

Einstein Manifolds,
Kähler Metrics, &
Conformal Geometry

Claude LeBrun
Stony Brook University

Special Geometries on Riemannian Manifolds
Montreal, October 13, 2021

Most recent references:

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Bach-Flat Kähler Surfaces

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Journal of Geometric Analysis
30 (2020) 2491–2514

Einstein Manifolds, Self-Dual Weyl Curvature,
and Conformally Kähler Geometry

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Mathematical Research Letters
28 (2021) 127–144

And

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Einstein Manifolds, Conformal Curvature, and
Anti-Holomorphic Involutions

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Annales Mathématiques du Québec
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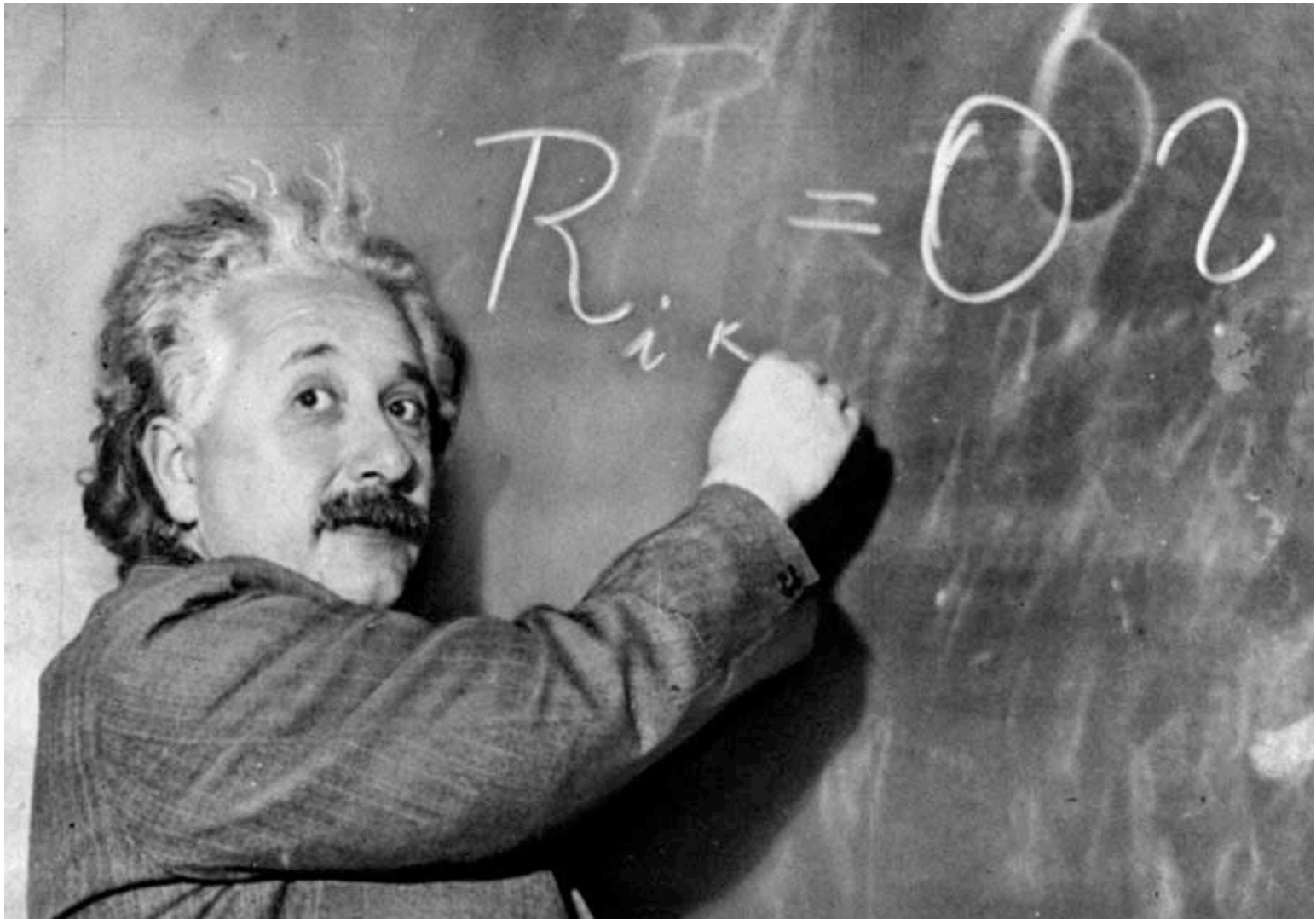
“... the greatest blunder of my life!”

— A. Einstein, to G. Gamow

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Has same sign as the *scalar curvature*

$$s = r_j^j = \mathcal{R}^{ij}_{ij}.$$

Dimension Four is Exceptional

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When $n = 4$, Einstein metrics satisfy a remarkable conformally-invariant condition.

On Riemannian n -manifold (M, g) ,

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$$\mathcal{R}^{ab}_{cd} = W^{ab}_{cd} + \frac{4}{n-2} \overset{\circ}{r} \begin{bmatrix} a & b \\ c & d \end{bmatrix} + \frac{2}{n(n-1)} \mathfrak{s} \delta \begin{bmatrix} a & b \\ c & d \end{bmatrix}$$

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W^a_{bcd} unchanged if $g \rightsquigarrow \hat{g} = u^2 g$.

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Proposition. Assume $n \geq 4$. Then

(M^n, g) locally conformally flat $\iff W \equiv 0$.

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Measures deviation $[g]$ from conformal flatness.

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Of course, conformally Einstein good enough!

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But when $n \neq 4$, Einstein $\not\Rightarrow$ critical point of \mathcal{W} !

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When $n = 4$, conf. Einstein \implies critical for \mathcal{W} .

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$$\mathcal{R} = \left(\begin{array}{c|c} W_+ + \frac{s}{12} & \mathring{r} \\ \hline \mathring{r} & W_- + \frac{s}{12} \end{array} \right)$$

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	Λ^{+*}	Λ^{-*}
Λ^+	$W_+ + \frac{s}{12}$	\mathring{r}
Λ^-	\mathring{r}	$W_- + \frac{s}{12}$

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$$\tau(M) = \frac{1}{12\pi^2} \int_M \left(|W_+|^2 - |W_-|^2 \right) d\mu$$

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Hence

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A case of special interest:

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(M^4, g, J) Kähler.

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Restriction of \mathcal{W}_+ to Kähler metrics?

On Kähler metrics,

$$\int |W_+|^2 d\mu = \int \frac{s^2}{24} d\mu$$

so any critical point of restriction must be **extremal** in sense of Calabi.

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Andrzej Derdziński : For Kähler metrics g ,

$$B = \frac{1}{12} \left[2s\dot{r} + \text{Hess}_0(s) + 3J^* \text{Hess}_0(s) \right]$$

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- $g_t = g + tB$ is Kähler metric for small t .

Action Function on Kähler Cone

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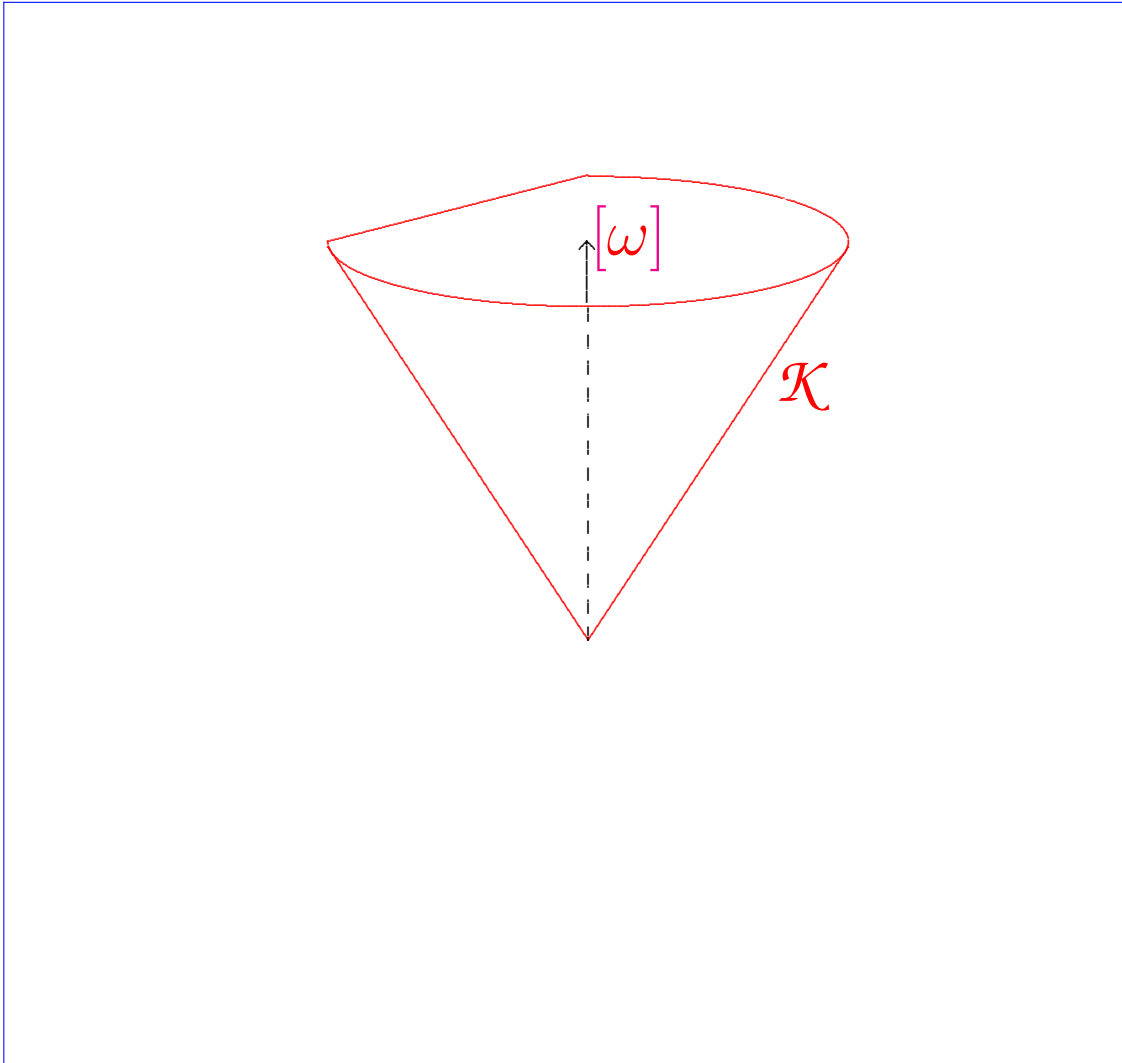
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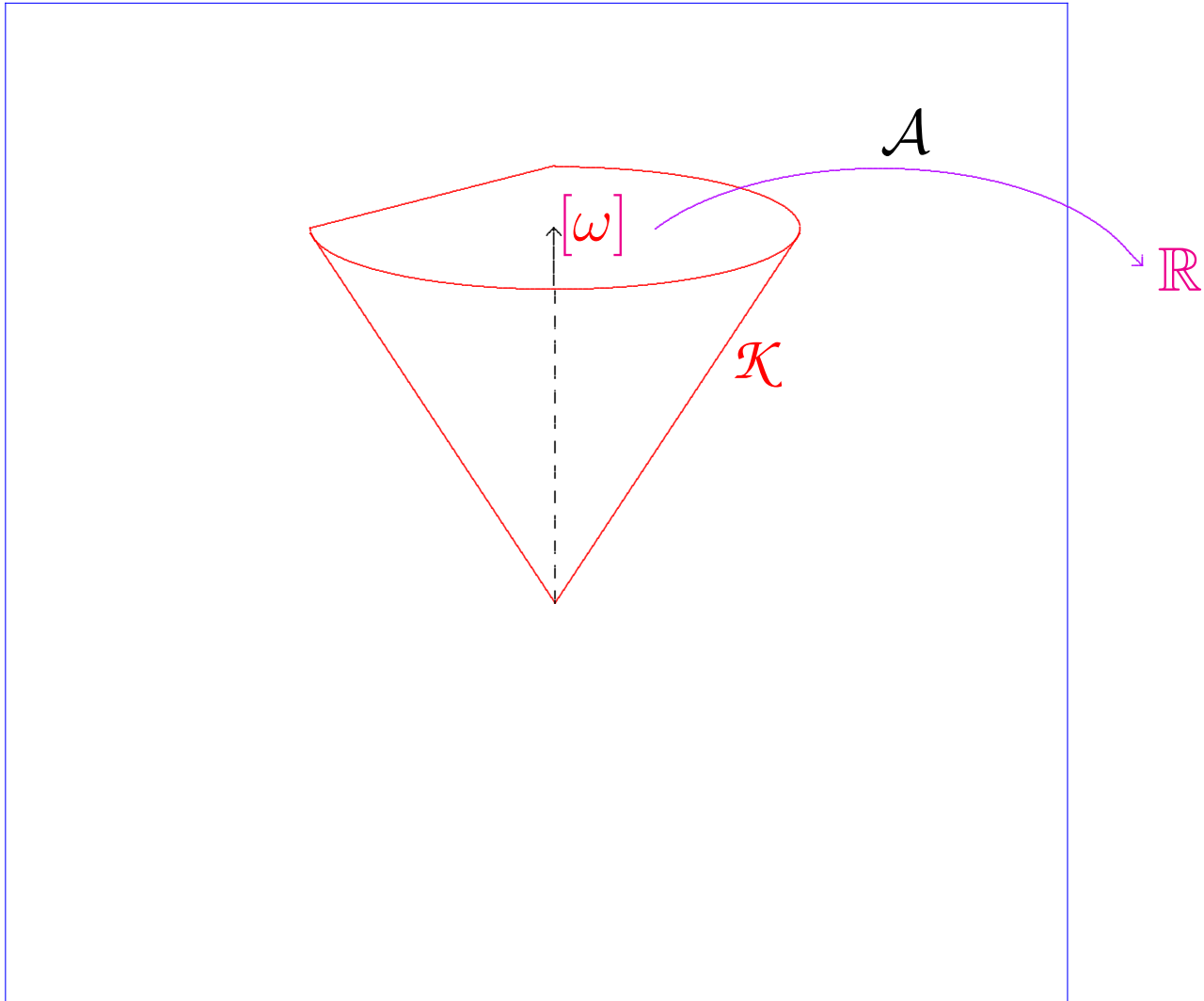
For any extremal Kähler (M^4, g, J) ,

$$\begin{aligned} \frac{1}{32\pi^2} \int s^2 d\mu_g &= \frac{(c_1 \cdot [\omega])^2}{[\omega]^2} + \frac{1}{32\pi^2} \|\mathcal{F}_{[\omega]}\|^2 \\ &=: \mathcal{A}([\omega]) \end{aligned}$$

where \mathcal{F} is Futaki invariant.



$$\mathcal{K} \subset H^{1,1}(M, \mathbb{R}) \subset H^2(M, \mathbb{R})$$



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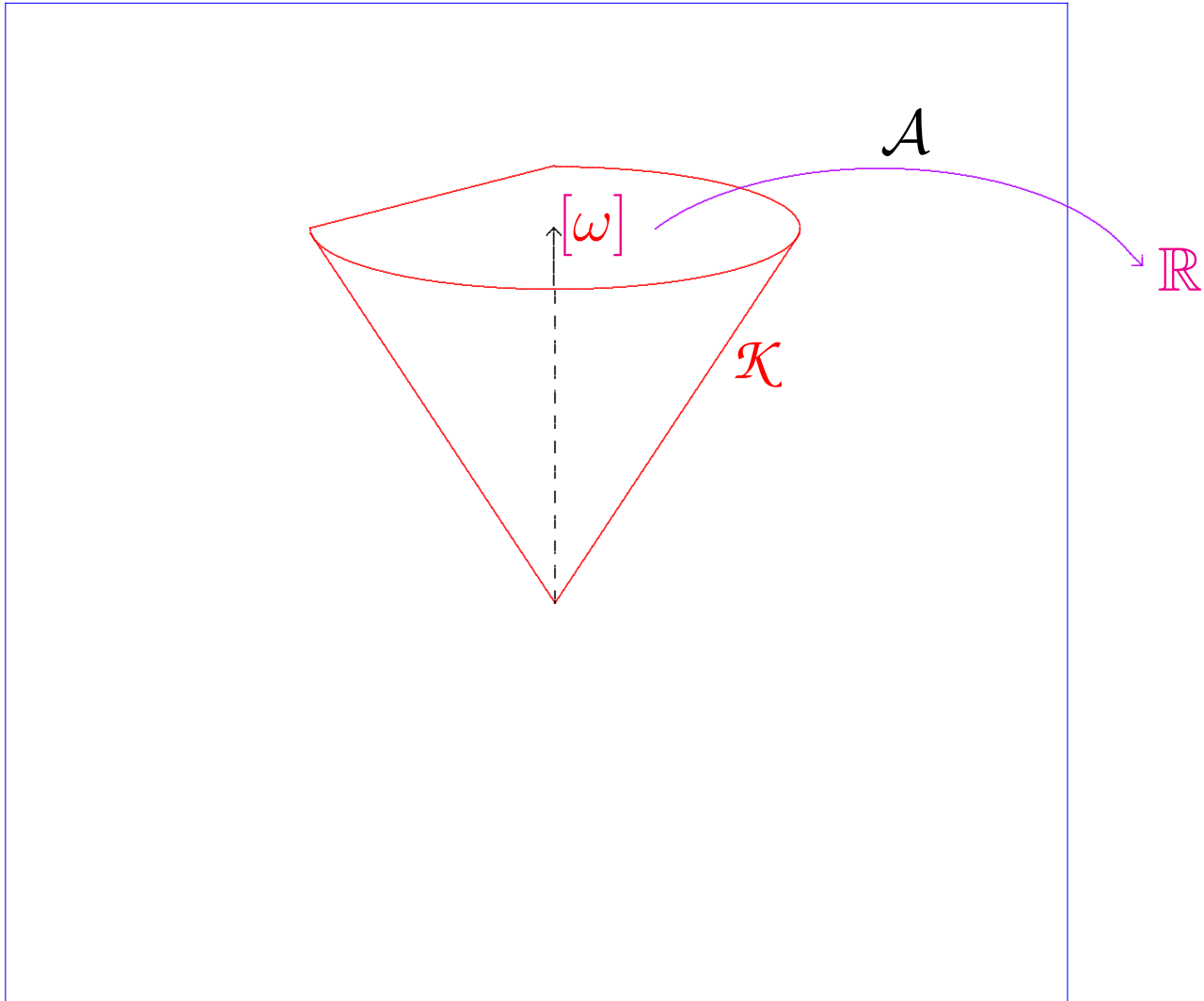
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Restriction of \mathcal{W}_+ to Kähler metrics?

On Kähler metrics,

$$\int |W_+|^2 d\mu = \int \frac{s^2}{24} d\mu$$

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Andrzej Derdziński : For Kähler metrics g ,

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Global implications?

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Moreover, each case actually occurs.

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III. $s < 0$ somewhere. Then

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If **not** Kähler-Einstein:

I. s is positive. Then

$(M, s^{-2}g)$ Einstein, $\lambda > 0$, $Hol = \mathbf{SO}(4)$.

II. s is zero. Then

(M, g, J) SFK, but not Ricci-flat.

III. s changes sign. Then

$(M, s^{-2}g)$ double Poincaré-Einstein. Here, $s = 0$ defines smooth connected \mathcal{Z}^3 , and $M - \mathcal{Z}$ has exactly two components.

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Main interest today:

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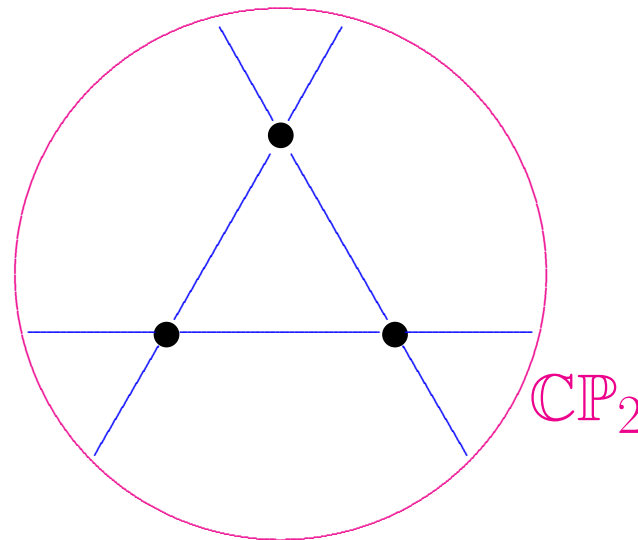
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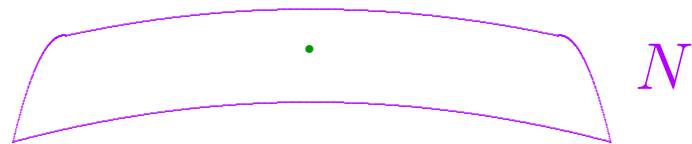
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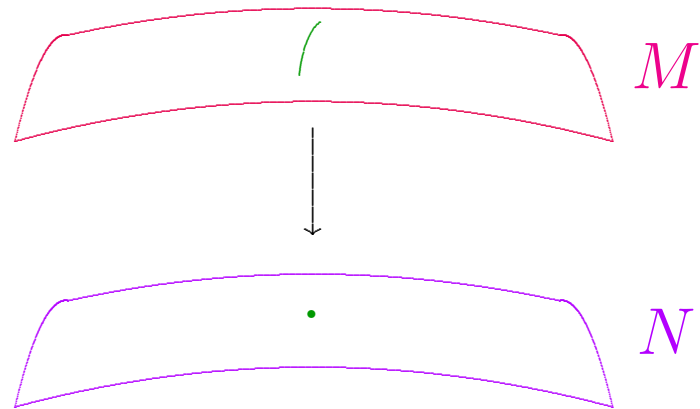
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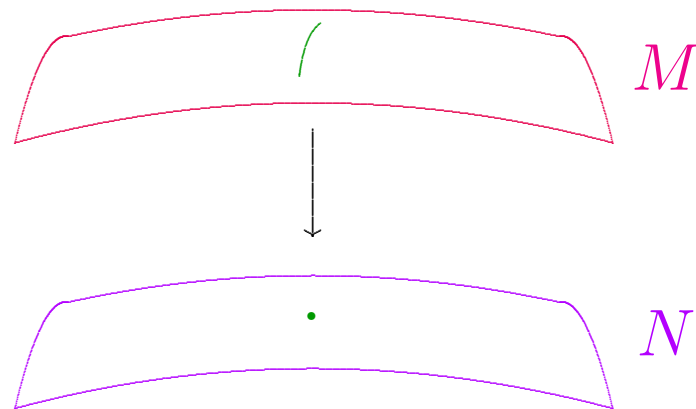
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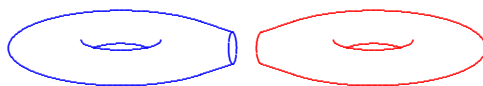
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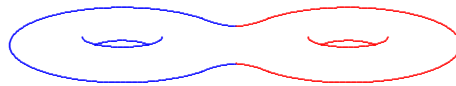
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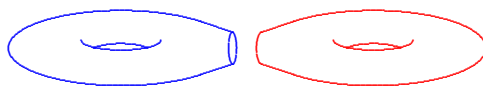
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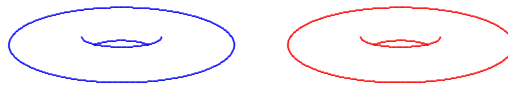
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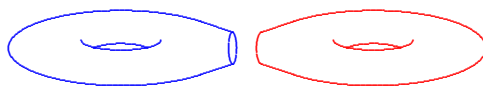
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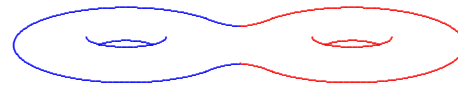
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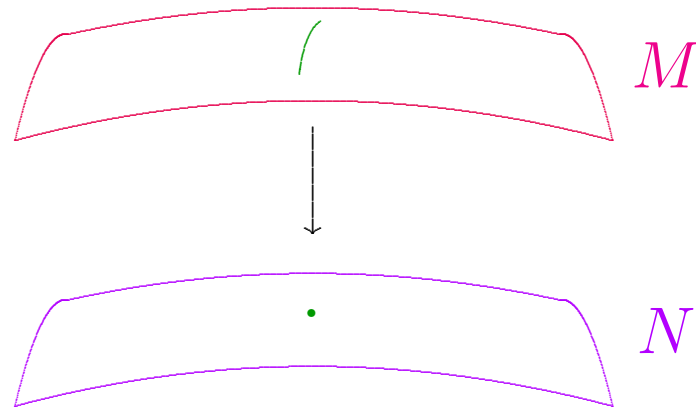
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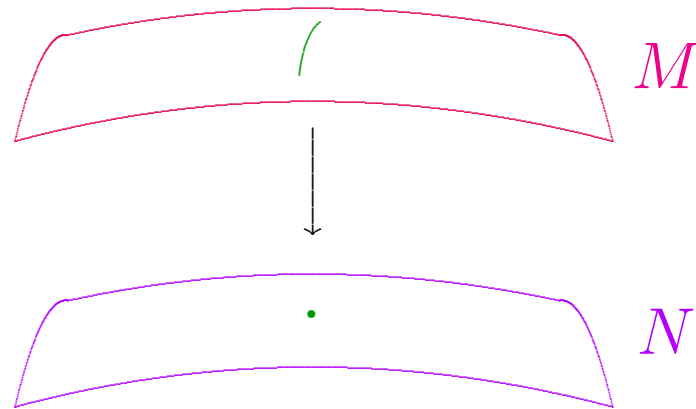


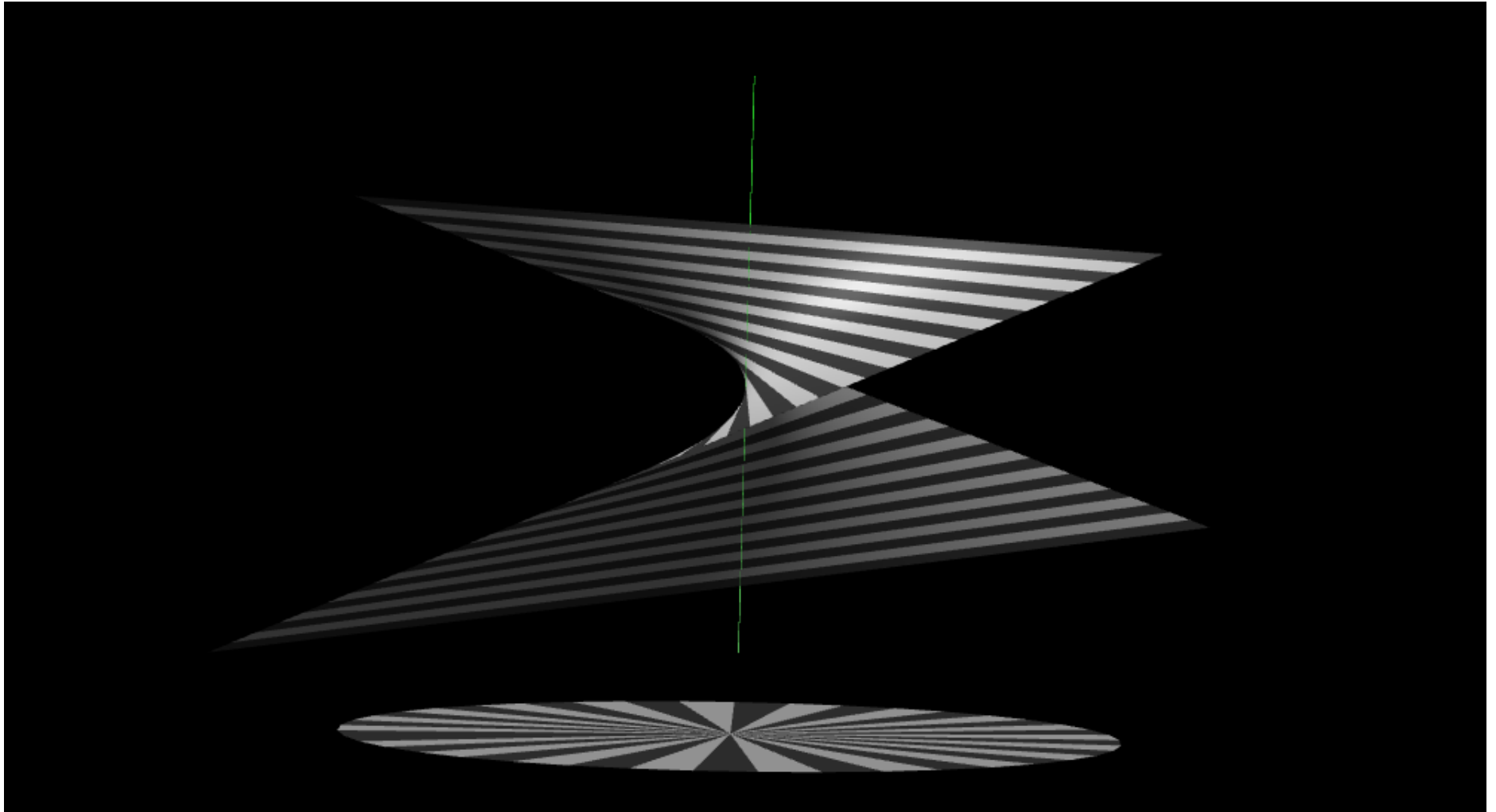
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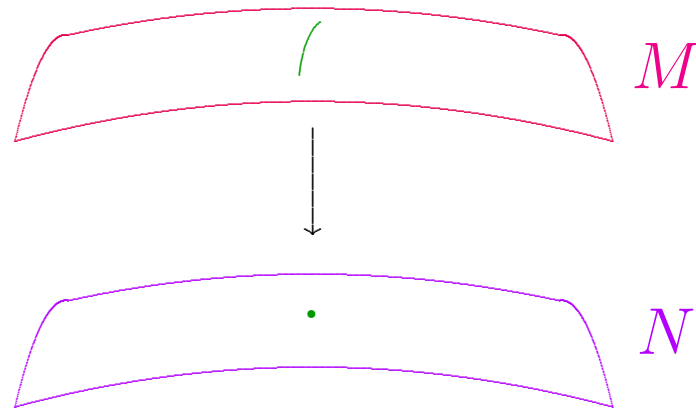


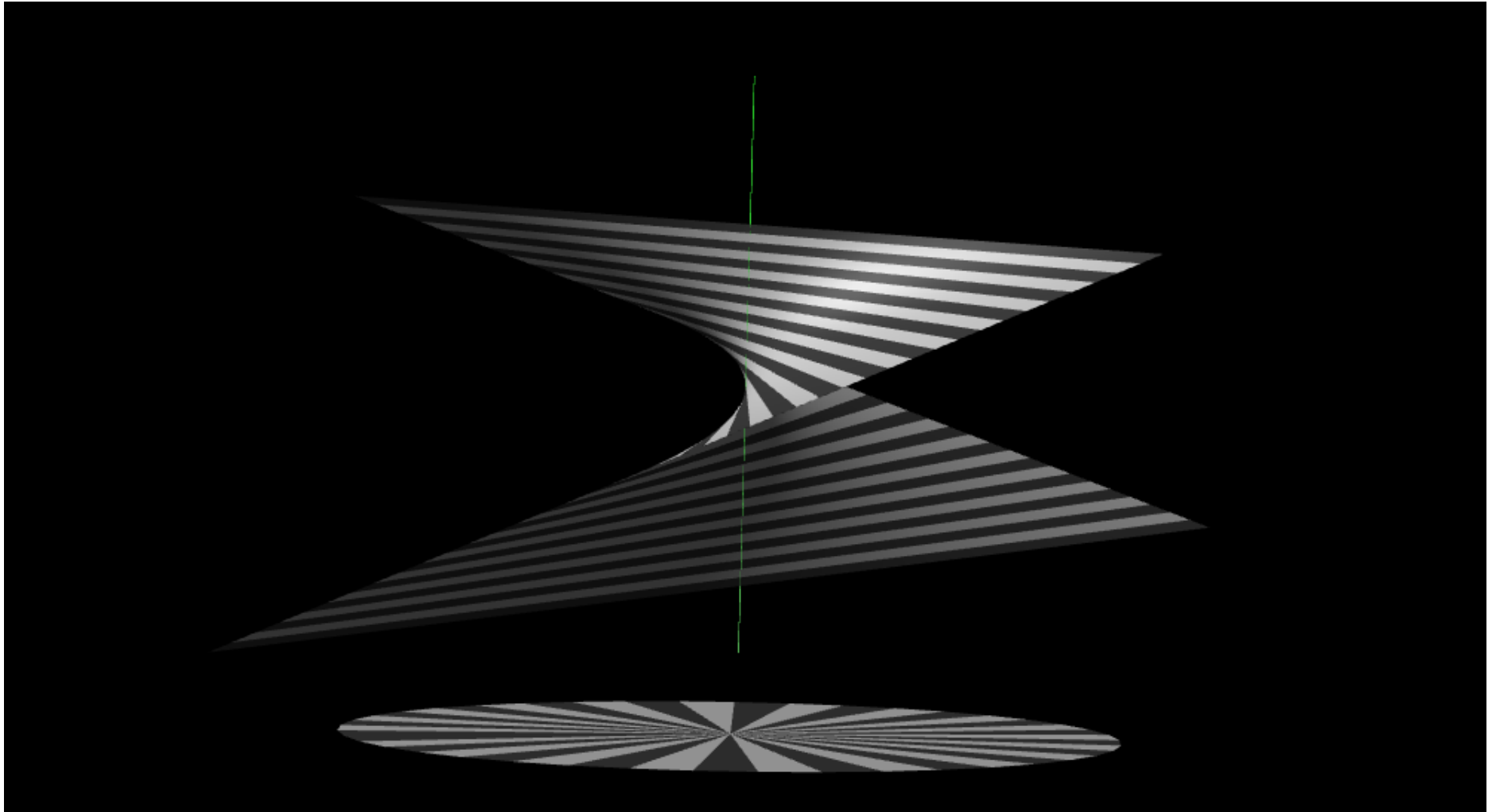
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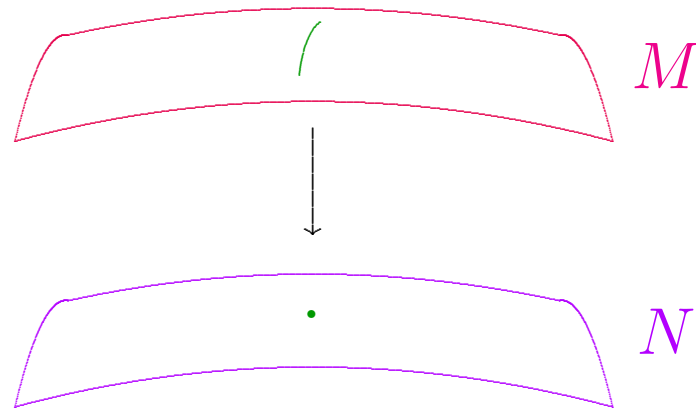


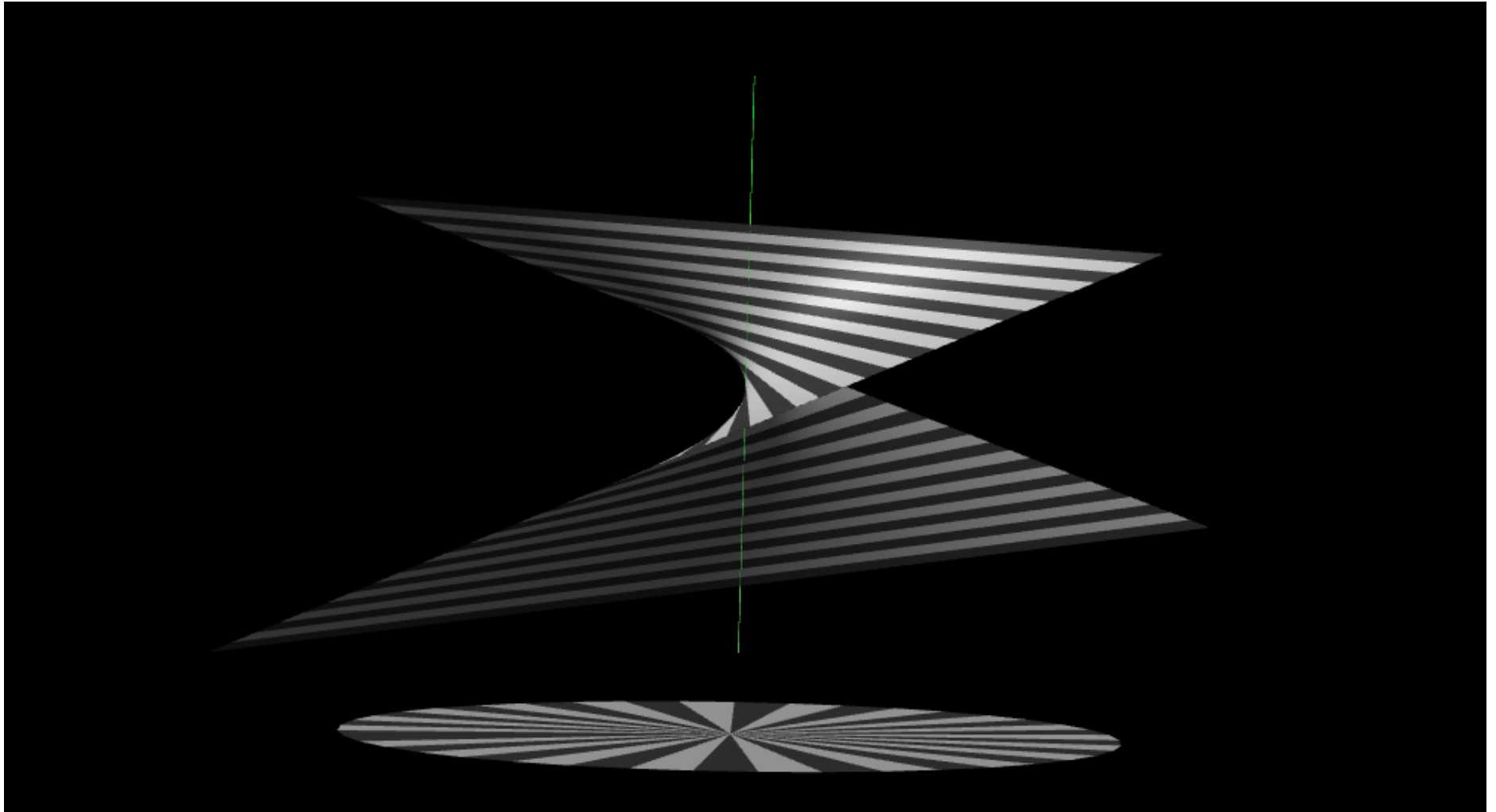
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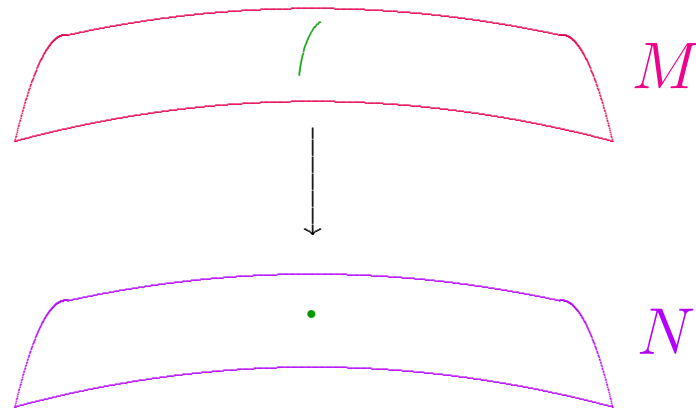


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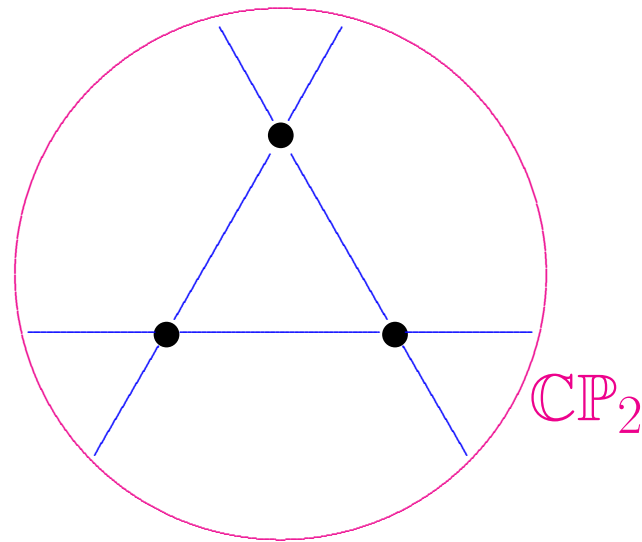


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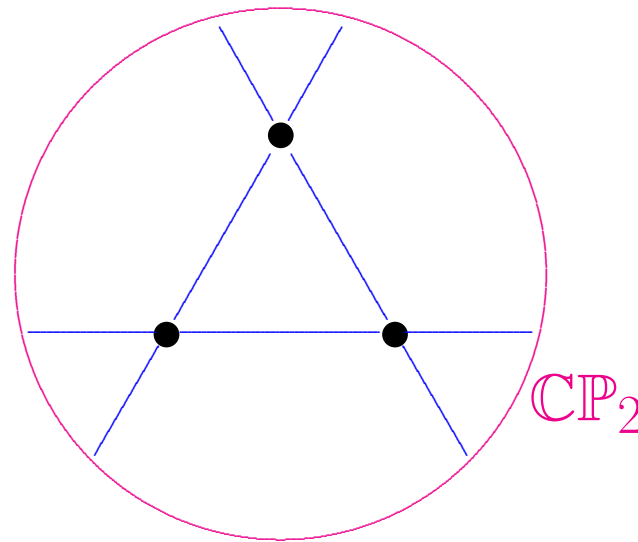
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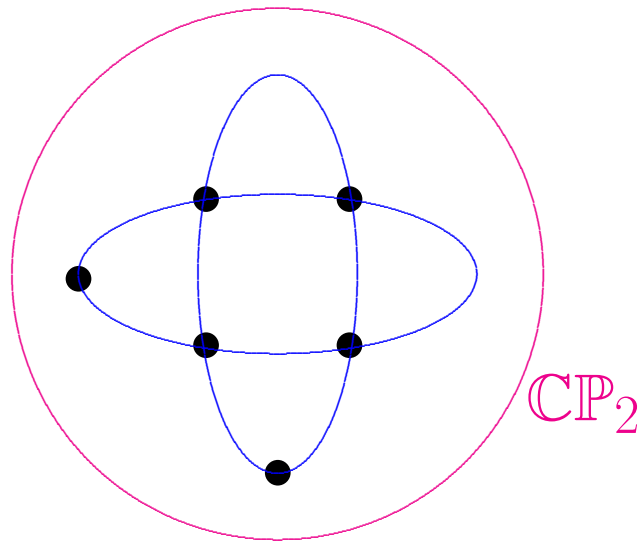


No 3 on a line,

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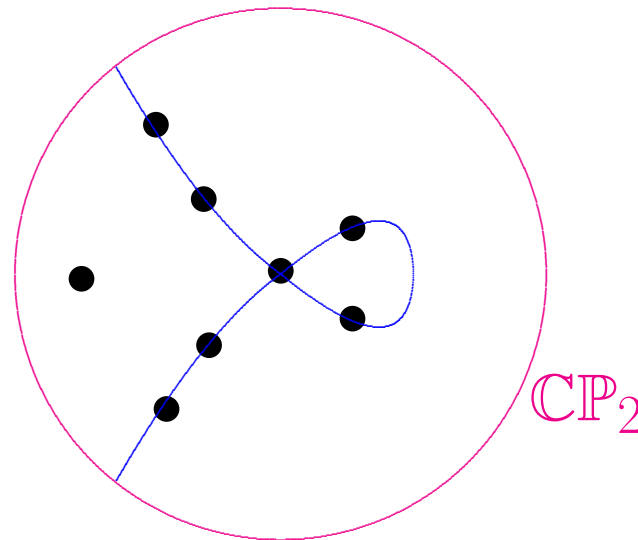


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Shorthand: “ $c_1 > 0$.”

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One reason this seems satisfying...

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But this is not needed in above result.

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Also true in non-K-E cases, among toric conformal classes $[g]$ of symplectic type.

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Understand all Einstein metrics on del Pezzos.

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Exactly one connected component of moduli space!

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Kähler $\implies \Lambda^+ = \mathbb{R}\omega \oplus \Re\Lambda^{2,0}$

$$W^+ = \text{trace-free part of } \begin{bmatrix} 0 & & \\ & 0 & \\ & & \frac{s}{4} \end{bmatrix}$$

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for these metrics & conformal rescalings:

$$g \rightsquigarrow h = f^2 g \implies \det(W^+) \rightsquigarrow f^{-6} \det(W^+).$$

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method also proves more general results.

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Claim: (M, h) compact Einstein $\implies J$ integrable.

Theorem B. *Let (M, h) be a simply-connected compact oriented Einstein 4-manifold, and suppose that its self-dual Weyl curvature*

$$W^+ : \Lambda^+ \rightarrow \Lambda^+$$

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at every point of M . Then h is conformal to an orientation-compatible Bach-flat extremal Kähler metric g with scalar curvature $s > 0$ on M .

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Corollary. *Every simply-connected compact oriented Einstein (M^4, h) with $\det(W^+) > 0$ is diffeomorphic to a del Pezzo surface. Conversely, every del Pezzo M^4 carries Einstein h with $\det(W^+) > 0$, and these sweep out exactly one connected component of moduli space $\mathcal{E}(M)$.*

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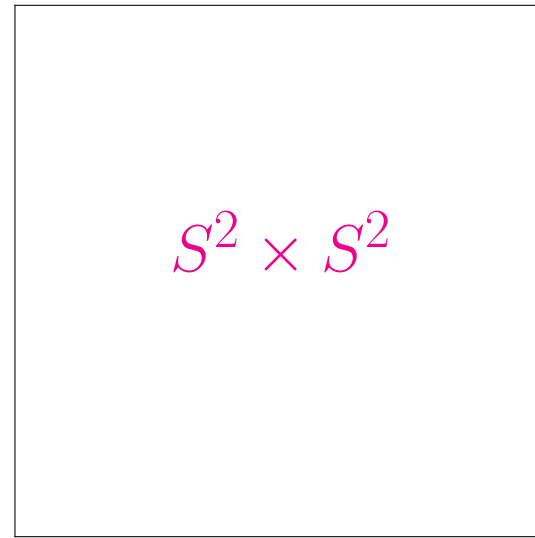
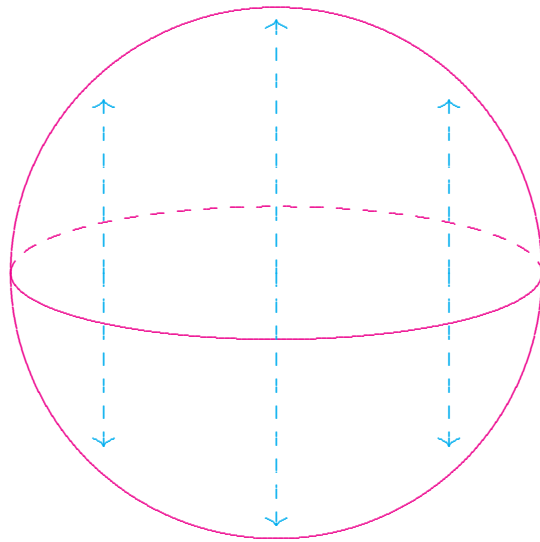
Simply connected hypothesis is essential!

Theorem C. *Let M be smooth compact oriented 4-manifold with $\pi_1 \neq 0$.*

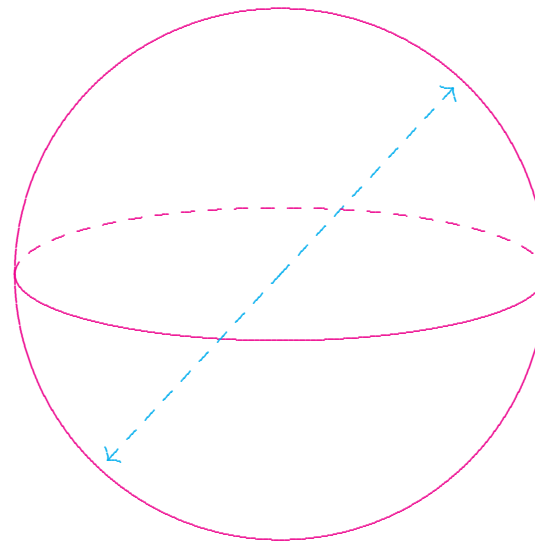
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Oriented spin 4-manifold
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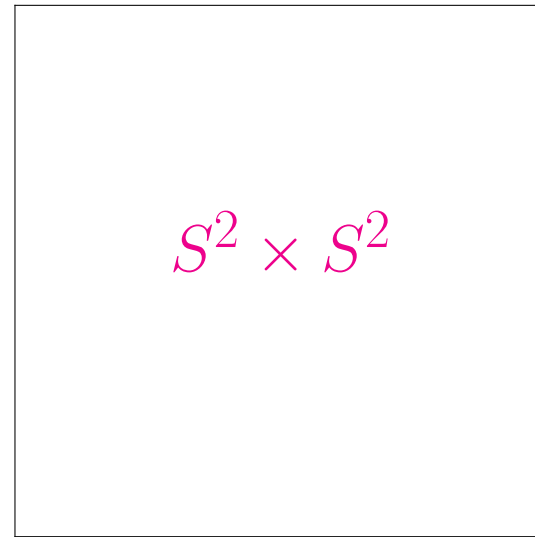
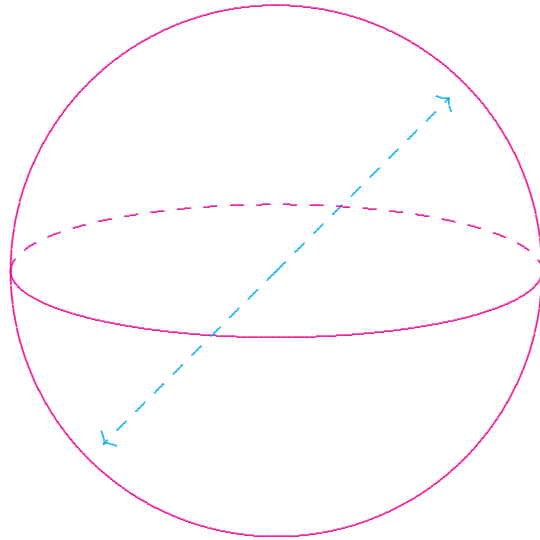


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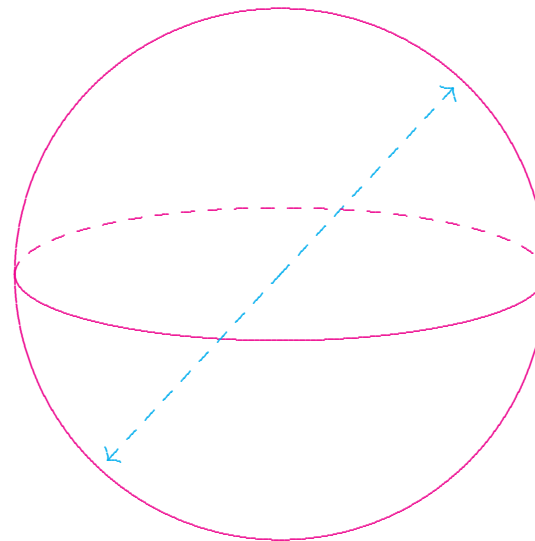
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Non-spin 4-manifold

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Why is $\mathcal{E}_{\det}(M) \subset \mathcal{E}(M)$ open and closed?

Open: $\det(W^+) > 0$.

Closed: $\det(W^+) = \frac{1}{3\sqrt{6}}|W^+|^3$ and $s \geq 0$.

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Theorem E. *Let (M, h) be a compact oriented Einstein 4-manifold. If*

$$\det(W^+) > -\frac{5\sqrt{2}}{21\sqrt{21}}|W^+|^3$$

everywhere on M , then actually $\det(W^+) > 0$. Consequently, all the results described remain true if we merely impose this ostensibly weaker hypothesis.

Some indication of the proof:

For clarity, let's just assume $\det(W^+) > 0 \dots$

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$$0 = \nabla^* \nabla (fW^+) + \frac{s}{2} fW^+ - 6fW^+ \circ W^+ + 2f|W^+|^2 I$$

with $\omega \otimes \omega$,

$$0 = \int_M [\langle \nabla^* \nabla (fW^+), \omega \otimes \omega \rangle + \dots] d\mu$$

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with $\omega \otimes \omega$, and integrate by parts. This yields:

$$0 = \int_M \left[\langle W^+, \nabla^* \nabla (\omega \otimes \omega) \rangle + \frac{s}{2} W^+ (\omega, \omega) - 6 |W^+ (\omega)|^2 + 2 |W^+|^2 |\omega|^2 \right] f \, d\mu$$

Let $\alpha \geq \beta \geq \gamma$ be eigenvalues of W^+ :

$$W^+ = \begin{bmatrix} \alpha & & \\ & \beta & \\ & & \gamma \end{bmatrix}$$

$$\alpha + \beta + \gamma = 0$$

$$\alpha > 0, \quad \gamma < 0, \quad \text{if } W^+ \neq 0$$

$$\det(W^+) = \alpha\beta\gamma$$

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Now choose $\omega \in \Gamma\Lambda^+$ so that

$$W_g^+(\omega) = \alpha \omega, \quad |\omega|_g \equiv \sqrt{2},$$

after at worst passing to double cover $\hat{M} \rightarrow M$.

$$0 = \int_{\hat{M}} \left[\langle W^+, \nabla^* \nabla (\omega \otimes \omega) \rangle + \frac{s}{2} W^+(\omega, \omega) - 6 |W^+(\omega)|^2 + 2 |W^+|^2 |\omega|^2 \right] f \, d\mu$$

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$$0 = \int_M \left[-2W^+(\nabla_e \omega, \nabla^e \omega) - 2W^+(\omega, \nabla^e \nabla_e \omega) \right. \\ \left. + \frac{s}{2} W^+(\omega, \omega) - 6|W^+(\omega)|^2 + 2|W^+|^2 |\omega|^2 \right] f \, d\mu$$

$$0 = \int_M \left[-2W^+(\nabla_e \omega, \nabla^e \omega) - 2\alpha \langle \omega, \nabla^e \nabla_e \omega \rangle \right. \\ \left. + \frac{s}{2} \alpha |\omega|^2 - 6\alpha^2 |\omega|^2 + 2|W^+|^2 |\omega|^2 \right] f \, d\mu$$

because

$$W_g^+(\omega) = \alpha \omega$$

$$0 = \int_M \left[-2W^+ (\nabla_e \omega, \nabla^e \omega) + 2\alpha \langle \omega, \nabla^* \nabla \omega \rangle \right. \\ \left. + \frac{s}{2} \alpha |\omega|^2 - 6\alpha^2 |\omega|^2 + 2|W^+|^2 |\omega|^2 \right] f \, d\mu$$

$$0 \geq \int_M \left[-2W^+(\nabla_e \omega, \nabla^e \omega) + 2\alpha \langle \omega, \nabla^* \nabla \omega \rangle \right. \\ \left. + \frac{s}{2} \alpha |\omega|^2 - 6\alpha^2 |\omega|^2 + 3\alpha^2 |\omega|^2 \right] f \, d\mu$$

because

$$|W_g^+|^2 \geq \frac{3}{2} \alpha^2$$

$$0 \geq \int_M \left[-2W^+(\nabla_e \omega, \nabla^e \omega) + 2\alpha \langle \omega, \nabla^* \nabla \omega \rangle \right. \\ \left. + \frac{s}{2} \alpha |\omega|^2 - 3\alpha^2 |\omega|^2 \right] f \, d\mu$$

$$|\omega|_g^2 = 2 \implies (\nabla_e \omega) \perp \omega$$

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$$|\omega|_g^2 = 2 \implies (\nabla_e \omega) \perp \omega$$

$$\det(W^+) > 0 \implies W^+ \sim \begin{bmatrix} + & & \\ & - & \\ & & - \end{bmatrix}$$

$$0 \geq \int_M \left[-2W^+(\nabla_e \omega, \nabla^e \omega) + 2\alpha \langle \omega, \nabla^* \nabla \omega \rangle + \frac{s}{2} \alpha |\omega|^2 - 3\alpha^2 |\omega|^2 \right] f \, d\mu$$

$$|\omega|_g^2 = 2 \implies (\nabla_e \omega) \perp \omega$$

$$\det(W^+) > 0 \implies W^+(\nabla_e \omega, \nabla^e \omega) \leq 0$$

$$0 \geq \int_M \left[\begin{aligned} &2\alpha \langle \omega, \nabla^* \nabla \omega \rangle \\ &+ \frac{s}{2} \alpha |\omega|^2 - 3\alpha^2 |\omega|^2 \end{aligned} \right] f \, d\mu$$

$$|\omega|_g^2 = 2 \implies (\nabla_e \omega) \perp \omega$$

$$\det(W^+) > 0 \implies -W^+(\nabla_e \omega, \nabla^e \omega) \geq 0$$

$$0 \geq \int_M \left[2\alpha \langle \omega, \nabla^* \nabla \omega \rangle + \frac{s}{2} \alpha |\omega|^2 - 3\alpha^2 |\omega|^2 \right] f \, d\mu$$

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But

$$\alpha f \equiv 1$$

$$0 \geq \int_M \left[2\langle \omega, \nabla^* \nabla \omega \rangle + \frac{s}{2} \alpha |\omega|^2 - 3 |\omega|^2 \alpha \right] d\mu$$

$$0 \geq \int_M \left[2\langle \omega, \nabla^* \nabla \omega \rangle - 3W^+(\omega, \omega) + \frac{s}{2} |\omega|^2 \right] d\mu$$

$$0 \geq \int_M \left[\frac{1}{2} |\nabla \omega|^2 + \frac{3}{2} \langle \omega, \left(\nabla^* \nabla - 2W^+ + \frac{s}{3} \right) \omega \rangle \right] d\mu$$

$$0 \geq \int_M \left[\frac{1}{2} |\nabla \omega|^2 + \frac{3}{2} \langle \omega, (d + d^*)^2 \omega \rangle \right] d\mu$$

Because

$$(d + d^*)^2 = \nabla^* \nabla - 2W^+ + \frac{s}{3}$$

on $\Gamma\Lambda^+$.

$$0 \geq \frac{1}{2} \int_M |\nabla \omega|^2 d\mu + 3 \int_M |d\omega|^2 d\mu$$

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So $\nabla \omega \equiv 0$, and g is Kähler!





Thanks for the invitation!



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It's such a pleasure to be here!