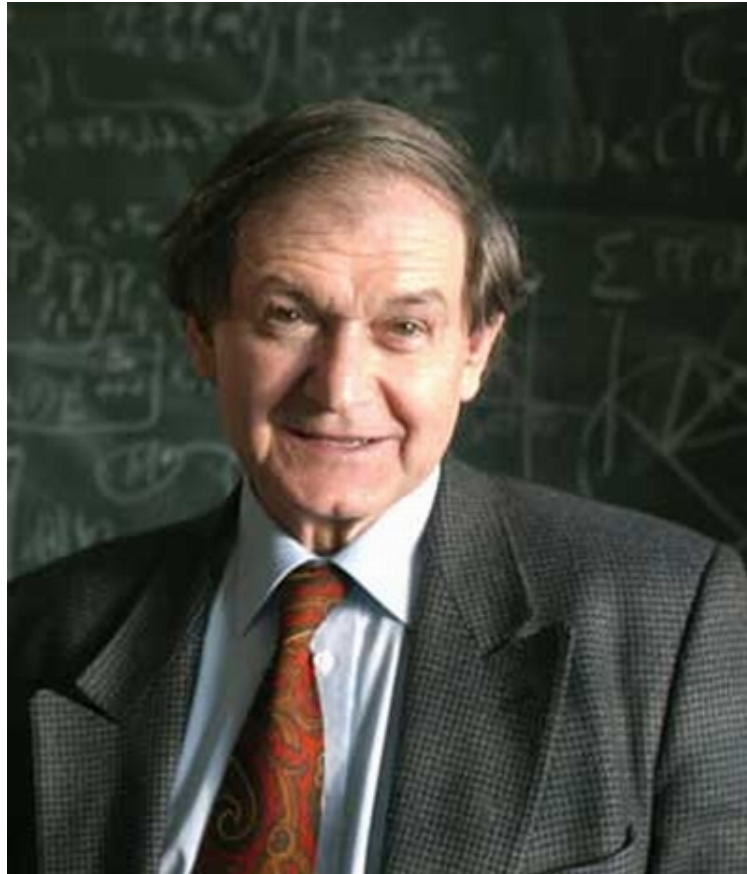


The Einstein-Weyl Equations,
Scattering Maps, and
Holomorphic Disks

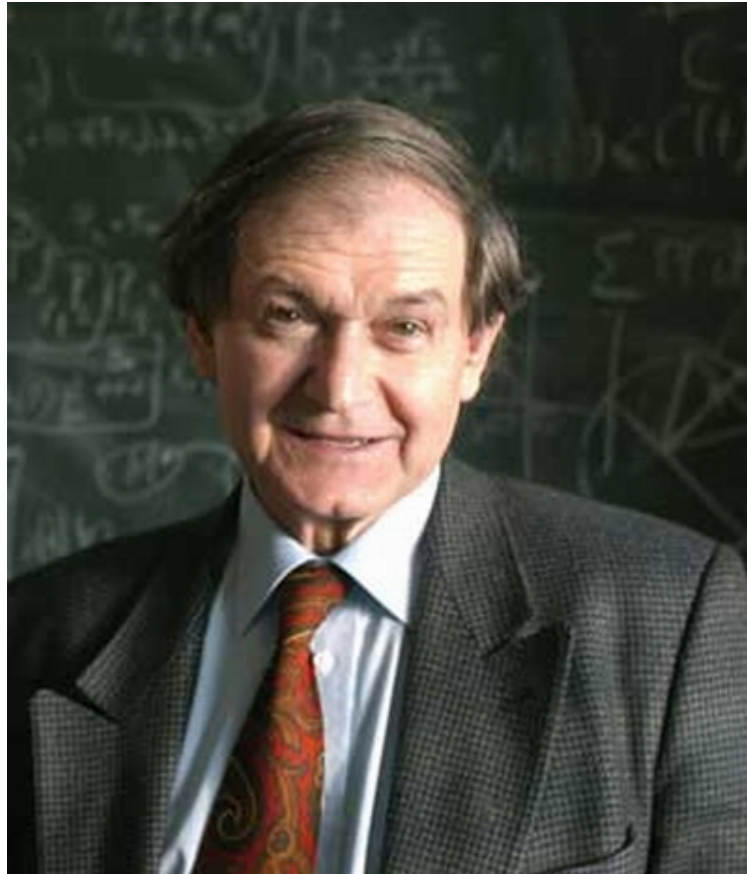
Claude LeBrun
Stony Brook University

Oxford, September 13, 2013

For Roger Penrose

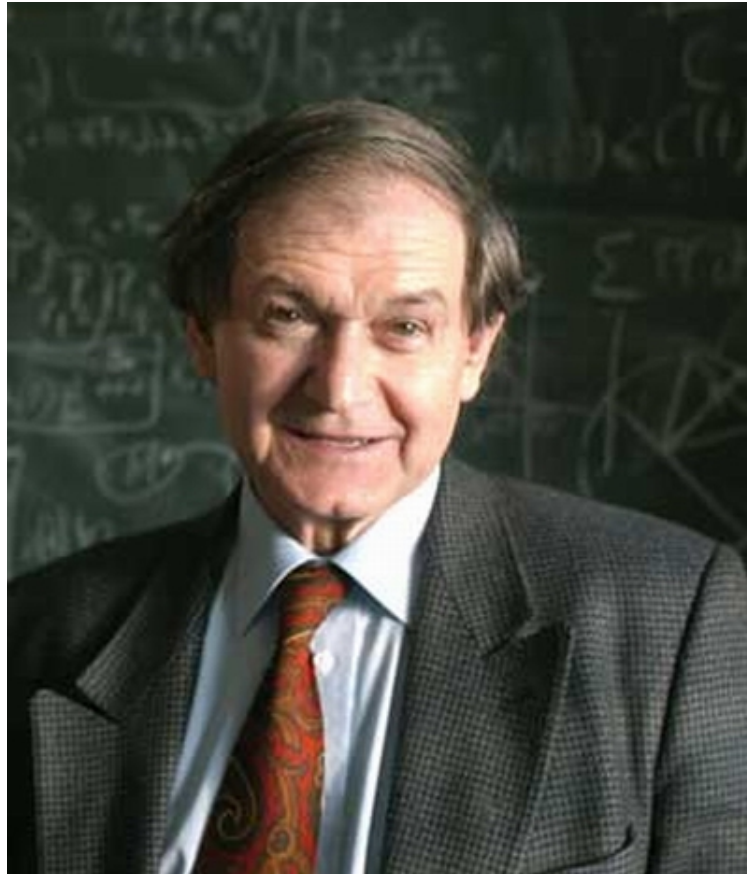


For Roger Penrose



who discovered so many remarkable links

For Roger Penrose



who discovered so many remarkable links between
complex manifolds and space-time geometry;

and Paul Tod



and Paul Tod



whose results revealed

and Paul Tod



whose results revealed key features
of the Einstein-Weyl equations.

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But — story takes place in $(2 + 1)$ -dimensions!

Joint work with

Lionel Mason

University of Oxford

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Main references:

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- The Einstein-Weyl Equations, Scattering Maps, and Holomorphic Disks, Math. Res. Lett. 16 (2009) 291–301.

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Weyl's 1918 gauge theory

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based on Weyl connections $([g], \nabla)$:

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$$\nabla_v g \propto g \quad \forall v$$

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$$\nabla \rightsquigarrow \nabla + \delta_k^j \nu_\ell + \delta_\ell^j \nu_k - \nu^j g_{k\ell}$$

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where $\nu = d \log u$.

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where $n = \dim M$.

Hermann Weyl



$$\mathfrak{B} = (\mathfrak{G} + \alpha \mathbf{I}) + \frac{\varepsilon^2}{4} V \bar{g} \{ \mathbf{1} - 3 (\varphi_i \varphi^i) \},$$

$$\Gamma_{ik}^r = \left\{ \begin{matrix} ik \\ r \end{matrix} \right\} + \frac{1}{2} \varepsilon^2 (\delta_i^r \varphi_k + \delta_k^r \varphi_i - g_{ik} \varphi^r).$$

Unter Vernachlässigung der winzigen kosmologischen Terme erhalten wir hier also genau die klassische Maxwell-Einsteinsche Theorie der Elektrizität und Gravitation. Um Übereinstimmung mit den in § 34 verwendeten

Ricci tensor

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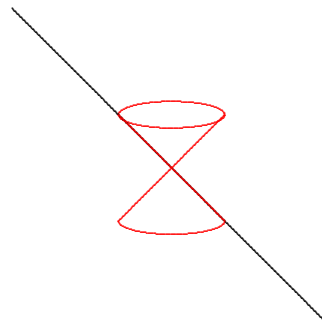
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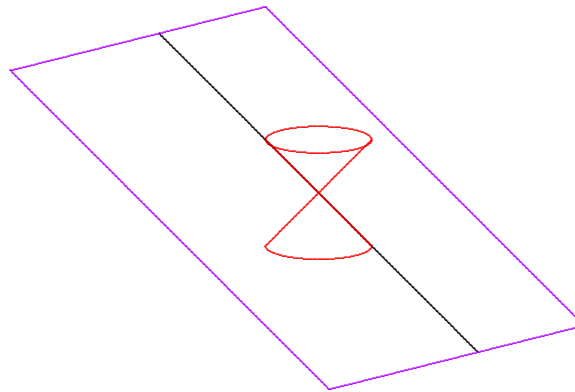
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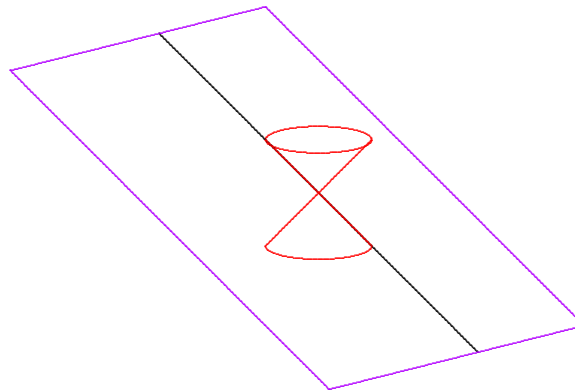
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Equations $\iff \exists$ totally geodesic null surfaces.

Élie Cartan



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THÉORÈME. — Les espaces de Weyl à trois dimensions qui admettent ∞^2 plans isotropes dépendent essentiellement de quatre fonctions arbitraires de deux arguments.

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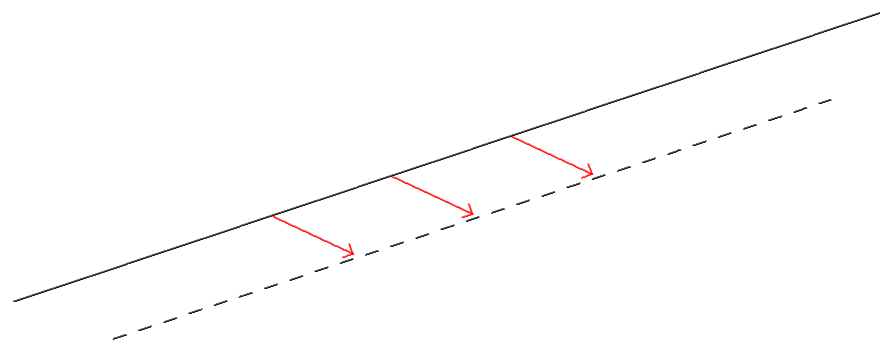
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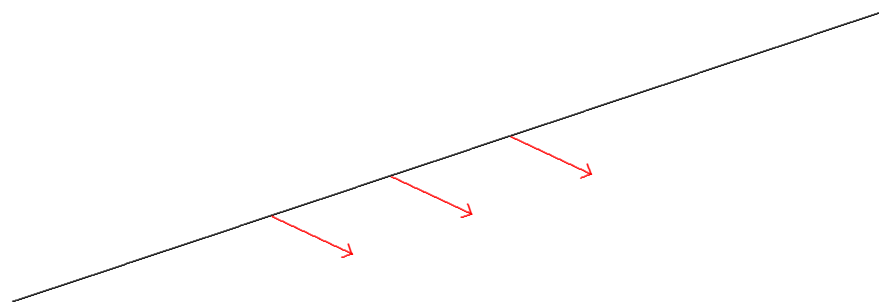
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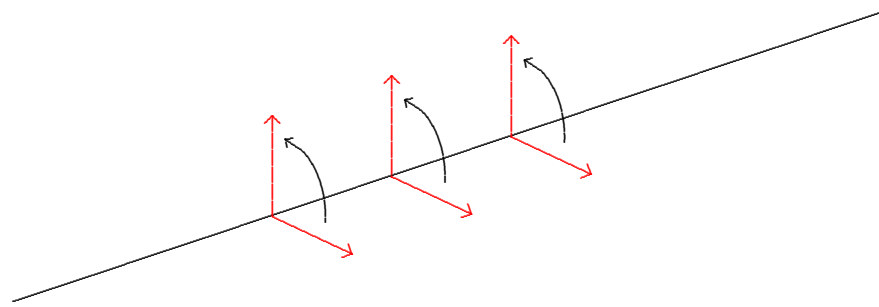
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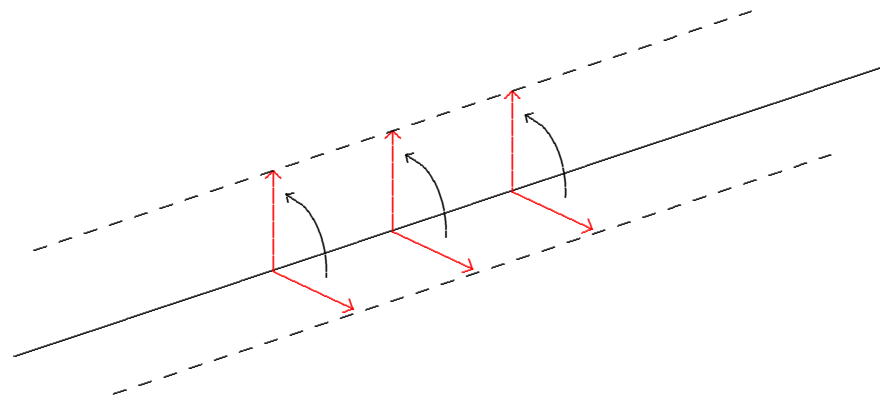
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Nigel Hitchin



and so if

$$R(U \times V, U)U = U \times R(V, U)U, \quad (2.2)$$

then we can define a linear map

$$J(V) = U \times V \quad (2.3)$$

which satisfies

$$J^2(V) = U \times (U \times V) = (U, V)U - (U, U)V = -V$$

We thus have a real complex surface G with a family of real lines of self-intersection number 2. It can be shown that any such surface may be obtained by the above geodesic construction, but using a Weyl structure rather than a Riemannian structure. The integrability condition (2.2) is then the analogue of Einstein's equations $(R_{(ij)} = \Lambda g_{ij})$ for the Weyl structure (see [10]). This is the

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- *orientation-reversing diffeomorphisms*

$$\psi : \mathbb{CP}_1 \rightarrow \mathbb{CP}_1.$$

space-time oriented

Conformal Lorentzian n -manifold $(M, [g])$ called
space-time oriented

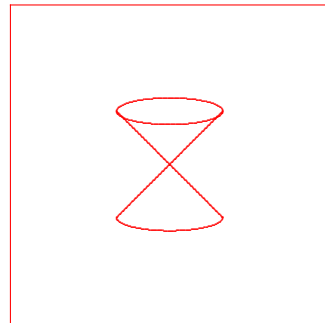
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\implies time-orientation: future vs. past.

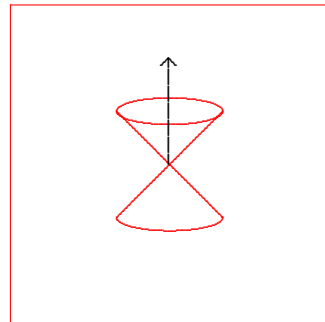
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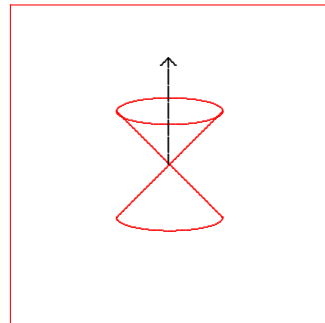
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$\implies M$ also oriented, in usual sense.

globally hyperbolic

Time-oriented conformal Lorentzian n -manifold $(M, [g])$
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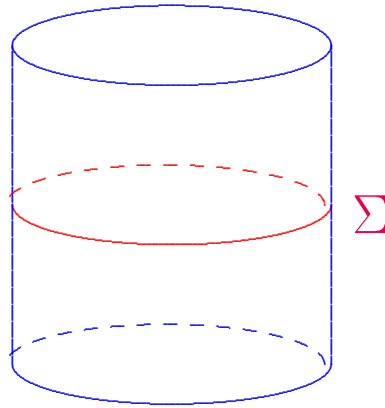
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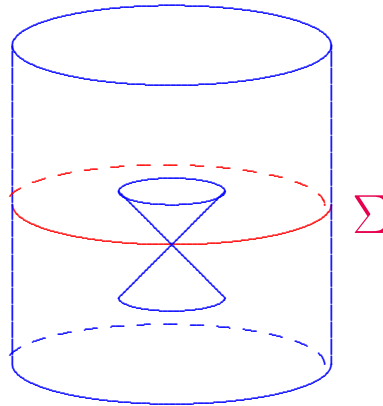
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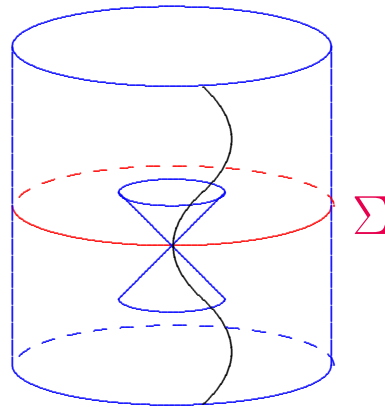
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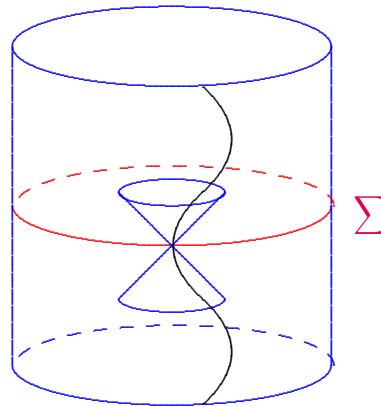
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$$\implies M \approx \Sigma \times \mathbb{R}$$

conformally compact

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$$M = SO(3, 1)/SO(2, 1)$$

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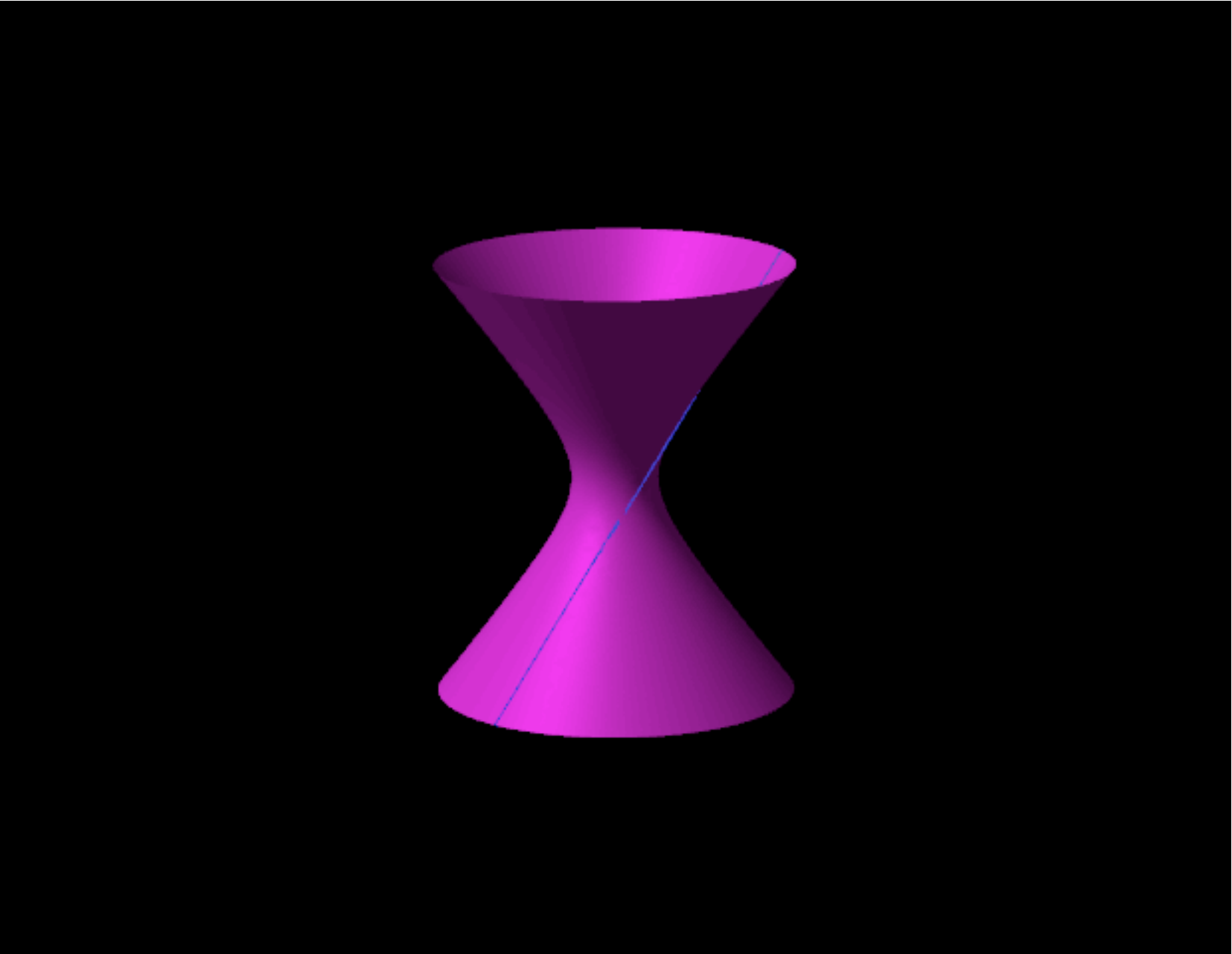
M = hypersurface

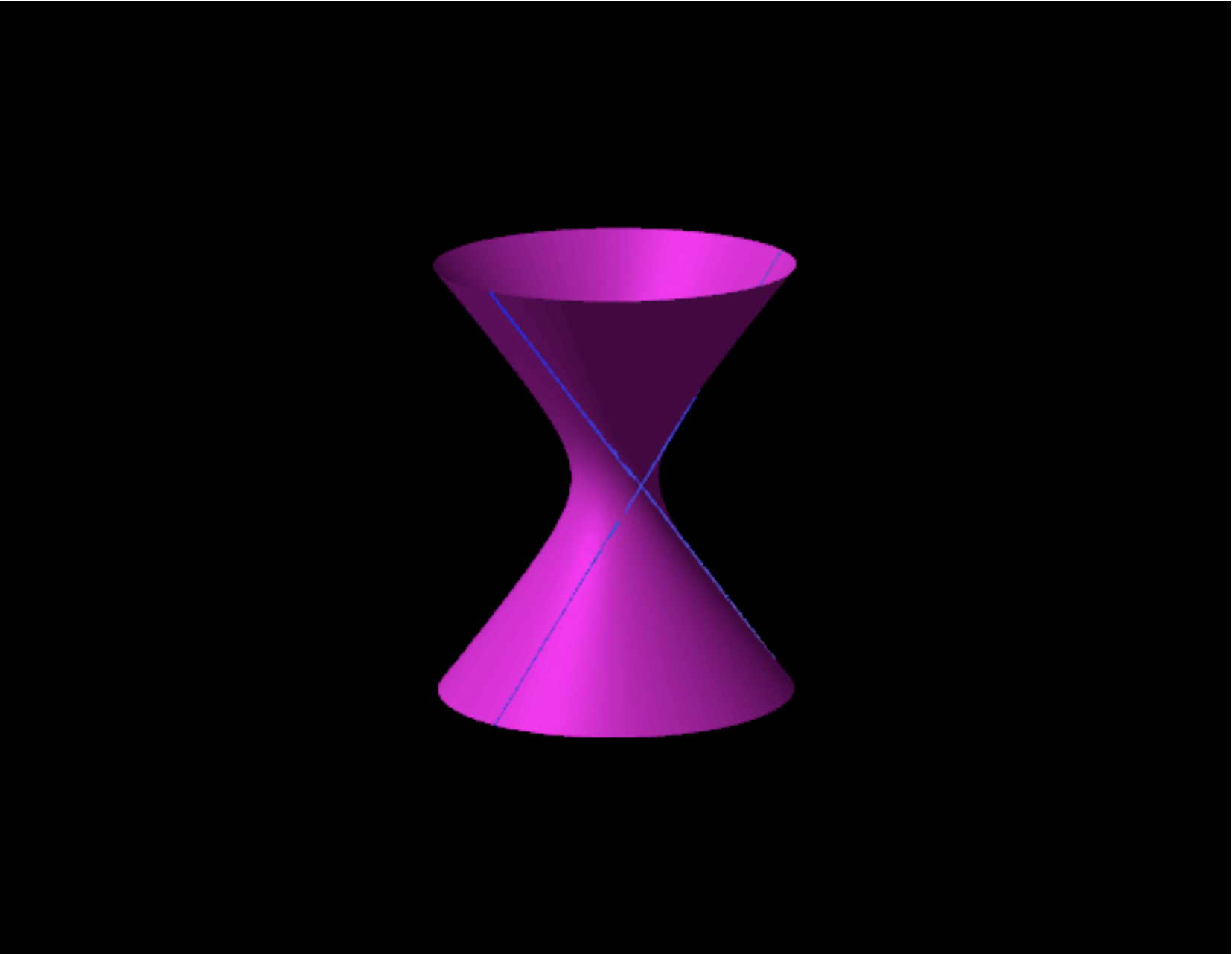
$$x^2 + y^2 + z^2 - t^2 = 1$$

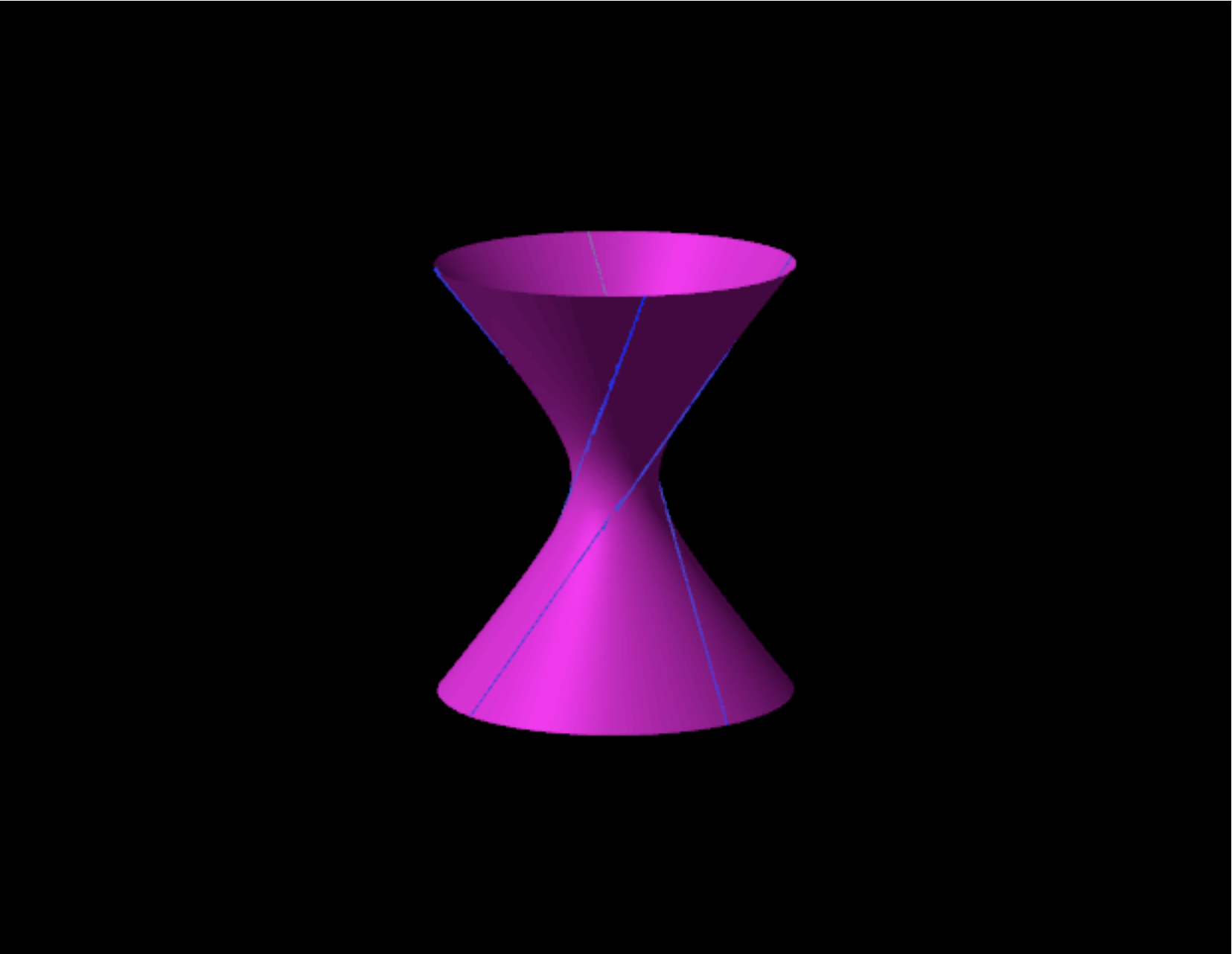
in Minkowski space \mathbb{R}^4

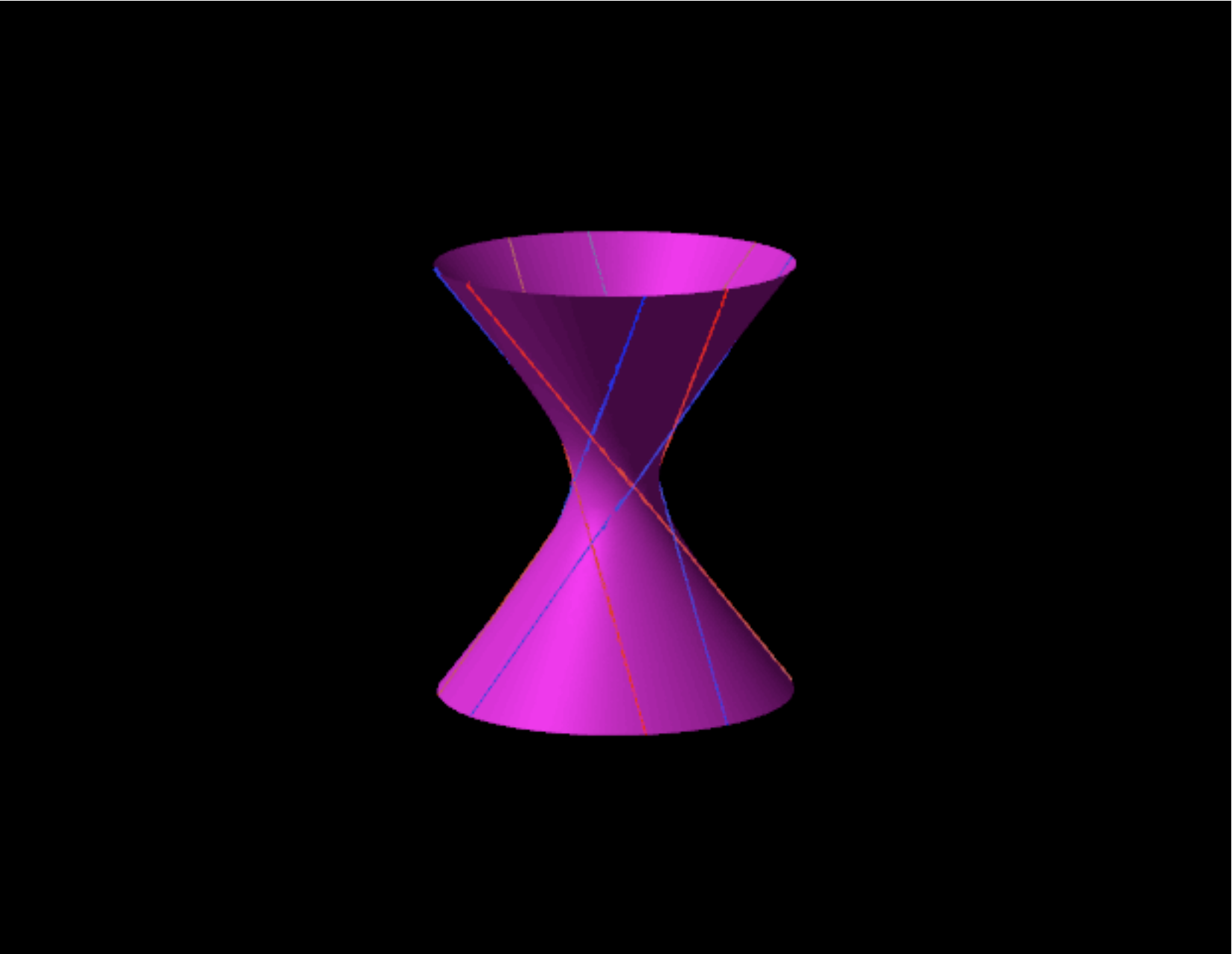
$$g = dx^2 + dy^2 + dz^2 - dt^2$$

$$M = SL(2, \mathbb{C}) / SL(2, \mathbb{R})$$









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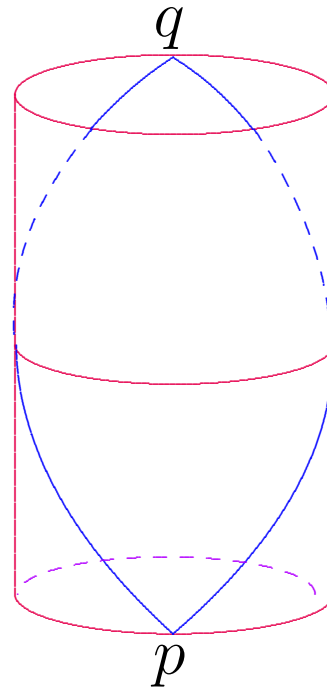
where $u = \sin \tau$.

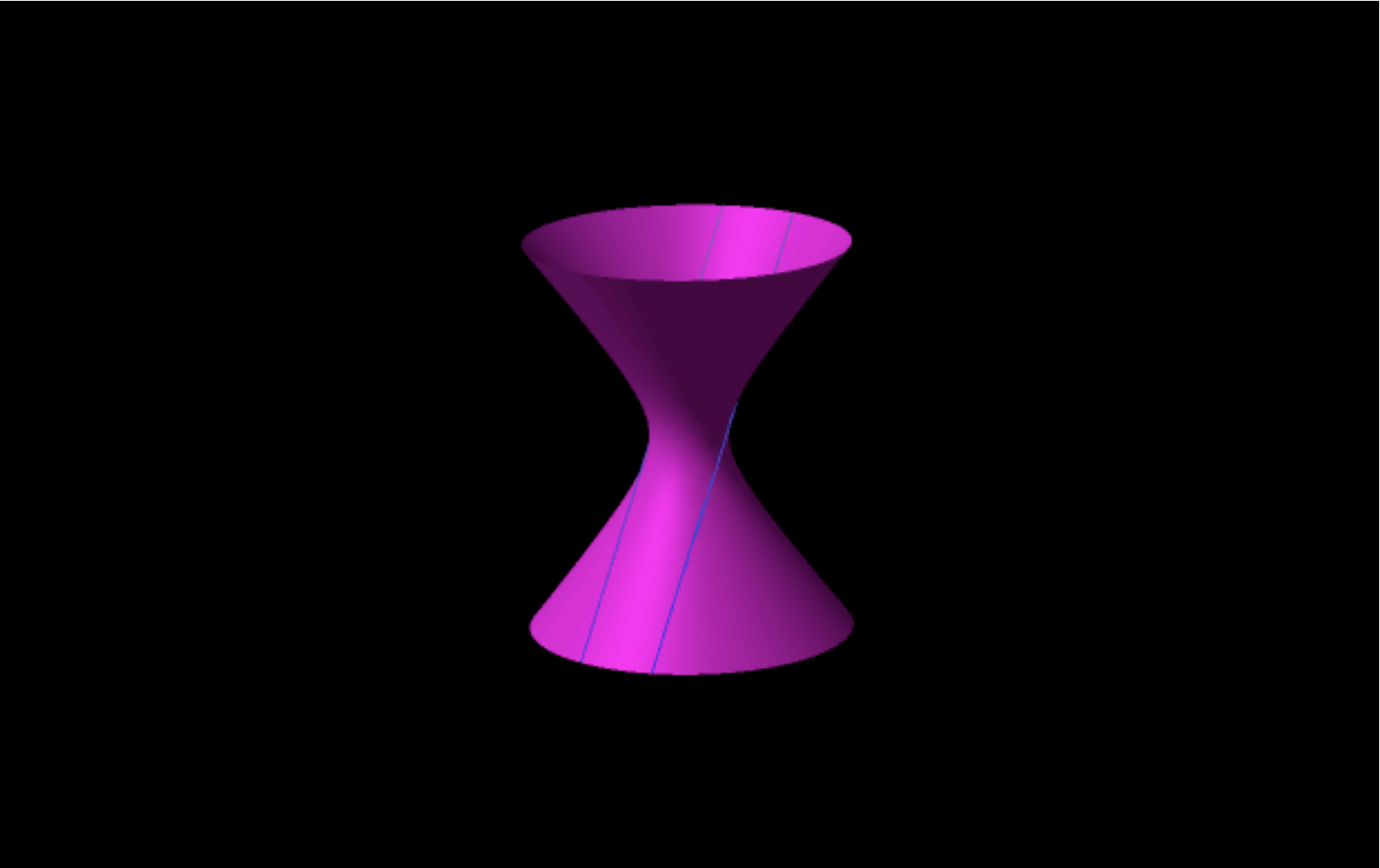
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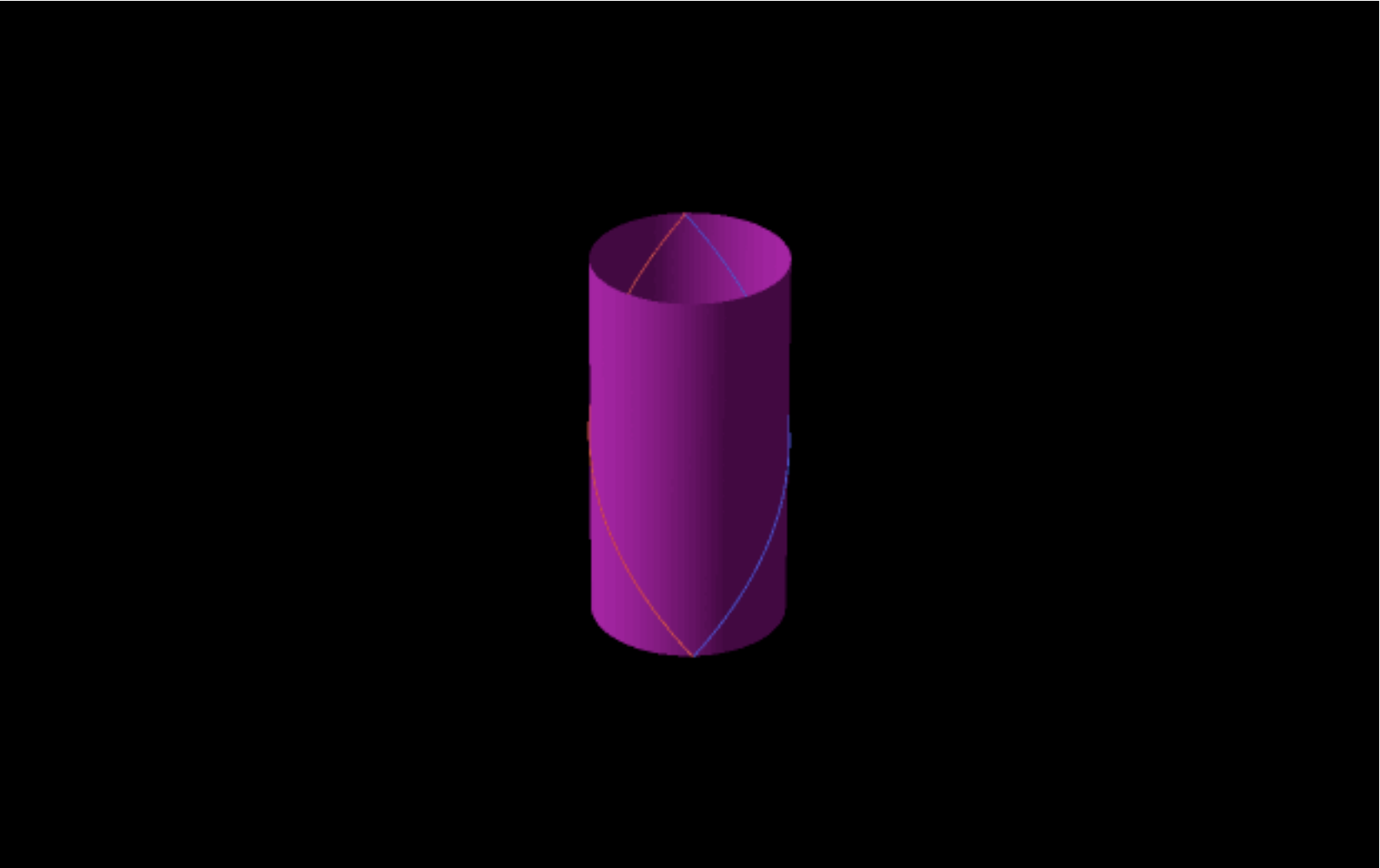
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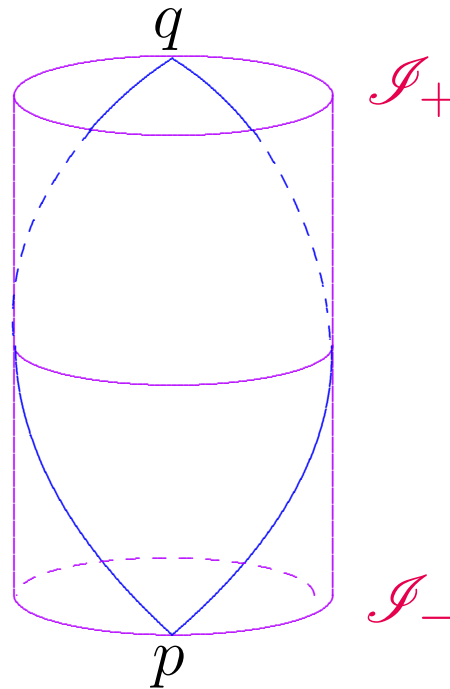


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Theorem. *There is a natural one-to-one correspondence between*

- *smooth, space-time-oriented, conformally compact, globally hyperbolic Lorentzian Einstein-Weyl 3-manifolds $(M, [g], \nabla)$; and*
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$$\psi_1, \psi_2 : \mathbb{CP}_1 \rightarrow \mathbb{CP}_1$$

considered same iff

$$\psi_1 = \varphi \circ \psi_2 \circ \phi^{-1}$$

for Möbius transformations $\varphi, \phi \in PSL(2, \mathbb{C})$.

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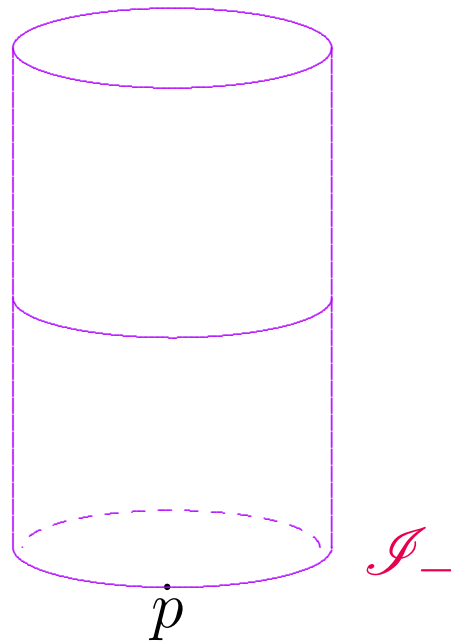
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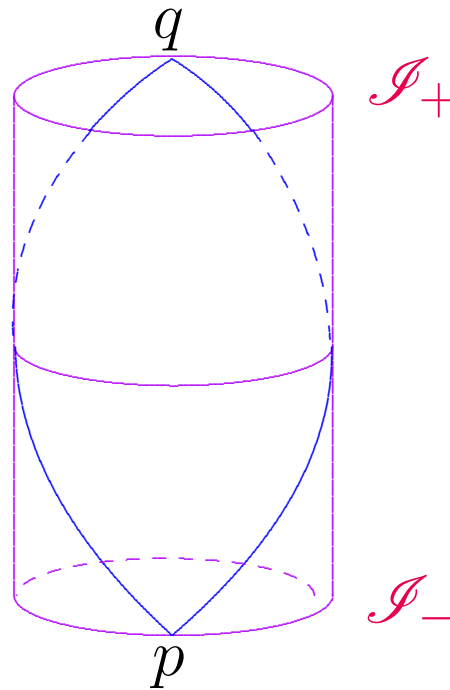
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to Einstein-Weyl $(M^3, [g], \nabla)$ satisfying hypotheses.

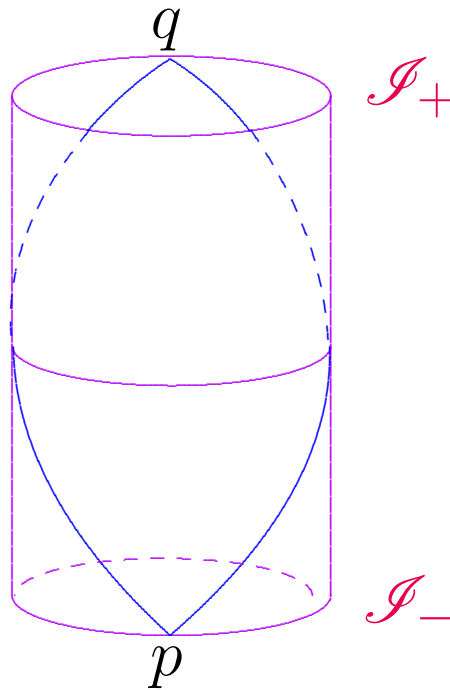
Lemma. *For an Einstein-Weyl manifold as above, let $p \in \mathcal{I}_-$ be any point of past infinity.*



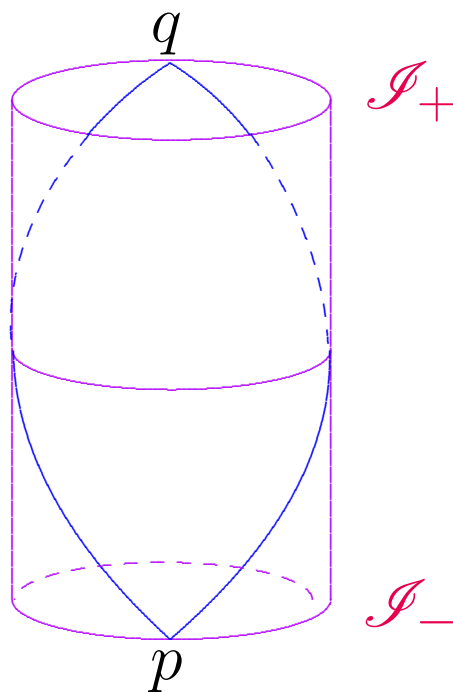
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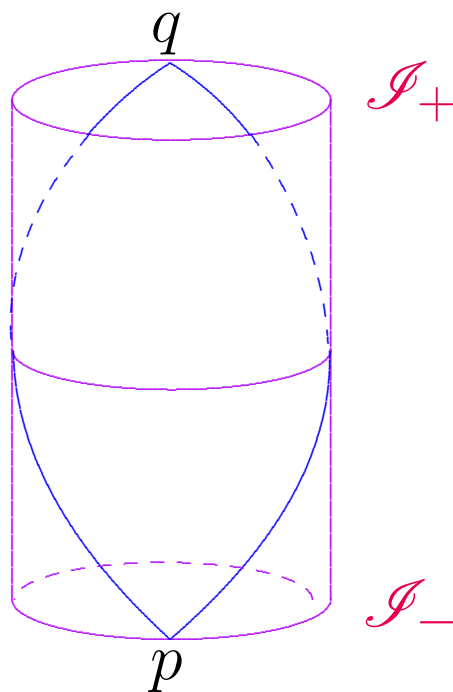
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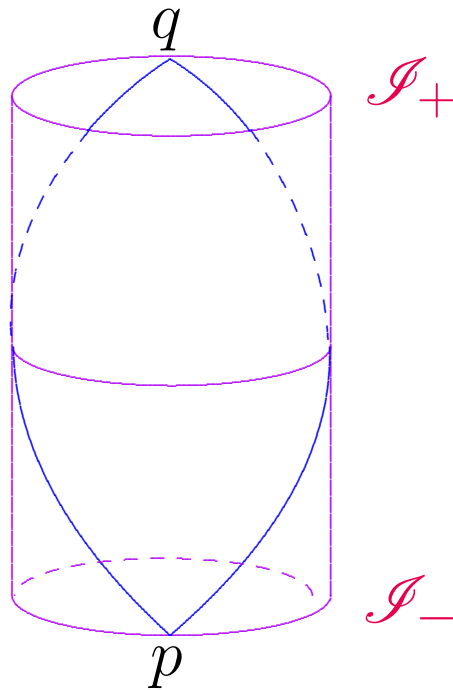
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Inverting correspondence:

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twistor disk construction

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Graph of orientation-reversing diffeomorphism

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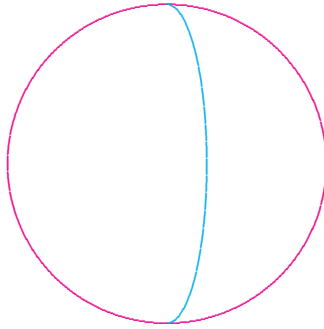
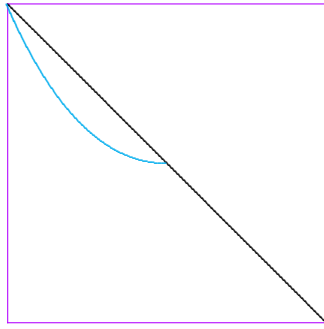
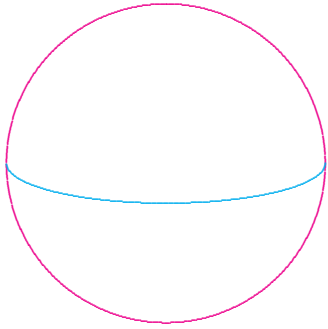
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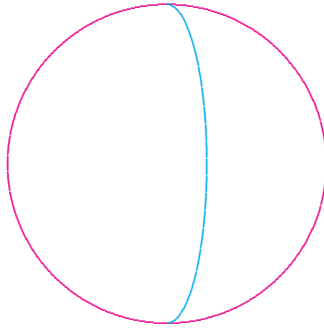
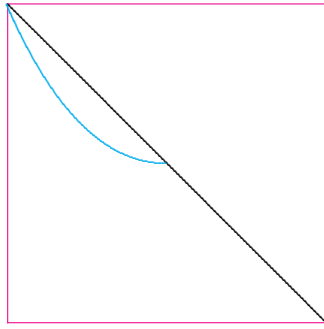
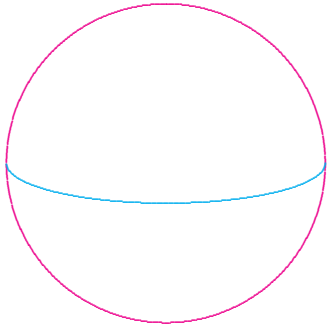
Strategy: construct 3-manifold $M = M_\psi$
as moduli space of holomorphic disks D
in $Z = \mathbb{CP}_1 \times \mathbb{CP}_1$ with ∂D on $P \subset Z$.



When ψ is the antipodal map,
disks are explicitly given by

$$\zeta \longmapsto ([a\zeta + b : c\zeta + d], [-\bar{d}\zeta - \bar{c} : \bar{b}\zeta + \bar{a}])$$

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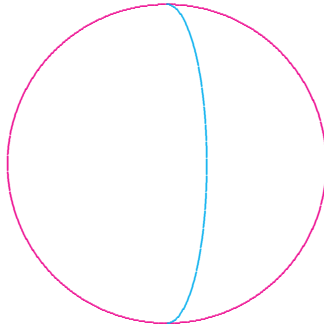
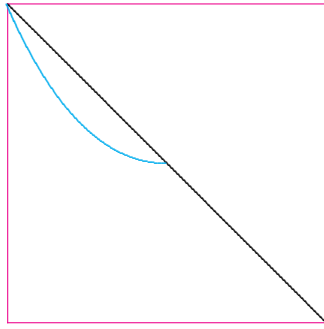
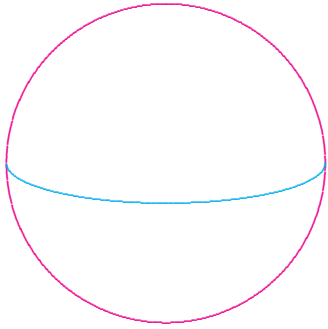
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Moduli space M of disks mod reparameterization:

de Sitter space $SL(2, \mathbb{C})/SL(2, \mathbb{R})$.

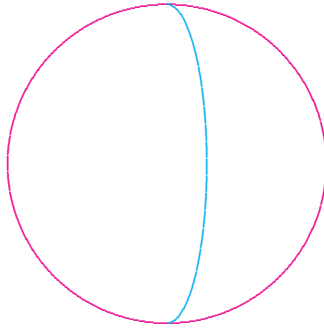
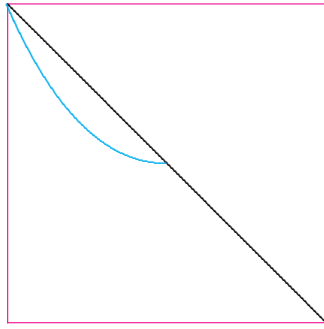
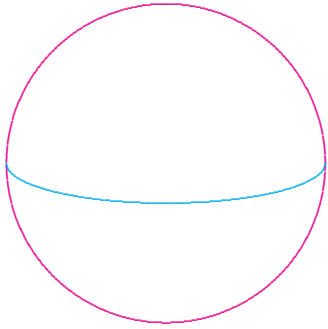


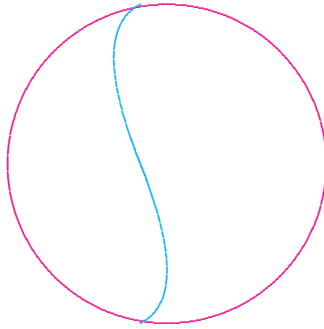
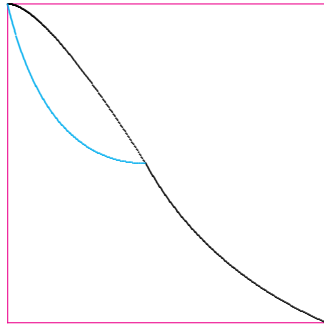
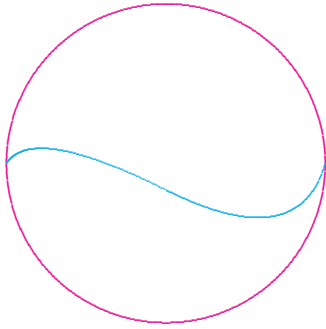
Now deform $P \hookrightarrow \mathbb{C}P_1 \times \mathbb{C}P_1$

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by replacing graph of anti-podal map with
graph of orientation-reversing diffeomorphism

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Lemma. *Let $\psi : \mathbb{C}\mathbb{P}_1 \rightarrow \mathbb{C}\mathbb{P}_1$ be any orientation-reversing diffeomorphism, and let*

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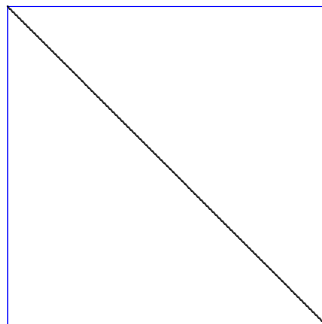
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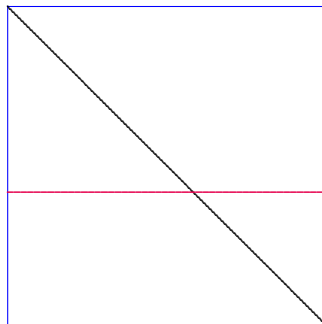
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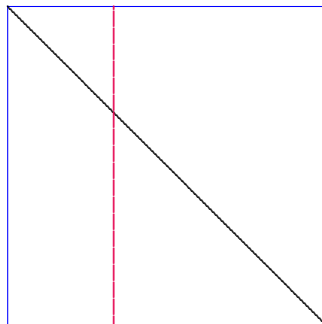
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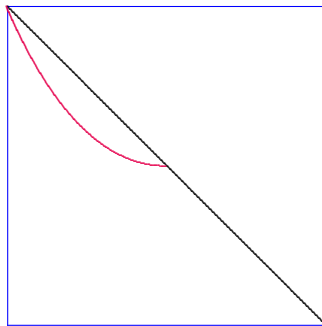
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If $(\Sigma, \partial\Sigma) \rightarrow (Z, P)$ is any holomorphic curve with boundary representing \mathbf{a} , then Σ is either a holomorphic disk as above, or is a factor $\mathbb{C}\mathbb{P}_1$ of $Z = \mathbb{C}\mathbb{P}_1 \times \mathbb{C}\mathbb{P}_1$.

Proof.

Proof. Regularity:

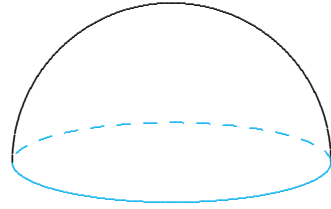
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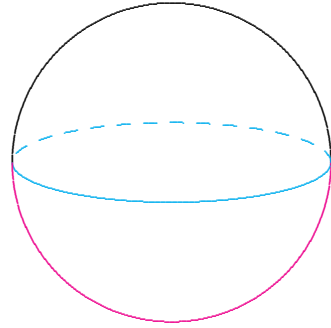
But need to show is that disk is actually a graph!

Proof. Consider abstract double $D \cup_{\partial D} \overline{D}$.

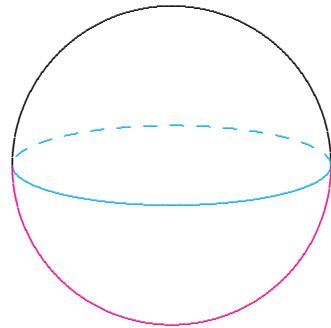
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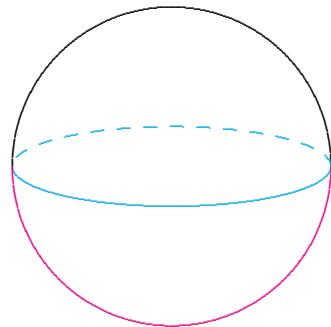
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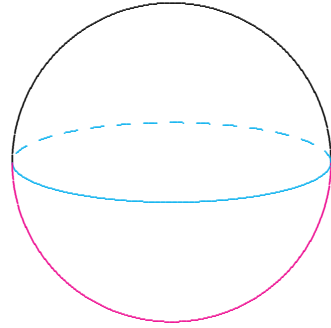
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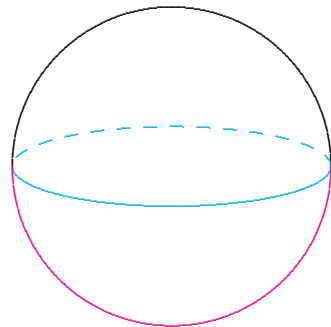
and continuous map (quasi-regular/quasi-conformal)

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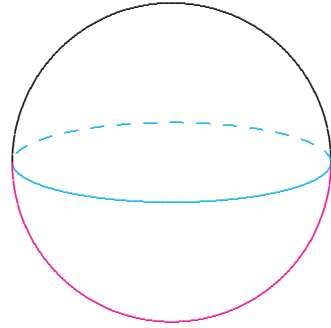
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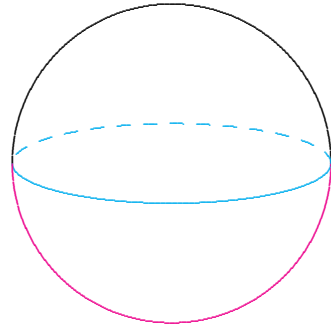
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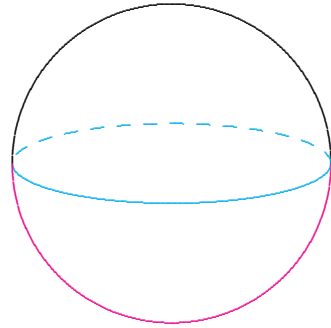
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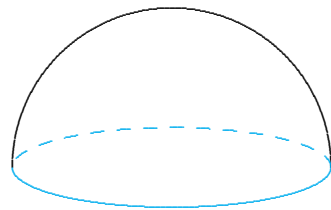
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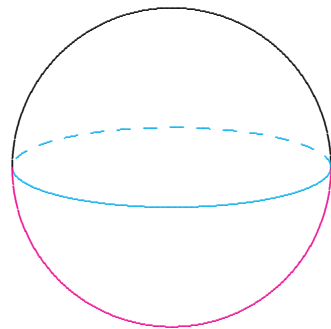
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Equals 2 in our case:

$$E \cong \mathcal{O}(2).$$

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$$\begin{aligned}h^1(\mathbb{C}\mathbb{P}_1, \mathcal{O}(2)) &= 0 \\h^0(\mathbb{C}\mathbb{P}_1, \mathcal{O}(2)) &= 3\end{aligned}$$

cf. Kodaira's Theorem
on deformation of complex submanifolds

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(Forsternic, Gromov, et al.)

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Disks all have same ω -area. \implies
any sequence has convergent subsequence...

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Tricky point: disks can degenerate to factor $\mathbb{C}\mathbb{P}_1$.

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Since ψ is continuous deformation of antipodal,

Continuity method \Rightarrow each level set **non-empty!**

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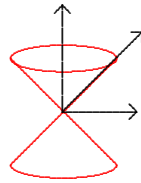
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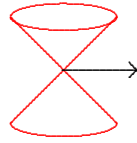
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\Rightarrow up to homothety $T_D M$ carries Lorentz metric, modelled on Killing form of $\mathfrak{sl}(2, \mathbb{R})$.

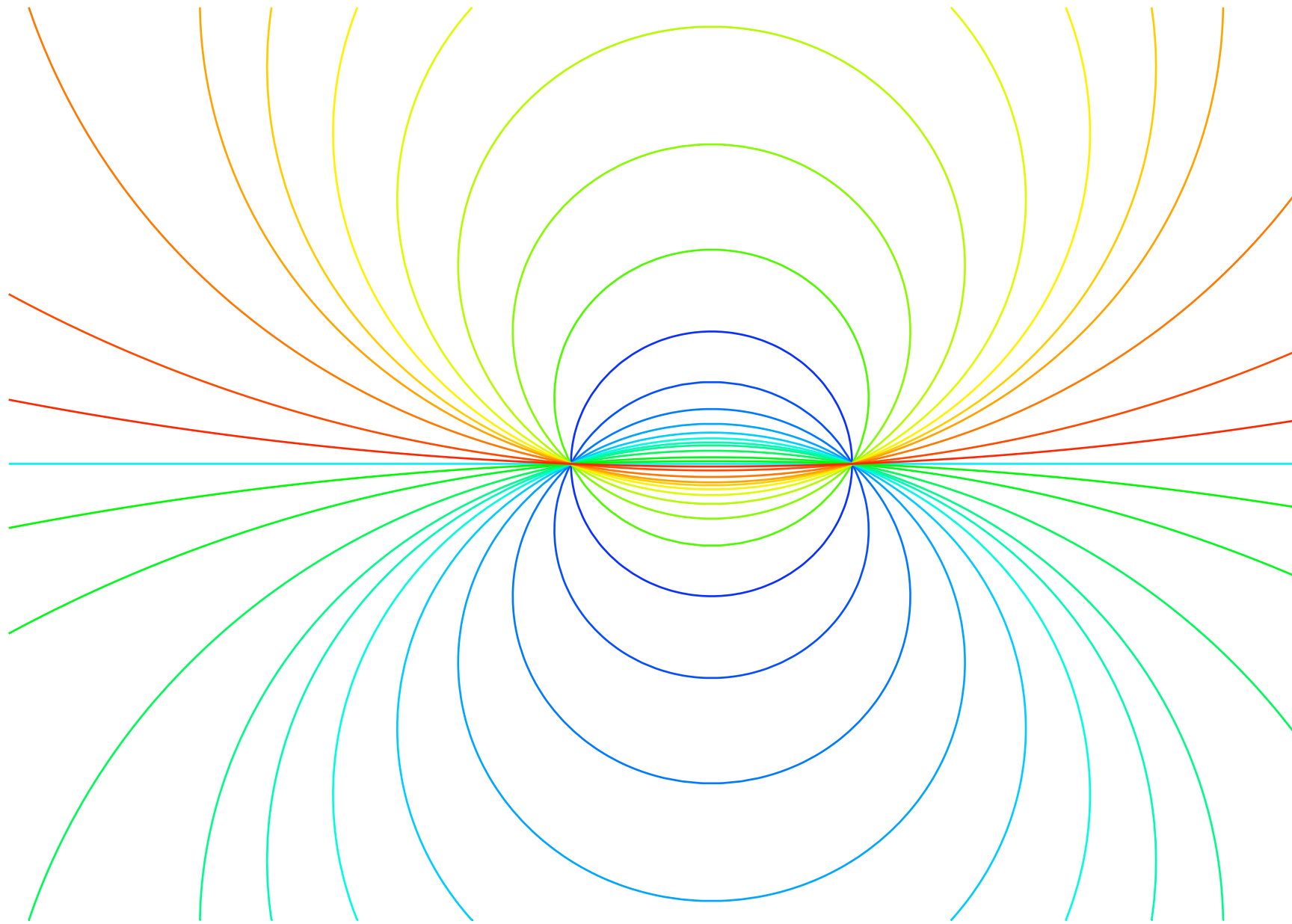
Trichotomy:

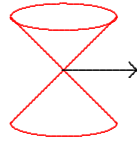
TM	$\mathfrak{sl}(2, \mathbb{R})$
space-like	hyperbolic
null	parabolic
time-like	elliptic



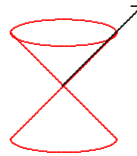


Space-like vector = infinitesimal variation with
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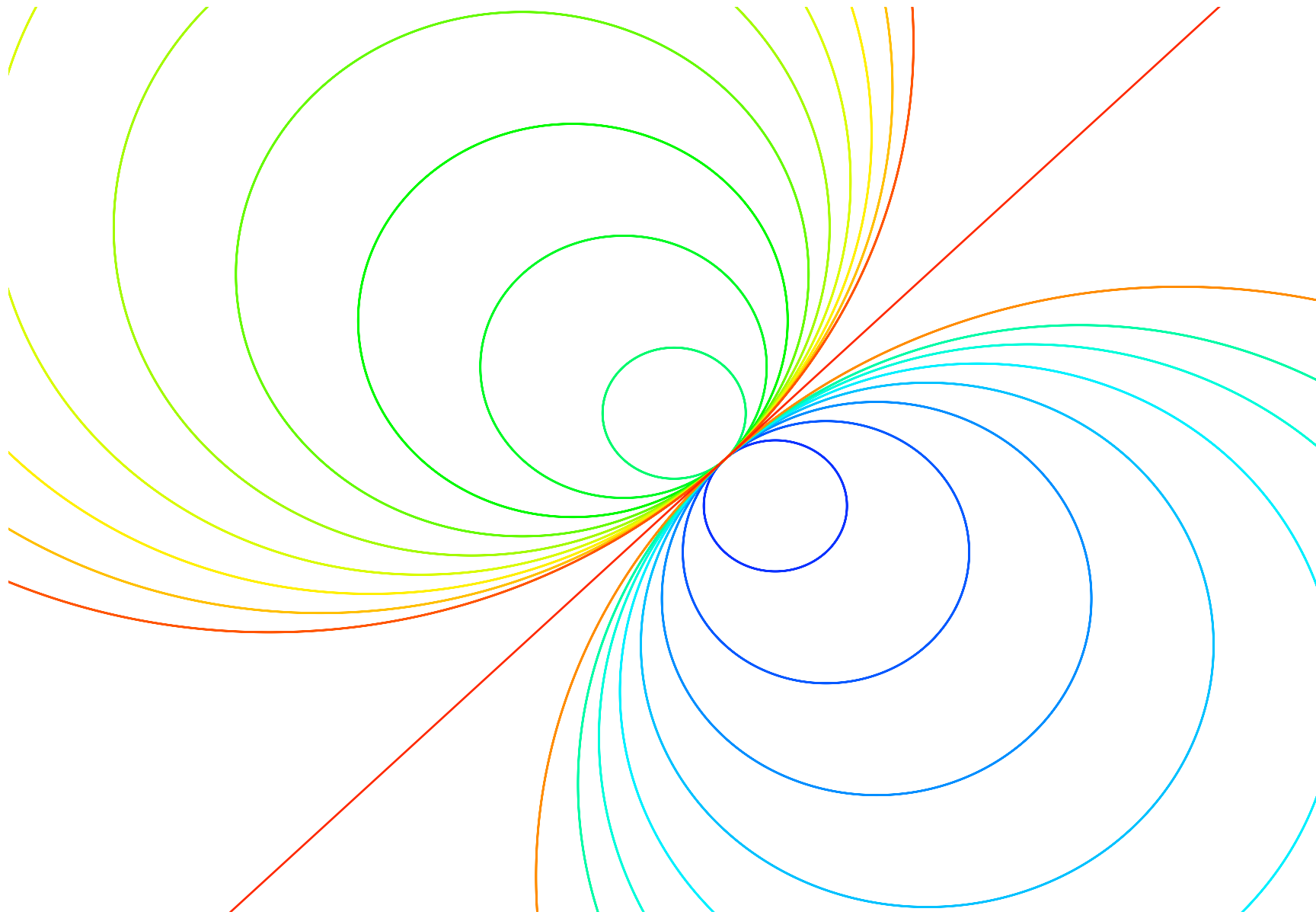


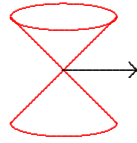


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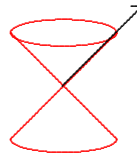


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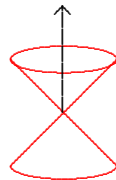




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By deformation: Cauchy surface topologically S^2 .

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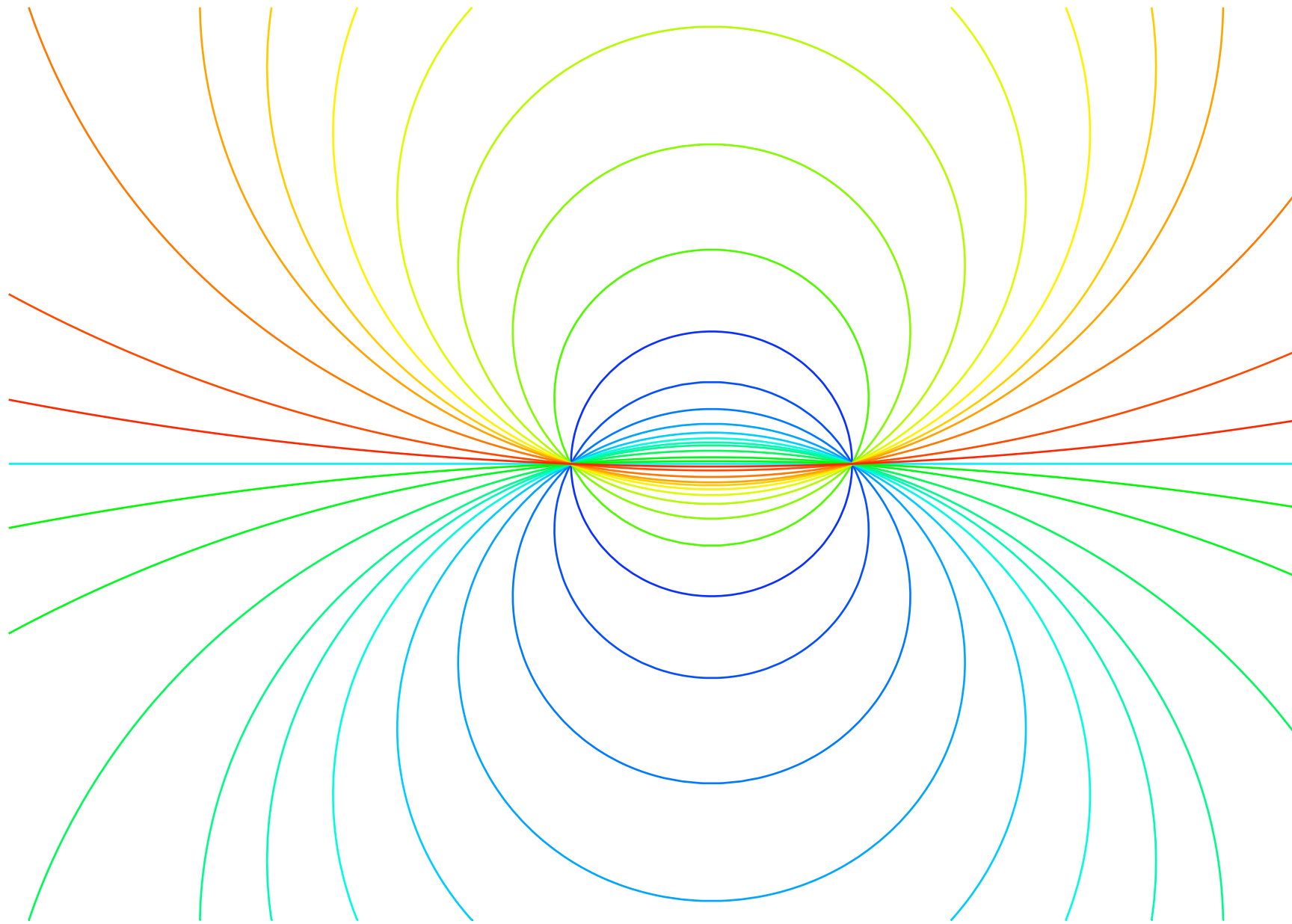
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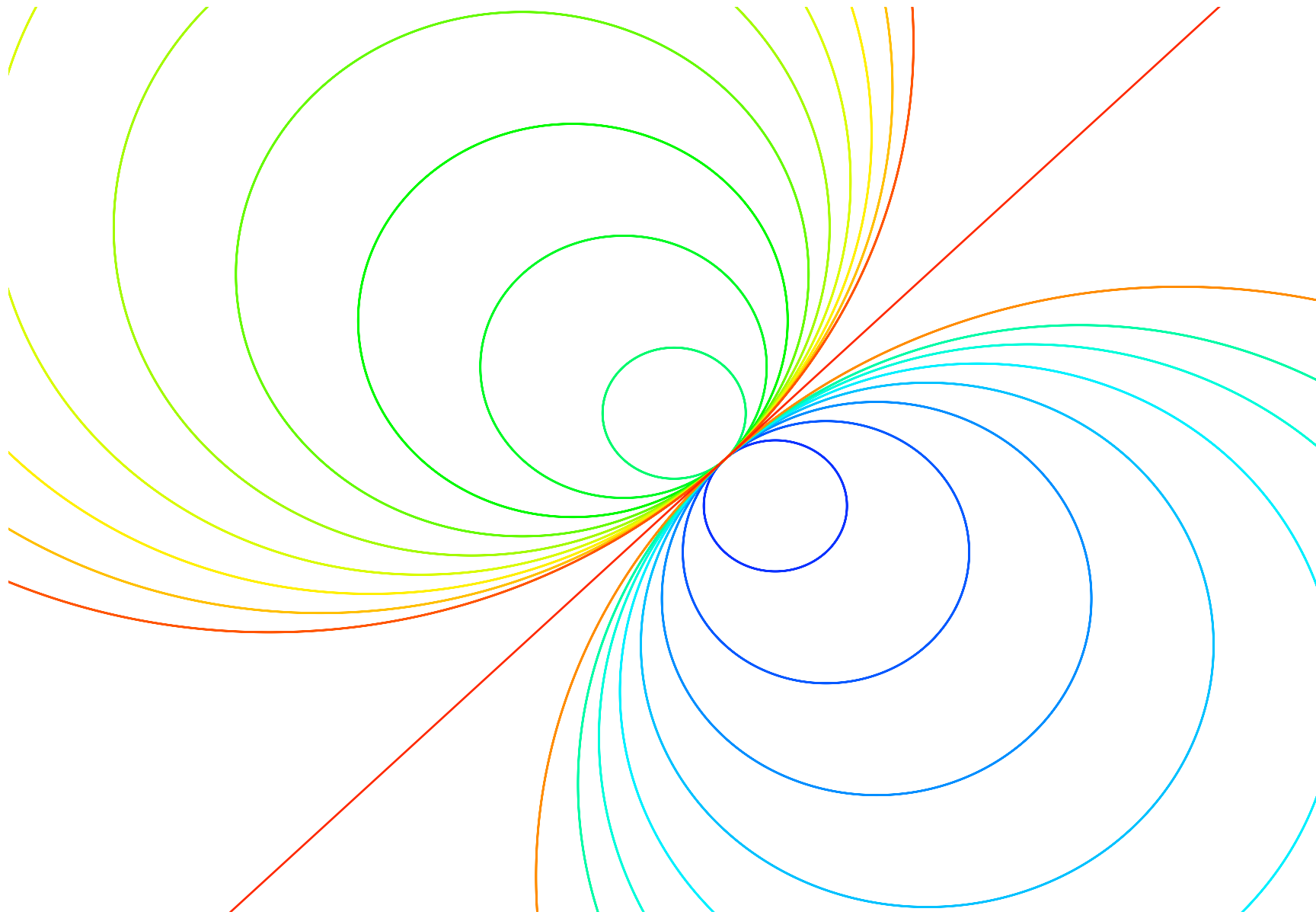
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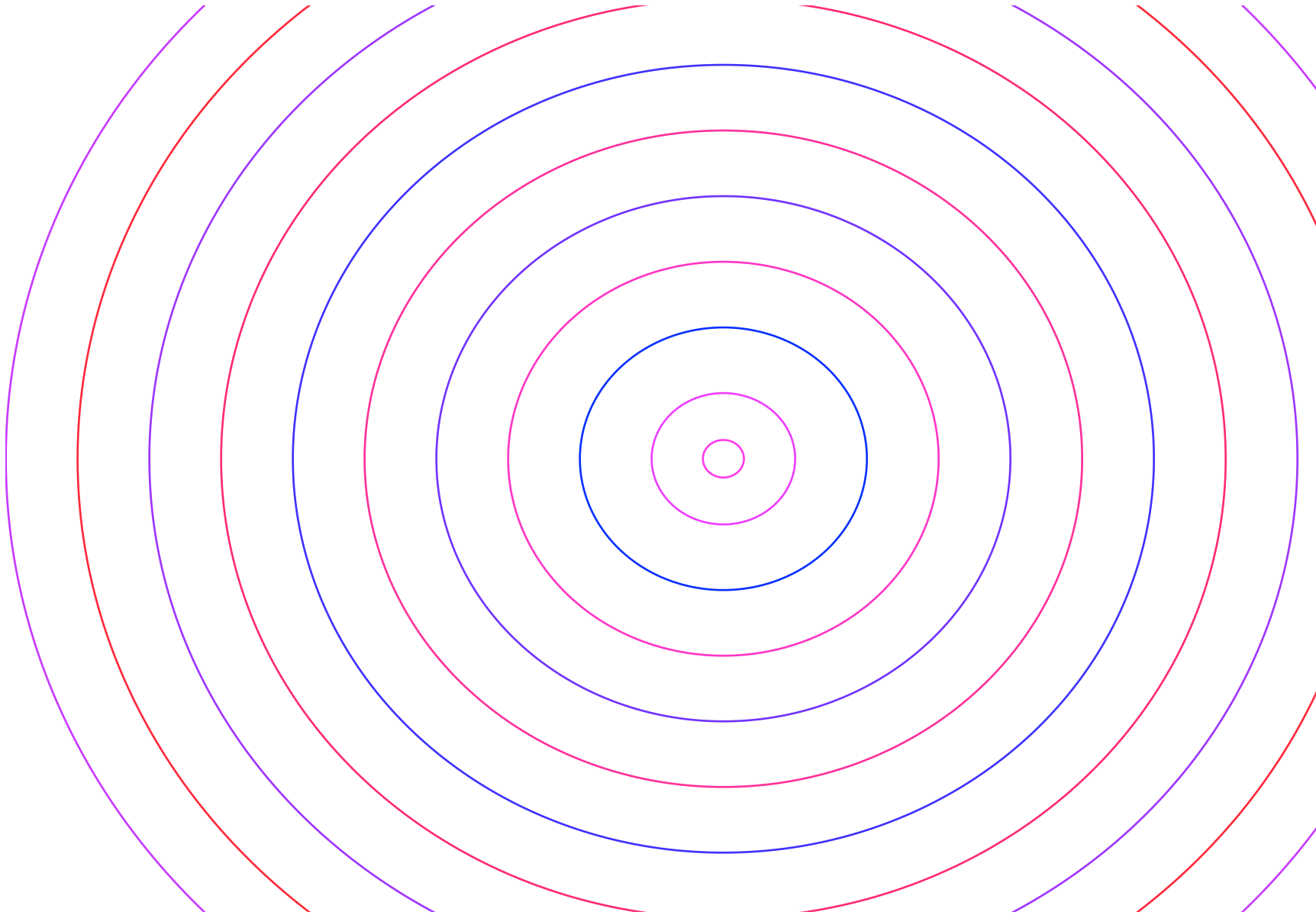
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Weyl connection ∇ :

Geodesics:

Space-like geodesic:

disks passing through a pair of distinct points $x \neq y$ in P .

Null geodesic:

disks through a given point $x \in P$
with specified tangent.

Time-like geodesic:

holomorphic disks through given point $x \in Z - P$.

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- — & Mason, Duke Math. J. 136 (2007) 205–273.

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Also gives direct proof of conformal compactness.

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