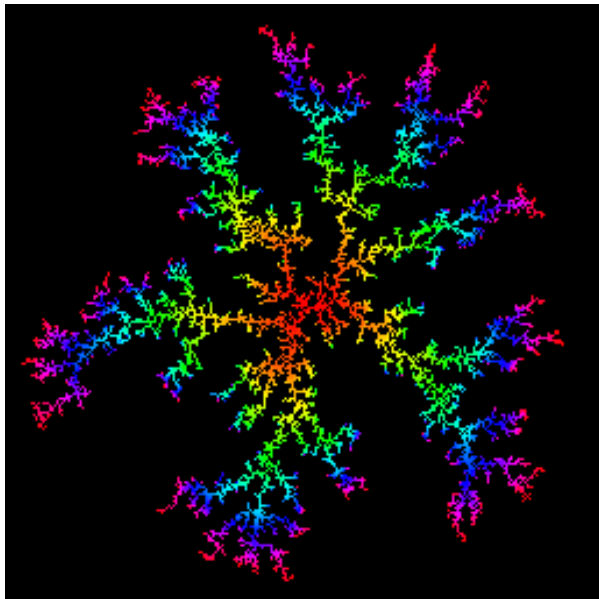
Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).

Markovian formulation

Eden exploration



Sample the cluster C_{n+1} from C_n by selecting an edge uniformly at random on ∂C_n , and then adding the vertex which is attached to it. **VARIANT:** Choose locations from harmonic measure (DLA) or harmonic measure to η power (η -DBM).



Euclidean Diffusion Limited Aggregation (DLA) introduced by Witten-Sander 1981.



DLA in nature: "A DLA cluster grown from a copper sulfate solution in an electrodeposition cell" (from Wikipedia)



DLA in nature: Magnese oxide patterns on the surface of a rock. (Halsey, Physics Today 2000)



DLA in nature: Magnese oxide patterns on the surface of a rock.

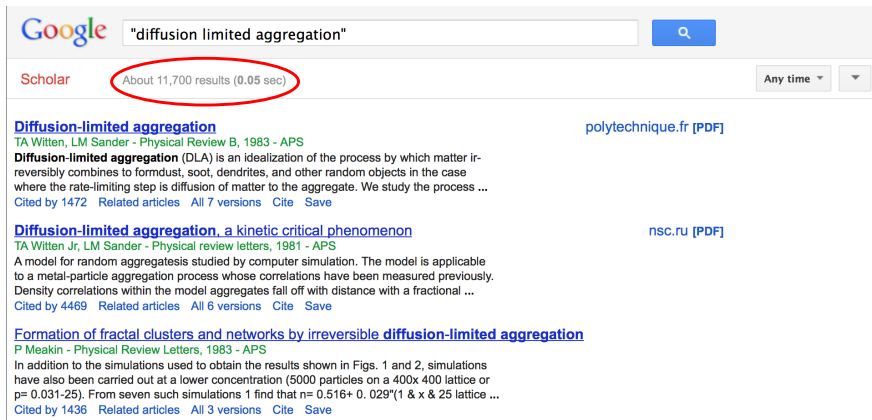


DLA in art: “High-voltage dielectric breakdown within a block of plexiglas” (from Wikipedia)

DLA in physics

Introduced by Witten and Sander in 1981 as a model for crystal growth. (Mineral deposits, Hele-Shaw flow, electrodeposition, lichen growth, lightning paths, coral, etc.)

An active area of research in physics for the last 33 years:



The image shows a Google Scholar search interface. The search bar contains the text "diffusion limited aggregation". Below the search bar, the results are displayed. The first result is "Diffusion-limited aggregation" by TA Witten and LM Sander, published in Physical Review B, 1983. The second result is "Diffusion-limited aggregation, a kinetic critical phenomenon" by TA Witten Jr and LM Sander, published in Physical review letters, 1981. The third result is "Formation of fractal clusters and networks by irreversible diffusion-limited aggregation" by P Meakin, published in Physical Review Letters, 1983. The search results are sorted by "Any time".

Google "diffusion limited aggregation"

Scholar About 11,700 results (0.05 sec) Any time ▾ ▾

Diffusion-limited aggregation polytechnique.fr [PDF]
TA Witten, LM Sander - *Physical Review B*, 1983 - APS
Diffusion-limited aggregation (DLA) is an idealization of the process by which matter irreversibly combines to form dust, soot, dendrites, and other random objects in the case where the rate-limiting step is diffusion of matter to the aggregate. We study the process ...
Cited by 1472 Related articles All 7 versions Cite Save

Diffusion-limited aggregation, a kinetic critical phenomenon nsc.ru [PDF]
TA Witten Jr, LM Sander - *Physical review letters*, 1981 - APS
A model for random aggregation studied by computer simulation. The model is applicable to a metal-particle aggregation process whose correlations have been measured previously. Density correlations within the model aggregates fall off with distance with a fractional ...
Cited by 4469 Related articles All 6 versions Cite Save

Formation of fractal clusters and networks by irreversible diffusion-limited aggregation
P Meakin - *Physical Review Letters*, 1983 - APS
In addition to the simulations used to obtain the results shown in Figs. 1 and 2, simulations have also been carried out at a lower concentration (5000 particles on a 400x 400 lattice or $p = 0.031-25$). From seven such simulations 1 find that $n = 0.516 + 0.029(1 + x + 25 \text{ lattice } \dots$
Cited by 1436 Related articles All 3 versions Cite Save

DLA in math?

DLA in math?

Not a lot of progress. (A related process called internal DLA is mathematically much more well understood.)

DLA in math?

Not a lot of progress. (A related process called internal DLA is mathematically much more well understood.)

Open questions

DLA in math?

Not a lot of progress. (A related process called internal DLA is mathematically much more well understood.)

Open questions

- ▶ Does DLA have a “scaling limit”?

DLA in math?

Not a lot of progress. (A related process called internal DLA is mathematically much more well understood.)

Open questions

- ▶ Does DLA have a “scaling limit”?
- ▶ Is the shape random at large scales?

DLA in math?

Not a lot of progress. (A related process called internal DLA is mathematically much more well understood.)

Open questions

- ▶ Does DLA have a “scaling limit”?
- ▶ Is the shape random at large scales?
- ▶ Does the macroscopic shape look like a tree?

DLA in math?

Not a lot of progress. (A related process called internal DLA is mathematically much more well understood.)

Open questions

- ▶ Does DLA have a “scaling limit”?
- ▶ Is the shape random at large scales?
- ▶ Does the macroscopic shape look like a tree?
- ▶ What is its asymptotic dimension? Simulation prediction: ≈ 1.71 on \mathbf{Z}^2

DLA in math?

Not a lot of progress. (A related process called internal DLA is mathematically much more well understood.)

Open questions

- ▶ Does DLA have a “scaling limit”?
- ▶ Is the shape random at large scales?
- ▶ Does the macroscopic shape look like a tree?
- ▶ What is its asymptotic dimension? Simulation prediction: ≈ 1.71 on \mathbf{Z}^2

Schramm 2006 ICM proceedings:

Given that the fractals produced by DLA are not conformally invariant, it is not too surprising that it is hard to faithfully model DLA using conformal maps. Harry Kesten [44] proved that the diameter of the planar DLA cluster after n steps grows asymptotically no faster than $n^{2/3}$, and this appears to be essentially the only theorem concerning two-dimensional DLA, though several very simplified variants of DLA have been successfully analysed.

DLA in math?

Not a lot of progress. (A related process called internal DLA is mathematically much more well understood.)

Open questions

- ▶ Does DLA have a “scaling limit”?
- ▶ Is the shape random at large scales?
- ▶ Does the macroscopic shape look like a tree?
- ▶ What is its asymptotic dimension? Simulation prediction: ≈ 1.71 on \mathbf{Z}^2

Schramm 2006 ICM proceedings:

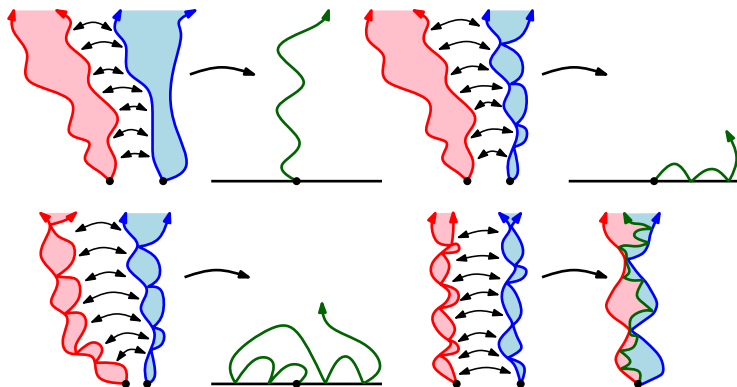
Given that the fractals produced by DLA are not conformally invariant, it is not too surprising that it is hard to faithfully model DLA using conformal maps. Harry Kesten [44] proved that the diameter of the planar DLA cluster after n steps grows asymptotically no faster than $n^{2/3}$, and this appears to be essentially the only theorem concerning two-dimensional DLA, though several very simplified variants of DLA have been successfully analysed.

What about DLA on random planar maps and Liouville quantum gravity surfaces?

Part II: DRAMA

WELDING RANDOM SURFACES

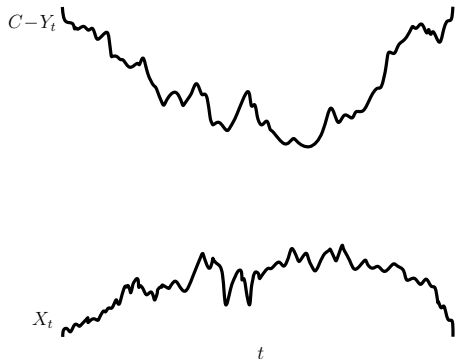
Can “weld” and “slice” special quantum surfaces called quantum wedges (with “weight” parameters indicating thickness) to obtain wedges (with other weights).



- ▶ Weight parameter $W = \gamma(\gamma + \frac{2}{\gamma} - \alpha)$ is additive under the welding operation.
- ▶ Interface between welding of independent wedges $\mathcal{W}_1, \mathcal{W}_2$ of weight W_1 and W_2 is an $\text{SLE}_\kappa(W_1 - 2; W_2 - 2)$ on combined surface.
- ▶ Glue **canonical random surfaces**, seam becomes **canonical random path**.

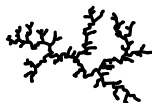
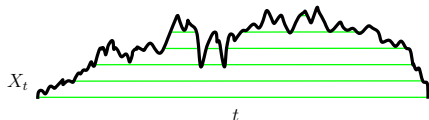
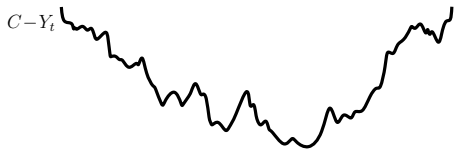
MATING RANDOM TREES

X, Y independent Brownian excursions on $[0, 1]$. Pick $C > 0$ large so that the graphs of X and $C - Y$ are disjoint.



MATING RANDOM TREES

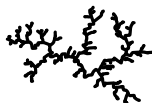
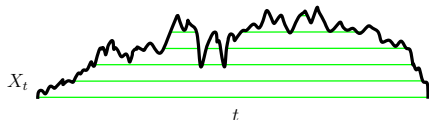
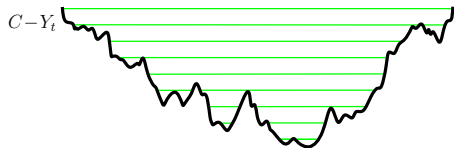
X, Y independent Brownian excursions on $[0, 1]$. Pick $C > 0$ large so that the graphs of X and $C - Y$ are disjoint.



- ▶ Identify points on the graph of X if they are connected by a **horizontal** line which is below the graph; yields a continuum random tree (CRT)

MATING RANDOM TREES

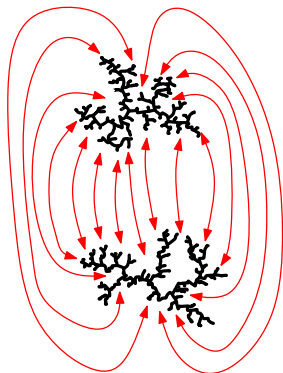
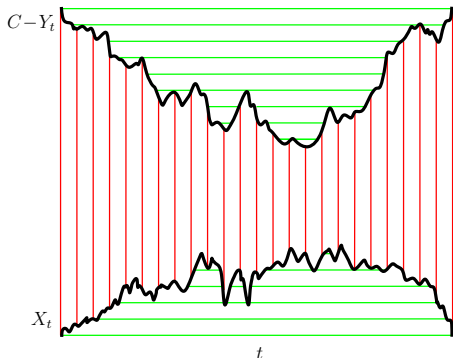
X, Y independent Brownian excursions on $[0, 1]$. Pick $C > 0$ large so that the graphs of X and $C - Y$ are disjoint.



- ▶ Identify points on the graph of X if they are connected by a **horizontal** line which is below the graph; yields a continuum random tree (CRT)
- ▶ Same for $C - Y_t$ yields an independent CRT

MATING RANDOM TREES

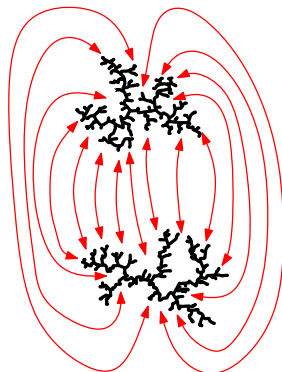
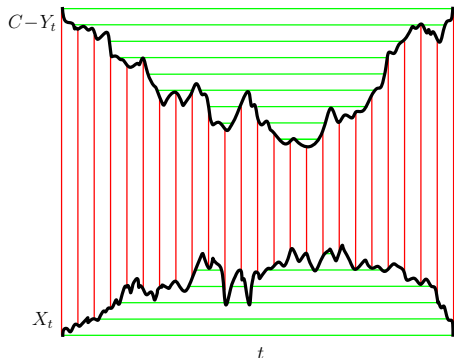
X, Y independent Brownian excursions on $[0, 1]$. Pick $C > 0$ large so that the graphs of X and $C - Y$ are disjoint.



- ▶ Identify points on the graph of X if they are connected by a **horizontal** line which is below the graph; yields a continuum random tree (CRT)
- ▶ Same for $C - Y_t$ yields an independent CRT
- ▶ Glue the CRTs together by declaring points on the **vertical** lines to be equivalent

MATING RANDOM TREES

X, Y independent Brownian excursions on $[0, 1]$. Pick $C > 0$ large so that the graphs of X and $C - Y$ are disjoint.

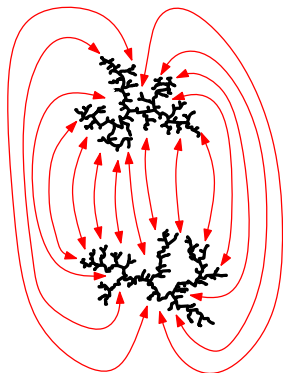
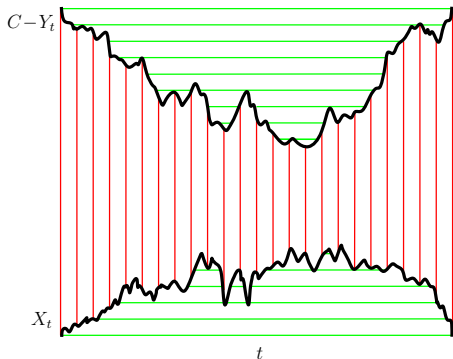


- ▶ Identify points on the graph of X if they are connected by a **horizontal** line which is below the graph; yields a continuum random tree (CRT)
- ▶ Same for $C - Y_t$ yields an independent CRT
- ▶ Glue the CRTs together by declaring points on the **vertical** lines to be equivalent

Q: What is the resulting structure?

MATING RANDOM TREES

X, Y independent Brownian excursions on $[0, 1]$. Pick $C > 0$ large so that the graphs of X and $C - Y$ are disjoint.

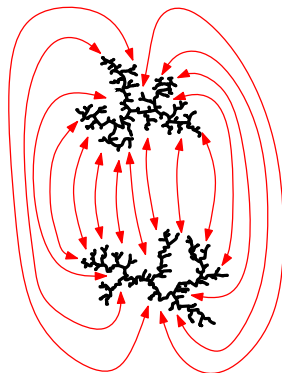
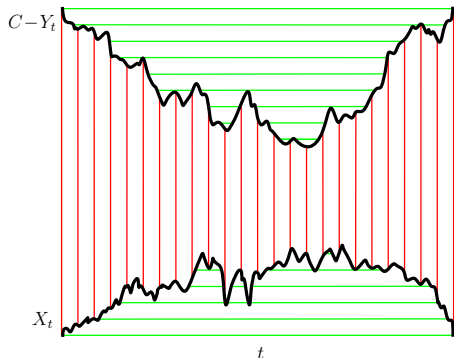


- ▶ Identify points on the graph of X if they are connected by a **horizontal** line which is below the graph; yields a continuum random tree (CRT)
- ▶ Same for $C - Y_t$ yields an independent CRT
- ▶ Glue the CRTs together by declaring points on the **vertical** lines to be equivalent

Q: What is the resulting structure? **A:** Sphere with a space-filling path.

MATING RANDOM TREES

X, Y independent Brownian excursions on $[0, 1]$. Pick $C > 0$ large so that the graphs of X and $C - Y$ are disjoint.



- ▶ Identify points on the graph of X if they are connected by a **horizontal** line which is below the graph; yields a continuum random tree (CRT)
- ▶ Same for $C - Y_t$ yields an independent CRT
- ▶ Glue the CRTs together by declaring points on the **vertical** lines to be equivalent

Q: What is the resulting structure? **A:** Sphere with a space-filling path. A **peanosphere**.

How to check this?

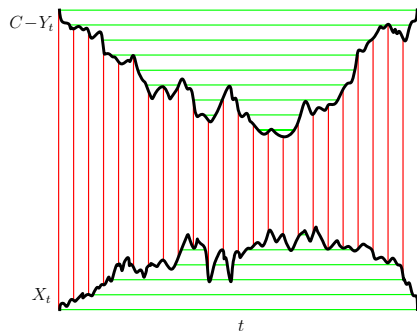
Theorem (Moore 1925)

Let \cong be any topologically closed equivalence relation on the sphere \mathbf{S}^2 . Assume that each equivalence class is connected and not equal to all of \mathbf{S}^2 . Then the quotient space \mathbf{S}^2 / \cong is homeomorphic to \mathbf{S}^2 if and only if no equivalence class separates the sphere into two or more connected components.

- ▶ An equivalence relation is topologically closed iff for any two sequences (x_n) and (y_n) with
 - ▶ $x_n \cong y_n$ for all n
 - ▶ $x_n \rightarrow x$ and $y_n \rightarrow y$
- ▶ we have that $x \cong y$.

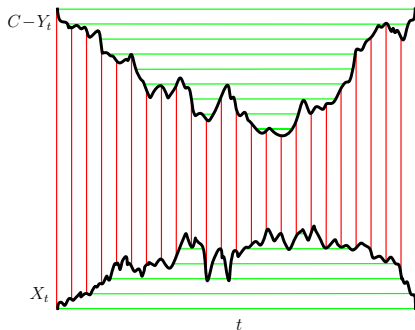
Constructing a sphere from a pair of trees

- ▶ X, Y ind. Brownian excursions on $[0, 1]$
- ▶ **Red/green** lines give an \cong -relation on \mathbf{S}^2



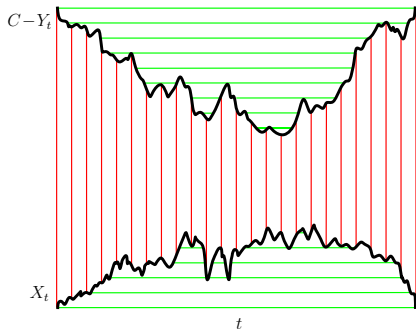
Constructing a sphere from a pair of trees

- ▶ X, Y ind. Brownian excursions on $[0, 1]$
- ▶ **Red/green** lines give an \cong -relation on \mathbf{S}^2
- ▶ Types of equivalence classes:



Constructing a sphere from a pair of trees

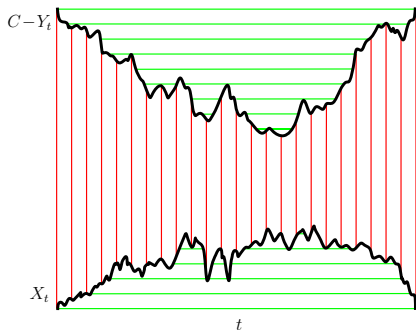
- ▶ X, Y ind. Brownian excursions on $[0, 1]$
- ▶ **Red/green** lines give an \cong -relation on \mathbf{S}^2
- ▶ Types of equivalence classes:
 1. Outer boundary of rectangle



Constructing a sphere from a pair of trees

- ▶ X, Y ind. Brownian excursions on $[0, 1]$
- ▶ **Red/green** lines give an \cong -relation on \mathbf{S}^2
- ▶ Types of equivalence classes:
 1. Outer boundary of rectangle
 2. **V** line which does not share an endpoint with a **H** line

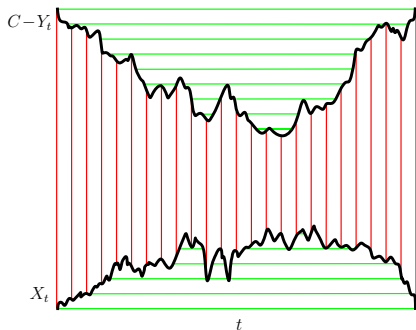
H = horizontal, **V** = vertical



Constructing a sphere from a pair of trees

- ▶ X, Y ind. Brownian excursions on $[0, 1]$
- ▶ **Red/green** lines give an \cong -relation on \mathbf{S}^2
- ▶ Types of equivalence classes:
 1. Outer boundary of rectangle
 2. **V** line which does not share an endpoint with a **H** line
 3. **H** line below X or above $C - Y$ with two **V** lines with common endpoint

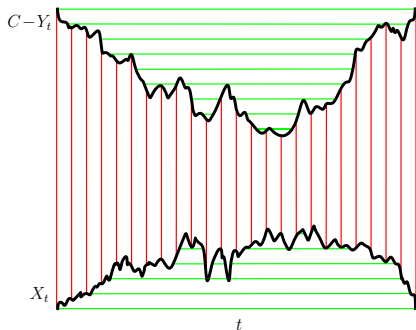
H = horizontal, **V** = vertical



Constructing a sphere from a pair of trees

- ▶ X, Y ind. Brownian excursions on $[0, 1]$
- ▶ **Red/green** lines give an \cong -relation on \mathbf{S}^2
- ▶ Types of equivalence classes:
 1. Outer boundary of rectangle
 2. **V** line which does not share an endpoint with a **H** line
 3. **H** line below X or above $C - Y$ with two **V** lines with common endpoint
 4. **H** line below X or above $C - Y$ with two **V** lines with common endpoint and a third **V** line hitting in the middle

H = horizontal, **V** = vertical



Constructing a sphere from a pair of trees

- ▶ X, Y ind. Brownian excursions on $[0, 1]$
- ▶ **Red/green** lines give an \cong -relation on \mathbf{S}^2
- ▶ Types of equivalence classes:
 1. Outer boundary of rectangle
 2. **V** line which does not share an endpoint with a **H** line
 3. **H** line below X or above $C - Y$ with two **V** lines with common endpoint
 4. **H** line below X or above $C - Y$ with two **V** lines with common endpoint and a third **V** line hitting in the middle
- ▶ \cong is topologically closed and does not separate \mathbf{S}^2 into two or more components, thus \mathbf{S}^2 / \cong is homeomorphic to \mathbf{S}^2

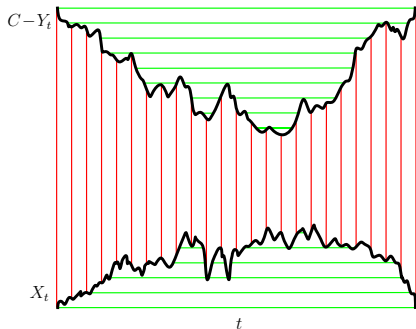
H = horizontal, **V** = vertical



Constructing a sphere from a pair of trees

- ▶ X, Y ind. Brownian excursions on $[0, 1]$
- ▶ **Red/green** lines give an \cong -relation on \mathbf{S}^2
- ▶ Types of equivalence classes:
 1. Outer boundary of rectangle
 2. **V** line which does not share an endpoint with a **H** line
 3. **H** line below X or above $C - Y$ with two **V** lines with common endpoint
 4. **H** line below X or above $C - Y$ with two **V** lines with common endpoint and a third **V** line hitting in the middle
- ▶ \cong is topologically closed and does not separate \mathbf{S}^2 into two or more components, thus \mathbf{S}^2 / \cong is homeomorphic to \mathbf{S}^2
- ▶ Following the **V** lines from left to right gives a space-filling path on \mathbf{S}^2 / \cong

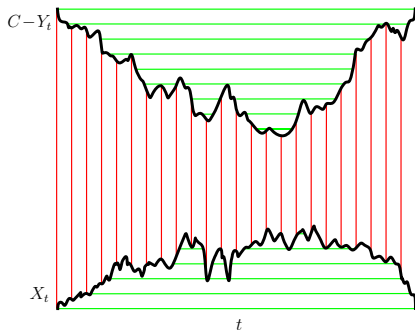
H = horizontal, **V** = vertical



Constructing a sphere from a pair of trees

- ▶ X, Y ind. Brownian excursions on $[0, 1]$
- ▶ **Red/green** lines give an \cong -relation on \mathbf{S}^2
- ▶ Types of equivalence classes:
 1. Outer boundary of rectangle
 2. **V** line which does not share an endpoint with a **H** line
 3. **H** line below X or above $C - Y$ with two **V** lines with common endpoint
 4. **H** line below X or above $C - Y$ with two **V** lines with common endpoint and a third **V** line hitting in the middle
- ▶ \cong is topologically closed and does not separate \mathbf{S}^2 into two or more components, thus \mathbf{S}^2 / \cong is homeomorphic to \mathbf{S}^2
- ▶ Following the **V** lines from left to right gives a space-filling path on \mathbf{S}^2 / \cong

H = horizontal, **V** = vertical

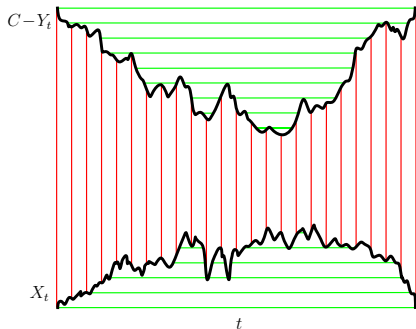


The sphere/space-filling path pair is a **peanoshere**

Constructing a sphere from a pair of trees

- ▶ X, Y ind. Brownian excursions on $[0, 1]$
- ▶ **Red/green** lines give an \cong -relation on \mathbf{S}^2
- ▶ Types of equivalence classes:
 1. Outer boundary of rectangle
 2. **V** line which does not share an endpoint with a **H** line
 3. **H** line below X or above $C - Y$ with two **V** lines with common endpoint
 4. **H** line below X or above $C - Y$ with two **V** lines with common endpoint and a third **V** line hitting in the middle
- ▶ \cong is topologically closed and does not separate \mathbf{S}^2 into two or more components, thus \mathbf{S}^2 / \cong is homeomorphic to \mathbf{S}^2
- ▶ Following the **V** lines from left to right gives a space-filling path on \mathbf{S}^2 / \cong

H = horizontal, **V** = vertical

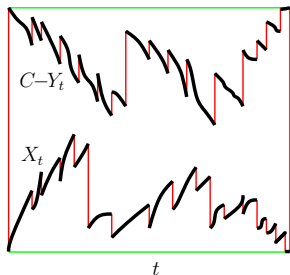


The sphere/space-filling path pair is a **peanoshere**

Peanosphere has **canonical embedding** in Euclidean sphere as **LQG, space-filling SLE**.

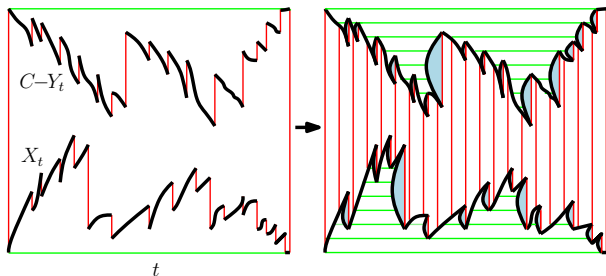
Gluing independent Lévy trees

Can view $\text{SLE}_{\kappa'}$ process, $\kappa' \in (4, 8)$ as a gluing of two $\frac{\kappa'}{4}$ -stable Lévy trees.



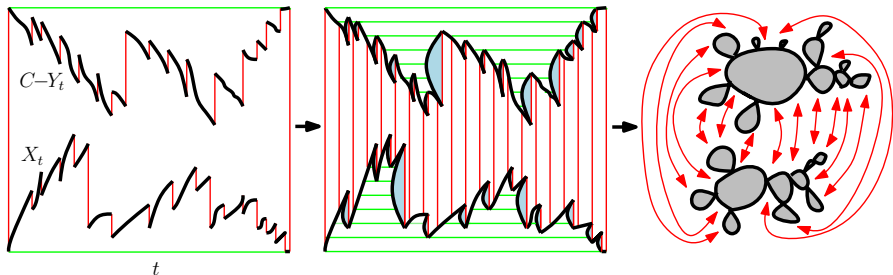
Gluing independent Lévy trees

Can view $\text{SLE}_{\kappa'}$ process, $\kappa' \in (4, 8)$ as a gluing of two $\frac{\kappa'}{4}$ -stable Lévy trees.



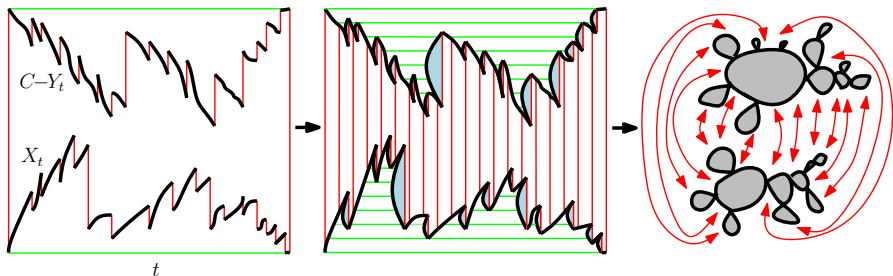
Gluing independent Lévy trees

Can view $\text{SLE}_{\kappa'}$ process, $\kappa' \in (4, 8)$ as a gluing of two $\frac{\kappa'}{4}$ -stable Lévy trees.



Gluing independent Lévy trees

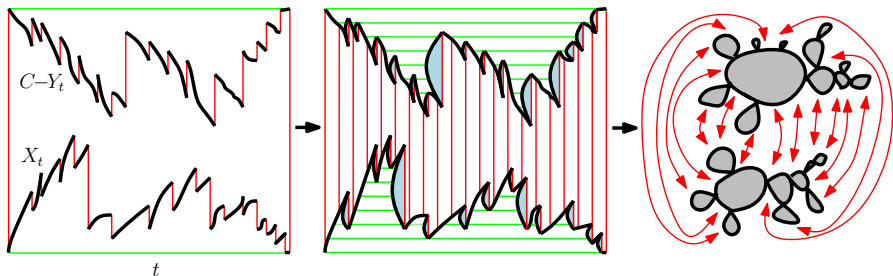
Can view $SLE_{\kappa'}$ process, $\kappa' \in (4, 8)$ as a gluing of two $\frac{\kappa'}{4}$ -stable Lévy trees.



- ▶ The two trees of quantum disks almost surely determine both the $SLE_{\kappa'}$ and the LQG surface on which it is drawn

Gluing independent Lévy trees

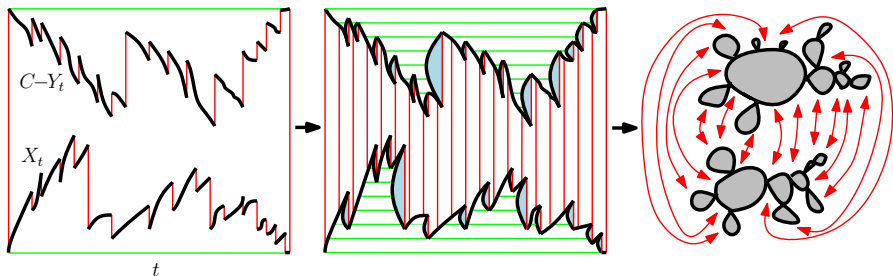
Can view $SLE_{\kappa'}$ process, $\kappa' \in (4, 8)$ as a gluing of two $\frac{\kappa'}{4}$ -stable Lévy trees.



- ▶ The two trees of quantum disks almost surely determine both the $SLE_{\kappa'}$ and the LQG surface on which it is drawn
- ▶ Can convert questions about $SLE_{\kappa'}$ into questions about $\frac{\kappa'}{4}$ -stable processes.

Gluing independent Lévy trees

Can view $\text{SLE}_{\kappa'}$ process, $\kappa' \in (4, 8)$ as a gluing of two $\frac{\kappa'}{4}$ -stable Lévy trees.



- ▶ The two trees of quantum disks almost surely determine both the $\text{SLE}_{\kappa'}$ and the LQG surface on which it is drawn
- ▶ Can convert questions about $\text{SLE}_{\kappa'}$ into questions about $\frac{\kappa'}{4}$ -stable processes.
- ▶ Scaling limit of “exploration path” on random planar map should be SLE_6 on a $\sqrt{8/3}$ -LQG. Using welding machinery, we can understand well the “bubbles” cut out by such an exploration process. We can understand conditional law of unexplored region given what we have seen.

RANDOM GROWTH ON RANDOM SURFACES

- ▶ Can we make sense of η -DBM on a γ -LQG? We have shown how to tile an LQG surface with dyadic squares of “about the same size” so we could run a DLA on this set of squares and try to take a fine mesh limit.

RANDOM GROWTH ON RANDOM SURFACES

- ▶ Can we make sense of η -DBM on a γ -LQG? We have shown how to tile an LQG surface with dyadic squares of “about the same size” so we could run a DLA on this set of squares and try to take a fine mesh limit.
- ▶ Or we could try η -DBM on corresponding RPM, which one would expect to behave similarly....

RANDOM GROWTH ON RANDOM SURFACES

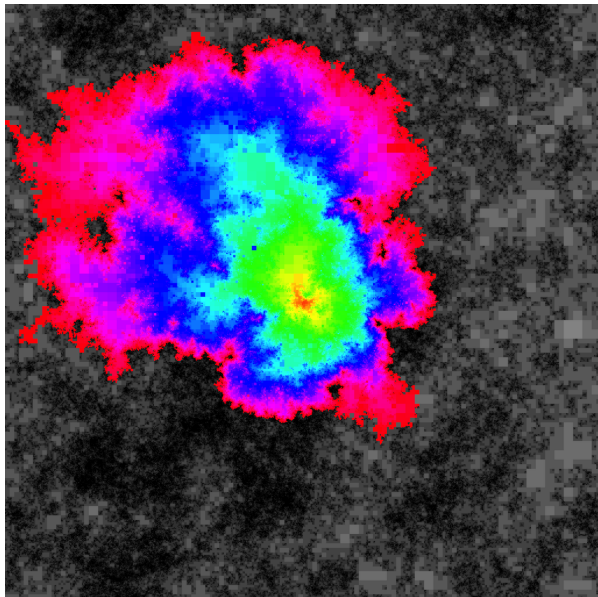
- ▶ Can we make sense of η -DBM on a γ -LQG? We have shown how to tile an LQG surface with dyadic squares of “about the same size” so we could run a DLA on this set of squares and try to take a fine mesh limit.
- ▶ Or we could try η -DBM on corresponding RPM, which one would expect to behave similarly....
- ▶ **Question:** Are there coral reefs, snowflakes, lichen, crystals, plants, lightning bolts, etc. whose growth rates are affected by a random medium (something like LQG)? The simulations look similar but have a bit more personality when γ is larger (as we will see). They look like Chinese dragons.

RANDOM GROWTH ON RANDOM SURFACES

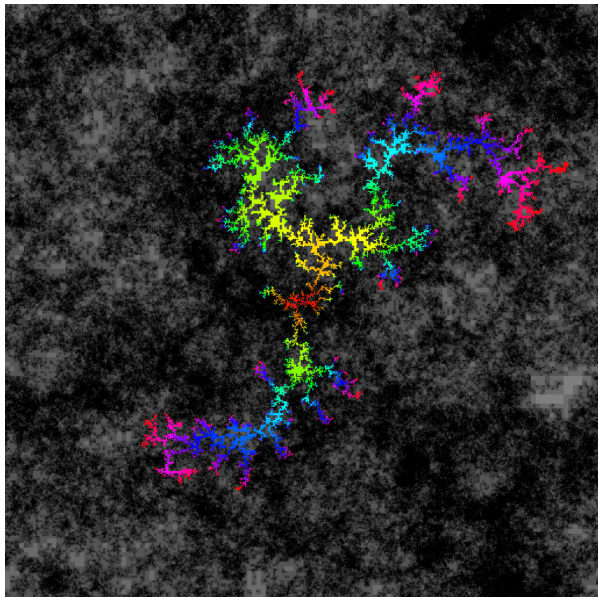
- ▶ Can we make sense of η -DBM on a γ -LQG? We have shown how to tile an LQG surface with dyadic squares of “about the same size” so we could run a DLA on this set of squares and try to take a fine mesh limit.
- ▶ Or we could try η -DBM on corresponding RPM, which one would expect to behave similarly....
- ▶ **Question:** Are there coral reefs, snowflakes, lichen, crystals, plants, lightning bolts, etc. whose growth rates are affected by a random medium (something like LQG)? The simulations look similar but have a bit more personality when γ is larger (as we will see). They look like Chinese dragons.
- ▶ We will ultimately want to construct a candidate for the scaling limit, which we will call (for reasons explained later) **quantum Loewner evolution:** $\text{QLE}(\gamma^2, \eta)$.

RANDOM GROWTH ON RANDOM SURFACES

- ▶ Can we make sense of η -DBM on a γ -LQG? We have shown how to tile an LQG surface with dyadic squares of “about the same size” so we could run a DLA on this set of squares and try to take a fine mesh limit.
- ▶ Or we could try η -DBM on corresponding RPM, which one would expect to behave similarly....
- ▶ **Question:** Are there coral reefs, snowflakes, lichen, crystals, plants, lightning bolts, etc. whose growth rates are affected by a random medium (something like LQG)? The simulations look similar but have a bit more personality when γ is larger (as we will see). They look like Chinese dragons.
- ▶ We will ultimately want to construct a candidate for the scaling limit, which we will call (for reasons explained later) **quantum Loewner evolution:** $\text{QLE}(\gamma^2, \eta)$.
- ▶ But first let's look at some computer generated images (and some animations), starting with an Eden exploration.



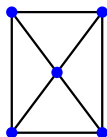
Eden model on $\sqrt{8/3}$ -LQG



DLA on a $\sqrt{2}$ -LQG

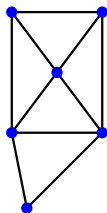
Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .



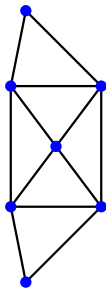
Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .



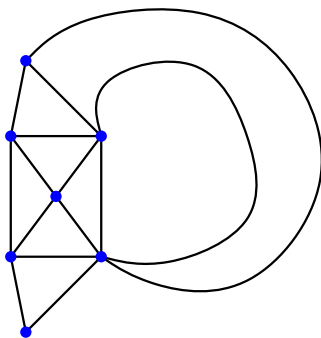
Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .



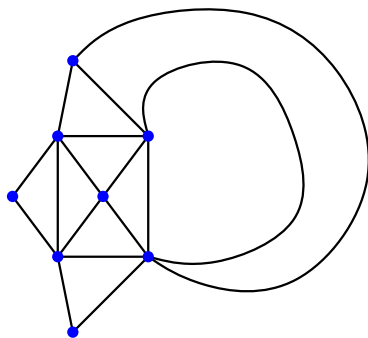
Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .



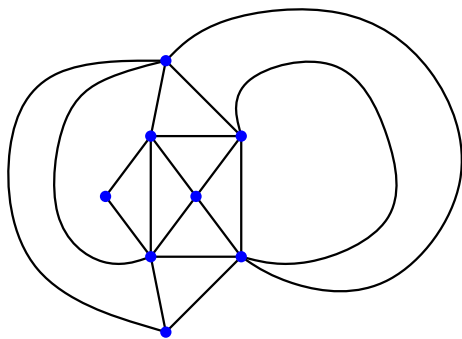
Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .



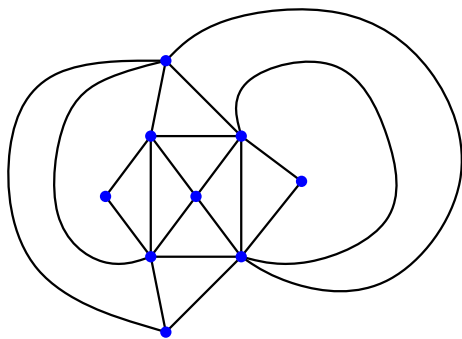
Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .



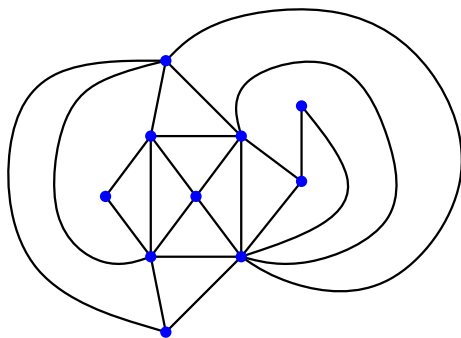
Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .



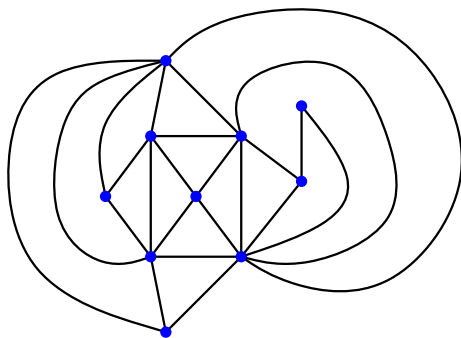
Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .



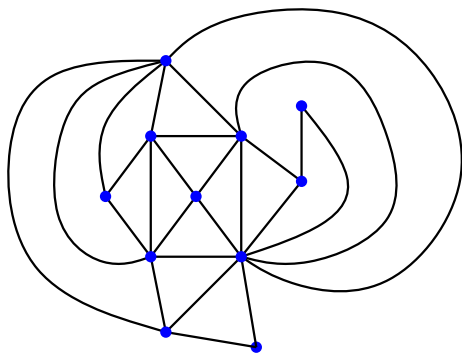
Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .



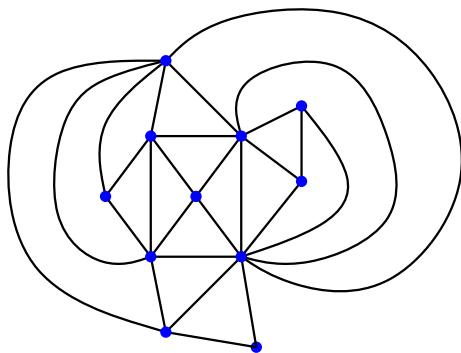
Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .



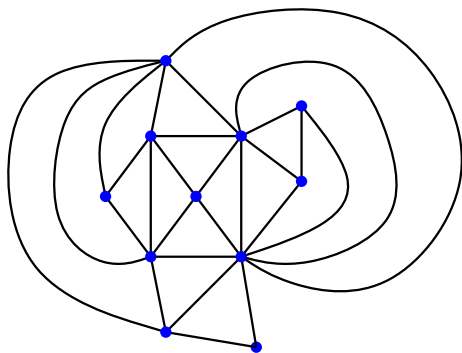
Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .



Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .

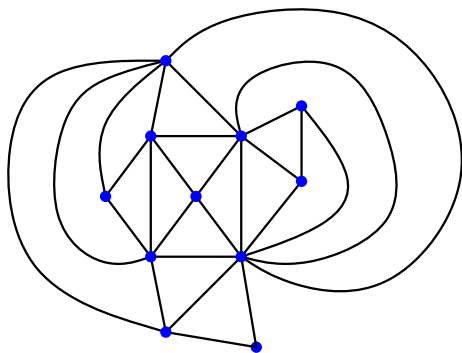


Important observations:

- ▶ Conditional law of map given ball at time n only depends on the boundary lengths of the outside components.

Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .

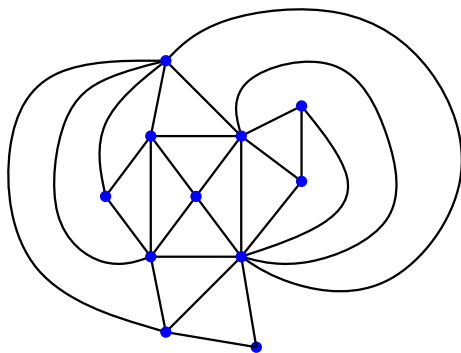


Important observations:

- ▶ Conditional law of map given ball at time n only depends on the boundary lengths of the outside components. *Exploration respects the Markovian structure of the map.*

Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .

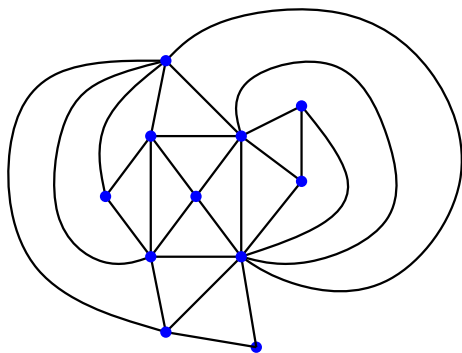


Important observations:

- ▶ Conditional law of map given ball at time n only depends on the boundary lengths of the outside components. *Exploration respects the Markovian structure of the map.*
- ▶ If we work on an “infinite” planar map, the conditional law of the map in the unbounded component only depends on the boundary length

Eden model on planar map

- ▶ Random planar map, random vertex x . Perform FPP from x .



Important observations:

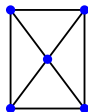
- ▶ Conditional law of map given ball at time n only depends on the boundary lengths of the outside components. *Exploration respects the Markovian structure of the map.*
- ▶ If we work on an “infinite” planar map, the conditional law of the map in the unbounded component only depends on the boundary length

Belief: Isotropic enough so that at large scales this is close to a ball in the graph metric

First passage percolation on random planar maps III

Variant:

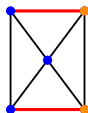
- ▶ Pick two **edges** on outer boundary of cluster



First passage percolation on random planar maps III

Variant:

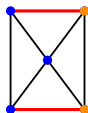
- ▶ Pick two **edges** on outer boundary of cluster
- ▶ Color vertices between edges blue and yellow



First passage percolation on random planar maps III

Variant:

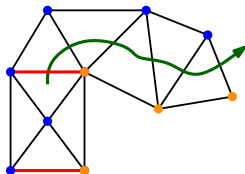
- ▶ Pick two **edges** on outer boundary of cluster
- ▶ Color vertices between edges blue and yellow
- ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$



First passage percolation on random planar maps III

Variant:

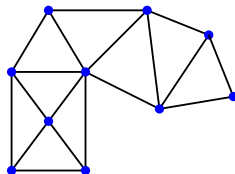
- ▶ Pick two **edges** on outer boundary of cluster
- ▶ Color vertices between edges blue and yellow
- ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$
- ▶ Explore percolation (blue/yellow) interface



First passage percolation on random planar maps III

Variant:

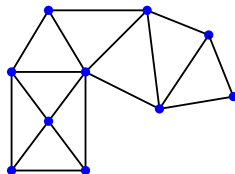
- ▶ Pick two **edges** on outer boundary of cluster
- ▶ Color vertices between edges blue and yellow
- ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$
- ▶ Explore percolation (blue/yellow) interface
- ▶ Forget colors



First passage percolation on random planar maps III

Variant:

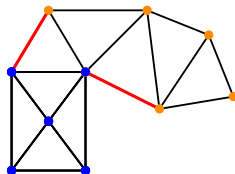
- ▶ Pick two **edges** on outer boundary of cluster
- ▶ Color vertices between edges blue and yellow
- ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$
- ▶ Explore percolation (blue/yellow) interface
- ▶ Forget colors
- ▶ Repeat



First passage percolation on random planar maps III

Variant:

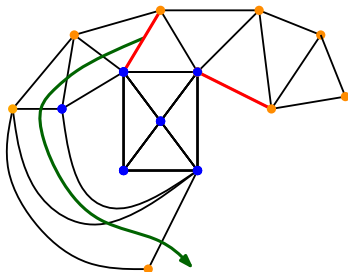
- ▶ Pick two **edges** on outer boundary of cluster
- ▶ Color vertices between edges blue and yellow
- ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$
- ▶ Explore percolation (blue/yellow) interface
- ▶ Forget colors
- ▶ Repeat



First passage percolation on random planar maps III

Variant:

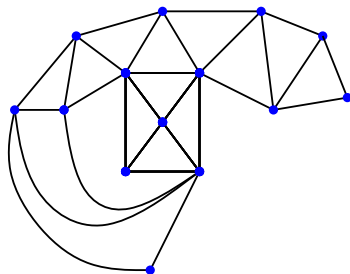
- ▶ Pick two **edges** on outer boundary of cluster
- ▶ Color vertices between edges blue and yellow
- ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$
- ▶ Explore percolation (blue/yellow) interface
- ▶ Forget colors
- ▶ Repeat



First passage percolation on random planar maps III

Variant:

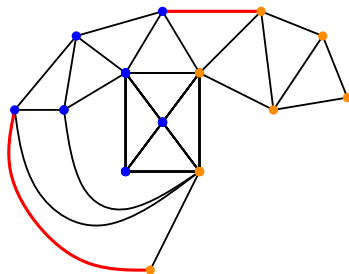
- ▶ Pick two **edges** on outer boundary of cluster
- ▶ Color vertices between edges blue and yellow
- ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$
- ▶ Explore percolation (blue/yellow) interface
- ▶ Forget colors
- ▶ Repeat



First passage percolation on random planar maps III

Variant:

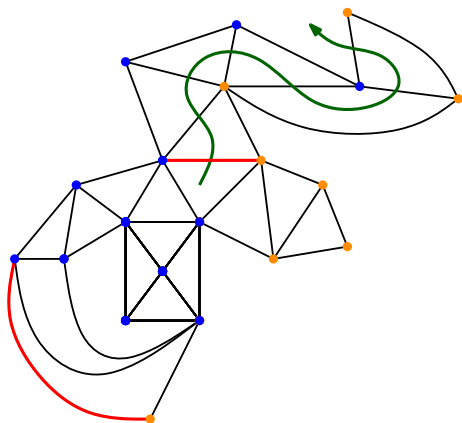
- ▶ Pick two **edges** on outer boundary of cluster
- ▶ Color vertices between edges blue and yellow
- ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$
- ▶ Explore percolation (blue/yellow) interface
- ▶ Forget colors
- ▶ Repeat



First passage percolation on random planar maps III

Variant:

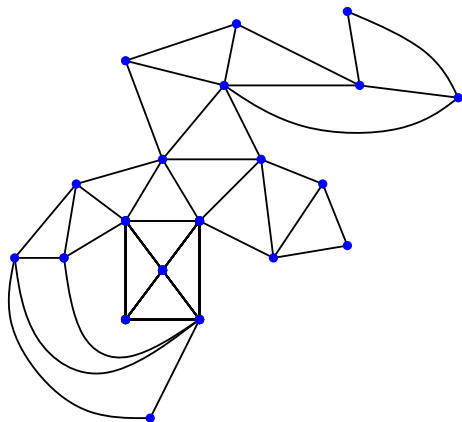
- ▶ Pick two **edges** on outer boundary of cluster
- ▶ Color vertices between edges blue and yellow
- ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$
- ▶ Explore percolation (blue/yellow) interface
- ▶ Forget colors
- ▶ Repeat



First passage percolation on random planar maps III

Variant:

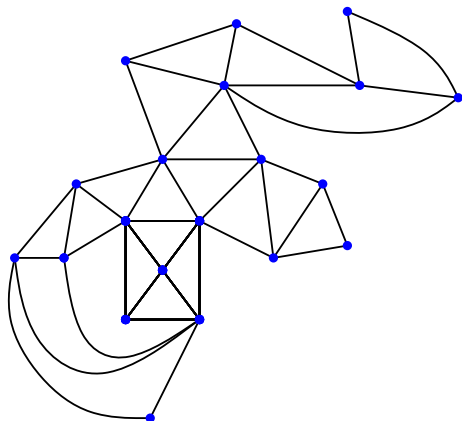
- ▶ Pick two **edges** on outer boundary of cluster
- ▶ Color vertices between edges blue and yellow
- ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$
- ▶ Explore percolation (blue/yellow) interface
- ▶ Forget colors
- ▶ Repeat



First passage percolation on random planar maps III

Variant:

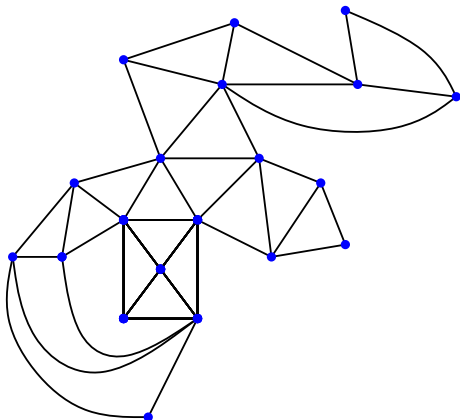
- ▶ Pick two **edges** on outer boundary of cluster
 - ▶ Color vertices between edges blue and yellow
 - ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$
 - ▶ Explore percolation (blue/yellow) interface
 - ▶ Forget colors
 - ▶ Repeat
- ▶ *This exploration also respects the Markovian structure of the map.*



First passage percolation on random planar maps III

Variant:

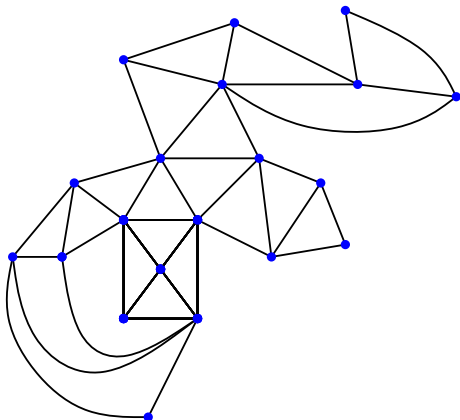
- ▶ Pick two **edges** on outer boundary of cluster
 - ▶ Color vertices between edges blue and yellow
 - ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$
 - ▶ Explore percolation (blue/yellow) interface
 - ▶ Forget colors
 - ▶ Repeat
- ▶ *This exploration also respects the Markovian structure of the map.*
- ▶ If we work on an “infinite” planar map, the conditional law of the map in the unbounded component only depends on the boundary length.



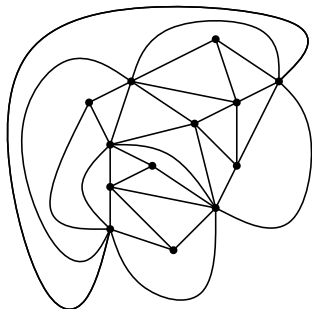
First passage percolation on random planar maps III

Variant:

- ▶ Pick two **edges** on outer boundary of cluster
 - ▶ Color vertices between edges blue and yellow
 - ▶ Color vertices on rest of map blue or yellow with prob. $\frac{1}{2}$
 - ▶ Explore percolation (blue/yellow) interface
 - ▶ Forget colors
 - ▶ Repeat
- ▶ *This exploration also respects the Markovian structure of the map.*
- ▶ If we work on an “infinite” planar map, the conditional law of the map in the unbounded component only depends on the boundary length.
- ▶ Expect that at large scales this growth process looks the same as FPP, hence the same as the graph metric ball

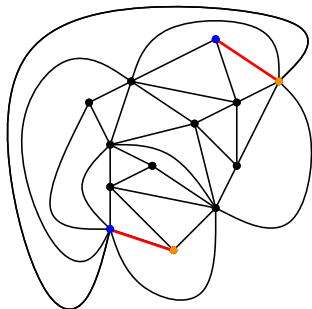


Continuum limit ansatz



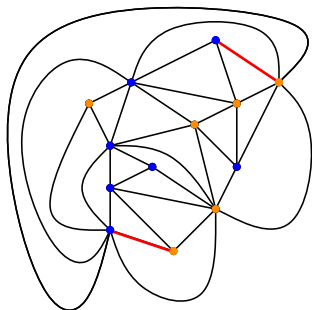
- ▶ Sample a random planar map

Continuum limit ansatz



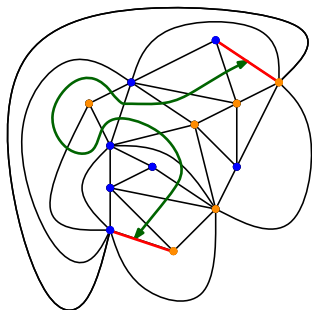
- ▶ Sample a random planar map and two edges uniformly at random

Continuum limit ansatz



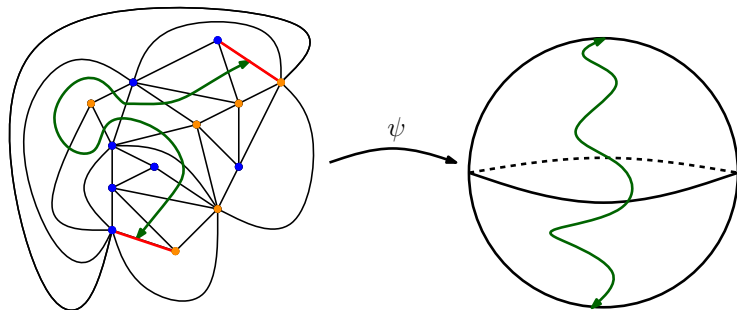
- ▶ Sample a random planar map and two edges uniformly at random
- ▶ Color vertices blue/yellow with probability $1/2$

Continuum limit ansatz



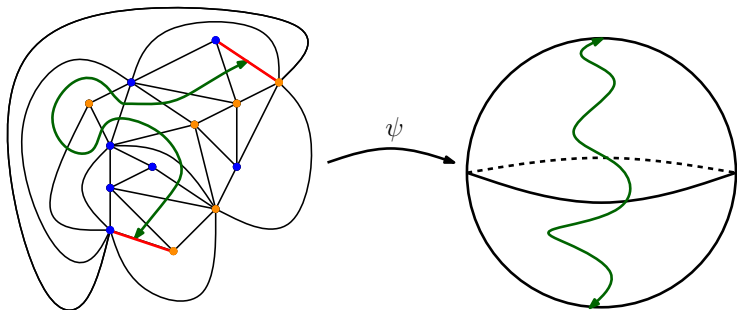
- ▶ Sample a random planar map and two edges uniformly at random
- ▶ Color vertices blue/yellow with probability $1/2$ and draw percolation interface

Continuum limit ansatz



- ▶ Sample a random planar map and two edges uniformly at random
- ▶ Color vertices blue/yellow with probability $1/2$ and draw percolation interface
- ▶ Conformally map to the sphere

Continuum limit ansatz



- ▶ Sample a random planar map and two edges uniformly at random
- ▶ Color vertices blue/yellow with probability $1/2$ and draw percolation interface
- ▶ Conformally map to the sphere

Ansatz Image of random map converges to a $\sqrt{8/3}$ -LQG surface and the image of the interface converges to an independent SLE_6 .

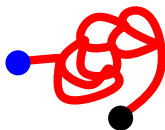
Continuum analog of first passage percolation on LQG

- ▶ Start off with $\sqrt{8/3}$ -LQG surface
- ▶ Fix $\delta > 0$ small and a starting point x



Continuum analog of first passage percolation on LQG

- ▶ Start off with $\sqrt{8/3}$ -LQG surface
- ▶ Fix $\delta > 0$ small and a starting point x
- ▶ Draw δ units of SLE_6



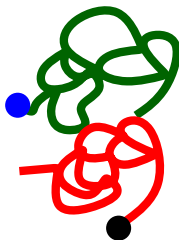
Continuum analog of first passage percolation on LQG

- ▶ Start off with $\sqrt{8/3}$ -LQG surface
- ▶ Fix $\delta > 0$ small and a starting point x
- ▶ Draw δ units of SLE_6
- ▶ Resample the tip according to boundary length



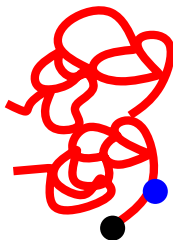
Continuum analog of first passage percolation on LQG

- ▶ Start off with $\sqrt{8/3}$ -LQG surface
- ▶ Fix $\delta > 0$ small and a starting point x
- ▶ Draw δ units of SLE_6
- ▶ Resample the tip according to boundary length
- ▶ Repeat



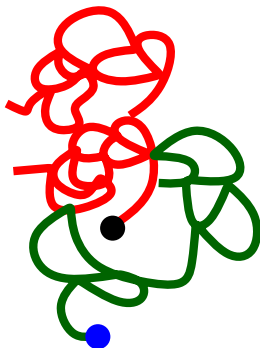
Continuum analog of first passage percolation on LQG

- ▶ Start off with $\sqrt{8/3}$ -LQG surface
- ▶ Fix $\delta > 0$ small and a starting point x
- ▶ Draw δ units of SLE_6
- ▶ Resample the tip according to boundary length
- ▶ Repeat



Continuum analog of first passage percolation on LQG

- ▶ Start off with $\sqrt{8/3}$ -LQG surface
- ▶ Fix $\delta > 0$ small and a starting point x
- ▶ Draw δ units of SLE_6
- ▶ Resample the tip according to boundary length
- ▶ Repeat



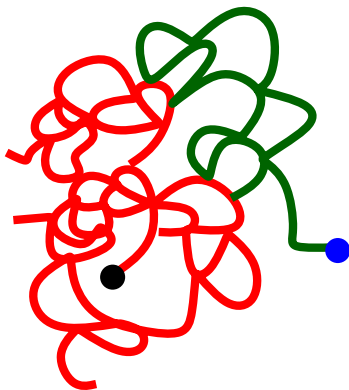
Continuum analog of first passage percolation on LQG

- ▶ Start off with $\sqrt{8/3}$ -LQG surface
- ▶ Fix $\delta > 0$ small and a starting point x
- ▶ Draw δ units of SLE_6
- ▶ Resample the tip according to boundary length
- ▶ Repeat



Continuum analog of first passage percolation on LQG

- ▶ Start off with $\sqrt{8/3}$ -LQG surface
- ▶ Fix $\delta > 0$ small and a starting point x
- ▶ Draw δ units of SLE_6
- ▶ Resample the tip according to boundary length
- ▶ Repeat



Continuum analog of first passage percolation on LQG

- ▶ Start off with $\sqrt{8/3}$ -LQG surface
- ▶ Fix $\delta > 0$ small and a starting point x
- ▶ Draw δ units of SLE_6
- ▶ Resample the tip according to boundary length
- ▶ Repeat



Continuum analog of first passage percolation on LQG

- ▶ Start off with $\sqrt{8/3}$ -LQG surface
- ▶ Fix $\delta > 0$ small and a starting point x
- ▶ Draw δ units of SLE_6
- ▶ Resample the tip according to boundary length
- ▶ Repeat
- ▶ Know the conditional law of the LQG surface at each stage, using exploration results



Continuum analog of first passage percolation on LQG

- ▶ Start off with $\sqrt{8/3}$ -LQG surface
- ▶ Fix $\delta > 0$ small and a starting point x
- ▶ Draw δ units of SLE_6
- ▶ Resample the tip according to boundary length
- ▶ Repeat
- ▶ Know the conditional law of the LQG surface at each stage, using exploration results



$QLE(8/3, 0)$ is the limit as $\delta \rightarrow 0$ of this growth process. It is described in terms of a radial Loewner evolution which is driven by a measure valued diffusion.

Continuum analog of first passage percolation on LQG

- ▶ Start off with $\sqrt{8/3}$ -LQG surface
- ▶ Fix $\delta > 0$ small and a starting point x
- ▶ Draw δ units of SLE_6
- ▶ Resample the tip according to boundary length
- ▶ Repeat
- ▶ Know the conditional law of the LQG surface at each stage, using exploration results



$QLE(8/3, 0)$ is the limit as $\delta \rightarrow 0$ of this growth process. It is described in terms of a radial Loewner evolution which is driven by a measure valued diffusion.

$QLE(8/3, 0)$ is SLE_6 with **tip re-randomization**. It can be understood as a “reshuffling” of the exploration procedure associated to the peanosphere.

What is $\text{QLE}(\gamma^2, \eta)$?

$\text{QLE}(8/3, 0)$ is a member of a two-parameter family of processes called $\text{QLE}(\gamma^2, \eta)$

- ▶ γ is the type of LQG surface on which the process grows
- ▶ η determines the manner in which it grows

What is $\text{QLE}(\gamma^2, \eta)$?

$\text{QLE}(8/3, 0)$ is a member of a two-parameter family of processes called $\text{QLE}(\gamma^2, \eta)$

- ▶ γ is the type of LQG surface on which the process grows
- ▶ η determines the manner in which it grows

Let μ_{HARM} (resp. μ_{LEN}) be harmonic (resp. length) measure on a γ -LQG surface. The rate of growth (i.e., rate at which microscopic particles are added) is proportional to

$$\left(\frac{d\mu_{\text{HARM}}}{d\mu_{\text{LEN}}} \right)^\eta d\mu_{\text{LEN}}.$$

What is $\text{QLE}(\gamma^2, \eta)$?

$\text{QLE}(8/3, 0)$ is a member of a two-parameter family of processes called $\text{QLE}(\gamma^2, \eta)$

- ▶ γ is the type of LQG surface on which the process grows
- ▶ η determines the manner in which it grows

Let μ_{HARM} (resp. μ_{LEN}) be harmonic (resp. length) measure on a γ -LQG surface. The rate of growth (i.e., rate at which microscopic particles are added) is proportional to

$$\left(\frac{d\mu_{\text{HARM}}}{d\mu_{\text{LEN}}} \right)^\eta d\mu_{\text{LEN}}.$$

- ▶ **First passage percolation:** $\eta = 0$

What is $\text{QLE}(\gamma^2, \eta)$?

$\text{QLE}(8/3, 0)$ is a member of a two-parameter family of processes called $\text{QLE}(\gamma^2, \eta)$

- ▶ γ is the type of LQG surface on which the process grows
- ▶ η determines the manner in which it grows

Let μ_{HARM} (resp. μ_{LEN}) be harmonic (resp. length) measure on a γ -LQG surface. The rate of growth (i.e., rate at which microscopic particles are added) is proportional to

$$\left(\frac{d\mu_{\text{HARM}}}{d\mu_{\text{LEN}}} \right)^\eta d\mu_{\text{LEN}}.$$

- ▶ **First passage percolation:** $\eta = 0$
- ▶ **Diffusion limited aggregation:** $\eta = 1$

What is $\text{QLE}(\gamma^2, \eta)$?

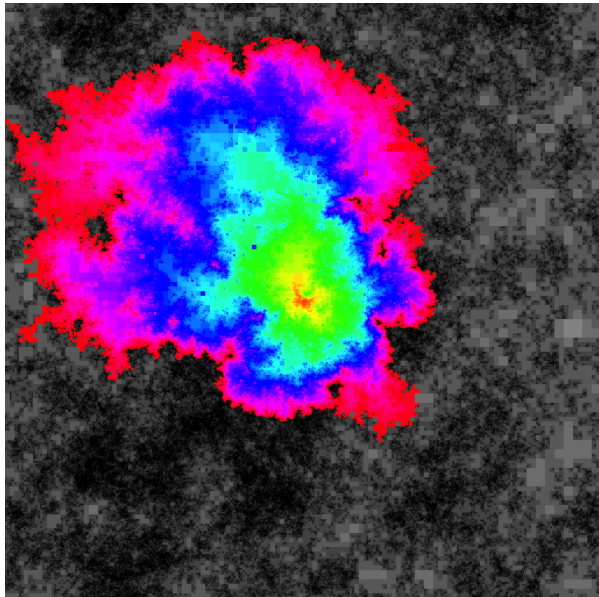
$\text{QLE}(8/3, 0)$ is a member of a two-parameter family of processes called $\text{QLE}(\gamma^2, \eta)$

- ▶ γ is the type of LQG surface on which the process grows
- ▶ η determines the manner in which it grows

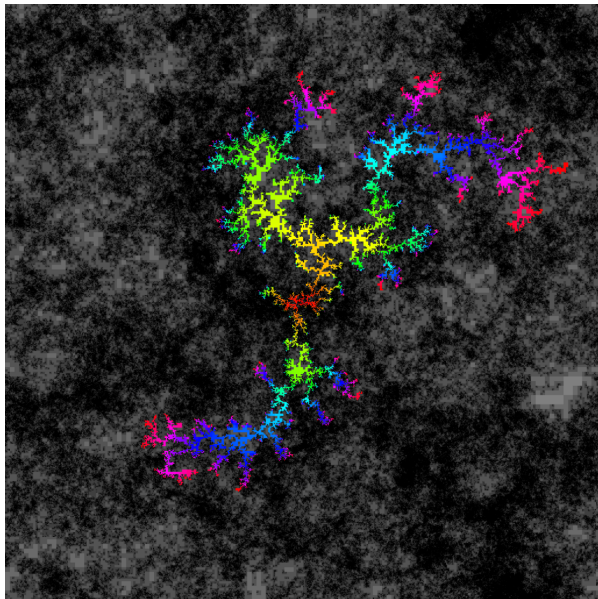
Let μ_{HARM} (resp. μ_{LEN}) be harmonic (resp. length) measure on a γ -LQG surface. The rate of growth (i.e., rate at which microscopic particles are added) is proportional to

$$\left(\frac{d\mu_{\text{HARM}}}{d\mu_{\text{LEN}}} \right)^\eta d\mu_{\text{LEN}}.$$

- ▶ **First passage percolation:** $\eta = 0$
- ▶ **Diffusion limited aggregation:** $\eta = 1$
- ▶ **η -dielectric breakdown model:** general values of η

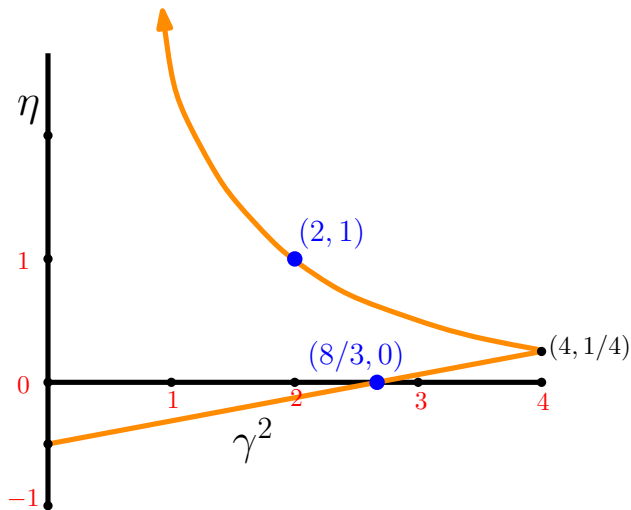


Discrete approximation of $QLE(8/3, 0)$. Metric ball on a $\sqrt{8/3}$ -LQG



Discrete approximation of $QLE(2, 1)$. DLA on a $\sqrt{2}$ -LQG

QLE(γ^2, η) processes we can construct



Each of the QLE(γ^2, η) processes with (γ^2, η) on the orange curves is built from an SLE $_{\kappa}$ process using tip re-randomization.

Results

What we can do:

- ▶ Existence of $\text{QLE}(\gamma^2, \eta)$ on the orange curves as a Markovian exploration of a γ -LQG surface.
- ▶ Derive an SPDE which the measure valued diffusion satisfies
- ▶ Continuity of the outer boundary of the growth at a given time

Results

What we can do:

- ▶ Existence of $\text{QLE}(\gamma^2, \eta)$ on the orange curves as a Markovian exploration of a γ -LQG surface.
- ▶ Derive an SPDE which the measure valued diffusion satisfies
- ▶ Continuity of the outer boundary of the growth at a given time

Work in progress:

- ▶ Results on phases for sample path behavior: which QLEs are trees, have holes, and fill space (joint also with Ewain Gwynne and Xin Sun)
- ▶ $\text{QLE}(8/3, 0)$ endows $\sqrt{8/3}$ -LQG with a distance function
- ▶ This metric space is isometric to the Brownian map: $\text{LQG} = \text{TBM}$

Results

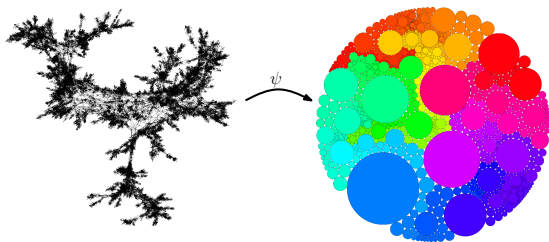
What we can do:

- ▶ Existence of $\text{QLE}(\gamma^2, \eta)$ on the orange curves as a Markovian exploration of a γ -LQG surface.
- ▶ Derive an SPDE which the measure valued diffusion satisfies
- ▶ Continuity of the outer boundary of the growth at a given time

Work in progress:

- ▶ Results on phases for sample path behavior: which QLEs are trees, have holes, and fill space (joint also with Ewain Gwynne and Xin Sun)
- ▶ $\text{QLE}(8/3, 0)$ endows $\sqrt{8/3}$ -LQG with a distance function
- ▶ This metric space is isometric to the Brownian map: $\text{LQG} = \text{TBM}$

What we would like to do: construct and study $\text{QLE}(\gamma^2, \eta)$ for (γ^2, η) pairs off the orange curves



Thanks!